



University of Dundee

Locally periodic unfolding method and two-scale convergence on surfaces of locally periodic microstructures

Ptashnyk, Mariya

Published in:
Multiscale Modeling and Simulation: A SIAM Interdisciplinary Journal

DOI:
[10.1137/140978405](https://doi.org/10.1137/140978405)

Publication date:
2015

Document Version
Peer reviewed version

[Link to publication in Discovery Research Portal](#)

Citation for published version (APA):
Ptashnyk, M. (2015). Locally periodic unfolding method and two-scale convergence on surfaces of locally periodic microstructures. *Multiscale Modeling and Simulation: A SIAM Interdisciplinary Journal*, 13(3), 1061-1105. <https://doi.org/10.1137/140978405>

General rights

Copyright and moral rights for the publications made accessible in Discovery Research Portal are retained by the authors and/or other copyright owners and it is a condition of accessing publications that users recognise and abide by the legal requirements associated with these rights.

Take down policy

If you believe that this document breaches copyright please contact us providing details, and we will remove access to the work immediately and investigate your claim.

LOCALLY PERIODIC UNFOLDING METHOD AND TWO-SCALE CONVERGENCE ON SURFACES OF LOCALLY PERIODIC MICROSTRUCTURES*

MARIYA PTASHNYK [†]

Abstract. In this paper we generalize the periodic unfolding method and the notion of two-scale convergence on surfaces of periodic microstructures to locally-periodic situations. The methods that we introduce allow us to consider a wide range of non-periodic microstructures, especially to derive macroscopic equations for problems posed in domains with perforations distributed non-periodically. Using the methods of locally periodic two-scale convergence (l-t-s) on oscillating surfaces and the locally periodic (l-p) boundary unfolding operator, we are able to analyze differential equations defined on boundaries of non-periodic microstructures and consider non-homogeneous Neumann conditions on the boundaries of perforations, distributed non-periodically.

1. Introduction. Many natural and man-made composite materials comprise non-periodic microscopic structures, e.g. fibrous microstructures in heart muscles [23, 48], exoskeletons [27], industrial filters [52], or space-dependent perforations in concrete [50]. An important special case of non-periodic microstructures is that of the so-called locally-periodic microstructures, where spatial changes are observed on a scale smaller than the size of the domain under consideration, but larger than the characteristic size of the microstructure. For many locally-periodic microstructures spatial changes cannot be represented by periodic functions depending on slow and fast variables, e.g. plywood-like structures of gradually rotated planes of parallel aligned fibers [13]. Thus, in these situations the standard two-scale convergence and periodic unfolding method cannot be applied. Hence, for a multiscale analysis of problems posed in domains with non-periodic perforations, in this paper we extend the periodic unfolding method and two-scale convergence on oscillating surfaces to locally-periodic situations (see Definition 3.4, Definition 3.2, Definition 3.3, and Definition 3.5). These generalizations are motivated by the locally-periodic two-scale convergence introduced in [49].

Two-scale convergence on surfaces of periodic microstructures was first introduced in [5, 43]. An extension of two-scale convergence associated with a fixed periodic Borel measure was considered in [55]. The unfolding operator maps functions defined on perforated domains, depending on small parameter ε , onto functions defined on the whole fixed domain, see [20, 22] and references therein. This helps to overcome one of the difficulties of perforated domains which is the use of extension operators. Using the boundary unfolding operator we can prove convergence results for nonlinear equations posed on oscillating boundaries of microstructures [22, 24, 36, 46]. The unfolding method is also an efficient tool to derive error estimates, see e.g. [28, 31, 32, 33, 47].

The main novelty of this article is the derivation of new techniques for the multiscale analysis of non-linear problems posed in domains with non-periodic perforations and on the surfaces of non-periodic microstructures. The l-p unfolding operator allows us to analyze nonlinear differential equations posed on domains with non-periodic perforations. The l-t-s convergence on oscillating surfaces and the l-p boundary unfolding operator allow us to show strong convergence for sequences defined on oscillat-

*This research was supported by EPSRC First Grant “Multiscale modelling and analysis of mechanical properties of plant cells and tissues”.

[†] Division of Mathematics, University of Dundee, DD1 4HN, Scotland, UK, (mptashnyk@maths.dundee.ac.uk).

ing boundaries of non-periodic microstructures and to derive macroscopic equations for nonlinear equations defined on boundaries of non-periodic microstructures. Until now, this was not possible using existing methods.

The paper is structured as follows. First, in Section 2, we present a mathematical description of locally periodic microstructures and state the definition of a locally periodic approximation for a function $\psi \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$. In Section 3 we introduce all the main definitions of the paper, i.e. the notion of a l-p unfolding operator, two-scale convergence for sequences defined on oscillating boundaries of locally periodic microstructures, and the l-p boundary unfolding operator. The main results are summarized in Section 4. The central results of this paper are convergence results for sequences bounded in L^p and $W^{1,p}$, with $p \in (1, \infty)$ (see Theorems 4.1, 4.2, 4.3, and 4.4). The proofs of the main results for the l-p unfolding operator are presented in Section 5. The properties of the decomposition of a $W^{1,p}$ -function with one part describing the macroscopic behavior and another part of order ε , are shown in Section 6. The proofs of the main results for the l-p unfolding operator in perforated domains are given in Section 7. The convergence results for locally-periodic two-scale convergence on oscillating surfaces and the l-p boundary unfolding operator are proved in Section 8. In Section 9 we apply the l-p unfolding operator to derive macroscopic problems for microscopic models of signaling processes in cell tissues comprising locally-periodic microstructures. As examples of tissues with locally-periodic microstructures we consider plant tissues, epithelial tissues, and non-periodic fibrous structure of heart tissue.

There are some existing results on the homogenization of problems posed on locally-periodic media. The homogenization of a heat-conductivity problem defined in domains with non-periodic microstructure consisting of spherical balls was studied in [14] using the Murat-Tartar H -convergence method [42], and in [3] by applying the θ -2 convergence. The non-periodic distribution of balls is given by a C^2 -diffeomorphism θ , transforming the centers of the balls. Estimates for a numerical approximation of this problem were derived in [53]. The notion of a Young measure was used in [38] to extend the concept of periodic two-scale convergence and to define the so-called *scale convergence*. The definition of scale convergence was motivated by the derivation of the Γ -limit for a sequence of nonlinear energy functionals involving non-periodic oscillations. Formal asymptotic expansions and the technique of two-scale convergence defined for periodic test functions, see e.g. [4, 44], were used to derive macroscopic equations for models posed on domains with locally periodic perforations, i.e. domains consisting of periodic cells with smoothly changing perforations [9, 17, 18, 37, 39, 45]. The H -convergence method [12, 13], the asymptotic expansion method [8], and the method of locally-periodic two-scale (l-t-s) convergence [49] were applied to analyze microscopic models posed on domains consisting of non-periodic fibrous materials. The optimization of the elastic properties of a material with locally-periodic microstructure was considered in [6, 7].

To illustrate the difference between the formulation of non-periodic microstructure by using periodic functions and the locally-periodic formulation of the problem, we consider a plywood-like structure, given as the superposition of gradually rotated planes of aligned parallel fibers. We consider layers of cylindrical fibers of radius εa orthogonal to the x_3 -axis and rotated around the x_3 -axis by an angle γ , constant in each layer and changing from one layer to another, see Fig.1. To describe the difference in the material properties of fibers and the inter-fibre space with the help of a periodic function, we define a function

$$(1.1) \quad A^\varepsilon(x) = A_1 \tilde{\eta}(R(\gamma(x_3))x/\varepsilon) + A_2 \left[1 - \tilde{\eta}(R(\gamma(x_3))x/\varepsilon) \right],$$

FIG. 1. Schematic representation of slow rotating and fast rotating plywood-like structures.

where A_1, A_2 are constant tensors and $\tilde{\eta}$ is the characteristic functions of a fibre of radius a in the direction of x_1 -axis, i.e.

$$(1.2) \quad \tilde{\eta}(y) = \begin{cases} 1 & \text{for } |\hat{y} - (1/2, 1/2)| \leq a, \\ 0 & \text{for } |\hat{y} - (1/2, 1/2)| > a, \end{cases}$$

and extended \hat{Y} -periodic to the whole \mathbb{R}^3 , with $a < 1/2$, $\hat{y} = (y_2, y_3)$, $Y = [0, 1]^3$, and $\hat{Y} = [0, 1]^2$. The inverse of the rotation matrix around the x_3 -axes with rotation angle α with the x_1 -axis is defined as

$$(1.3) \quad R(\alpha) = \begin{pmatrix} \cos(\alpha) & \sin(\alpha) & 0 \\ -\sin(\alpha) & \cos(\alpha) & 0 \\ 0 & 0 & 1 \end{pmatrix},$$

and $\gamma \in C^1(\mathbb{R})$ is a given function, such that $0 \leq \gamma(s) \leq \pi$ for all $s \in \mathbb{R}$. Then, considering for example an elliptic problem with a diffusion coefficient or elasticity tensor in the form (1.1) and using a change of variables $\tilde{x} = R(\gamma(x_3))x$, we can apply periodic homogenization techniques to derive corresponding macroscopic equations (see [10, 12] for details). However, in the representation of the microscopic structure by (1.1), every point of a fibre is rotated differently and the cylindrical structure of the fibers is deformed. Hence, A^ε represent the properties of a material with a different microstructure than the plywood-like structure, and for a correct representation of a plywood-like structure, a locally-periodic formulation of the microscopic problem is essential. Also, applying periodic homogenization techniques we obtain effective macroscopic coefficients different from the one obtained by using methods of locally-periodic homogenization (see [13, 49] for more details).

To define the characteristic function of the domain occupied by fibers in a domain with a locally-periodic plywood-like structure, we divide \mathbb{R}^3 in layers $L_k^\varepsilon = \mathbb{R}^2 \times ((k-1)\varepsilon^r, k\varepsilon^r)$ of height ε^r and perpendicular to the x_3 -axis, where $k \in \mathbb{Z}$ and $0 < r < 1$. In each L_k^ε we choose an arbitrary point $x_k^\varepsilon \in L_k^\varepsilon$. Using the locally-periodic approximation of $\eta \in C(\bar{\Omega}, L_{\text{per}}^\infty(Y_x))$, with $\eta(x, y) = \tilde{\eta}(R(x)y)$ for $x \in \Omega$ and $y \in Y_x$, given by

$$(\mathcal{L}^\varepsilon \eta)(x) = \sum_{k \in \mathbb{Z}} \tilde{\eta}(R(\gamma(x_{k,3}^\varepsilon))x/\varepsilon) \chi_{L_k^\varepsilon}(x) \quad \text{for } x \in \Omega,$$

the characteristic function of the domain occupied by fibers is given by

$$(1.4) \quad \chi_{\Omega_f^\varepsilon}(x) = \chi_\Omega(x)(\mathcal{L}^\varepsilon \eta)(x).$$

Here $\tilde{\eta} \in L_{\text{per}}^\infty(Y)$ is as in (1.2) and $Y_x = R^{-1}(\gamma(x_3))Y$. For a microstructure composed of fast rotating planes of parallel aligned fibrous, see Fig. 1, we consider an approximation by locally-periodic plywood-like structure with shifted periodicity $D(x)Y = R^{-1}(x)W(x)Y$, see [13, 49] for more details.

2. Locally periodic microstructures and locally periodic perforated domains. In this section we give a mathematical formulation of locally periodic microstructures. We also define the approximation of functions, where the periodicity with respect to the fast variable is dependent on the slow variable, by locally-periodic functions, i.e. periodic in subdomains smaller than the domain under consideration but larger than the representative size of the microstructure.

Let $\Omega \subset \mathbb{R}^d$ be a bounded Lipschitz domain. For each $x \in \mathbb{R}^d$ we consider a transformation matrix $D(x) \in \mathbb{R}^{d \times d}$ and its inverse $D^{-1}(x)$, such that $D, D^{-1} \in \text{Lip}(\mathbb{R}^d; \mathbb{R}^{d \times d})$ and $0 < D_1 \leq |\det D(x)| \leq D_2 < \infty$ for all $x \in \bar{\Omega}$. We consider the continuous family of parallelepipeds $Y_x = D_x Y$ on $\bar{\Omega}$, where $Y = (0, 1)^d$ is the ‘unit cell’ and denote $D_x := D(x)$ and $D_x^{-1} := D^{-1}(x)$.

For $\varepsilon > 0$, in a manner similar to [14, 49], we consider the partition covering of Ω by a family of open non-intersecting cubes $\{\Omega_n^\varepsilon\}_{1 \leq n \leq N_\varepsilon}$ of side ε^r , with $0 < r < 1$,

$$\Omega \subset \bigcup_{n=1}^{N_\varepsilon} \bar{\Omega}_n^\varepsilon \quad \text{and} \quad \Omega_n^\varepsilon \cap \Omega \neq \emptyset.$$

For arbitrary chosen fixed points $x_n^\varepsilon, \tilde{x}_n^\varepsilon \in \Omega_n^\varepsilon \cap \Omega$ we consider a covering of Ω_n^ε by parallelepipeds $\varepsilon D_{x_n^\varepsilon} Y$

$$\Omega_n^\varepsilon \subset \tilde{x}_n^\varepsilon + \bigcup_{\xi \in \Xi_n^\varepsilon} \varepsilon D_{x_n^\varepsilon} (\bar{Y} + \xi), \quad \text{where } \Xi_n^\varepsilon = \{\xi \in \mathbb{Z}^d : \tilde{x}_n^\varepsilon + \varepsilon D_{x_n^\varepsilon} (Y + \xi) \cap \Omega_n^\varepsilon \neq \emptyset\},$$

with $D_{x_n^\varepsilon} = D(x_n^\varepsilon)$ and $1 \leq n \leq N_\varepsilon$. For each $n = 1, \dots, N_\varepsilon$, \tilde{x}_n^ε is a fixed shift in the representation of the microscopic structure of Ω_n^ε . Often we can consider $\tilde{x}_n^\varepsilon = \varepsilon D_{x_n^\varepsilon} \xi$ for some $\xi \in \mathbb{Z}^d$.

We consider the space $C(\bar{\Omega}; C_{\text{per}}(Y_x))$ given in a standard way, i.e. for any $\tilde{\psi} \in C(\bar{\Omega}; C_{\text{per}}(Y))$ the relation $\psi(x, y) = \tilde{\psi}(x, D_x^{-1}y)$ with $x \in \Omega$ and $y \in Y_x$ yields $\psi \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$. In the same way the spaces $L^p(\Omega; C_{\text{per}}(Y_x))$, $L^p(\Omega; L_{\text{per}}^q(Y_x))$ and $C(\bar{\Omega}; L_{\text{per}}^q(Y_x))$, for $1 \leq p \leq \infty$, $1 \leq q < \infty$, are defined.

To describe locally-periodic microscopic properties of a composite material and to specify test functions associated with the locally-periodic microstructure of a material, as well as for the definition of the locally-periodic two-scale convergence, we shall consider a locally-periodic approximation of functions with space-dependent periodicity, functions in $C(\bar{\Omega}; C_{\text{per}}(Y_x))$, $L^p(\Omega; C_{\text{per}}(Y_x))$, or $C(\bar{\Omega}; L_{\text{per}}^q(Y_x))$. The locally-periodic approximated function is $Y_{x_n^\varepsilon}$ -periodic in each subdomain Ω_n^ε , with $n = 1, \dots, N_\varepsilon$, and is related to a test function associated with the periodic structure of Ω_n^ε . Since the microscopic structure of Ω_n^ε is represented by a union of periodicity cells $\varepsilon Y_{x_n^\varepsilon}$ shifted by a fixed point $\tilde{x}_n^\varepsilon \in \Omega_n^\varepsilon \cap \Omega$, with $n = 1, \dots, N_\varepsilon$, this shift is also reflected in the definition of the locally-periodic approximation.

Often coefficients in a microscopic model posed in a domain with locally-periodic microstructure depend only on the microscopic fast variables x/ε and the points $x_n^\varepsilon, \tilde{x}_n^\varepsilon \in \Omega_n^\varepsilon \cap \Omega$, describing the periodic microstructure in each Ω_n^ε , with $n = 1, \dots, N_\varepsilon$, and are independent of the macroscopic slow variables x . To define such functions we shall introduce a notion of a locally-periodic approximation $\mathcal{L}_0^\varepsilon$ of a function

$\psi \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ (or in $L^p(\Omega; C_{\text{per}}(Y_x))$, $C(\bar{\Omega}; L^q_{\text{per}}(Y_x))$). In each Ω_n^ε the function $\mathcal{L}_0^\varepsilon(\psi)$ is $Y_{x_n^\varepsilon}$ -periodic and depend only on the fast variables x/ε . This specific locally-periodic approximation is important for the derivation of macroscopic equations for a microscopic problem with coefficients discontinuous with respect to the fast variable, since for $\psi \in C(\bar{\Omega}; L^p(Y_x))$ we have that $\mathcal{L}_0^\varepsilon(\psi)$ converges strongly locally-periodic (l-p) two-scale, see [49].

As a locally periodic (l-p) approximation of ψ we name $\mathcal{L}^\varepsilon : C(\bar{\Omega}; C_{\text{per}}(Y_x)) \rightarrow L^\infty(\Omega)$ given by

$$(2.1) \quad (\mathcal{L}^\varepsilon \psi)(x) = \sum_{n=1}^{N_\varepsilon} \tilde{\psi}\left(x, \frac{D_{x_n^\varepsilon}^{-1}(x - \tilde{x}_n^\varepsilon)}{\varepsilon}\right) \chi_{\Omega_n^\varepsilon}(x) \quad \text{for } x \in \Omega.$$

We consider also the map $\mathcal{L}_0^\varepsilon : C(\bar{\Omega}; C_{\text{per}}(Y_x)) \rightarrow L^\infty(\Omega)$ defined for $x \in \Omega$ as

$$(\mathcal{L}_0^\varepsilon \psi)(x) = \sum_{n=1}^{N_\varepsilon} \psi\left(x_n^\varepsilon, \frac{x - \tilde{x}_n^\varepsilon}{\varepsilon}\right) \chi_{\Omega_n^\varepsilon}(x) = \sum_{n=1}^{N_\varepsilon} \tilde{\psi}\left(x_n^\varepsilon, \frac{D_{x_n^\varepsilon}^{-1}(x - \tilde{x}_n^\varepsilon)}{\varepsilon}\right) \chi_{\Omega_n^\varepsilon}(x).$$

If we choose $\tilde{x}_n^\varepsilon = D_{x_n^\varepsilon} \varepsilon \xi$ for some $\xi \in \mathbb{Z}^d$, then the periodicity of $\tilde{\psi}$ implies

$$(\mathcal{L}^\varepsilon \psi)(x) = \sum_{n=1}^{N_\varepsilon} \tilde{\psi}\left(x, \frac{D_{x_n^\varepsilon}^{-1}x}{\varepsilon}\right) \chi_{\Omega_n^\varepsilon}(x) \quad \text{and} \quad (\mathcal{L}_0^\varepsilon \psi)(x) = \sum_{n=1}^{N_\varepsilon} \tilde{\psi}\left(x_n^\varepsilon, \frac{D_{x_n^\varepsilon}^{-1}x}{\varepsilon}\right) \chi_{\Omega_n^\varepsilon}(x)$$

for $x \in \Omega$. In the following, we shall consider the case $\tilde{x}_n^\varepsilon = \varepsilon D_{x_n^\varepsilon} \xi$, with $\xi \in \mathbb{Z}^d$. However, all results hold for arbitrary chosen $\tilde{x}_n^\varepsilon \in \Omega_n^\varepsilon$ with $n = 1, \dots, N_\varepsilon$, see [49]. In a similar way we define $\mathcal{L}^\varepsilon \psi$ and $\mathcal{L}_0^\varepsilon \psi$ for ψ in $C(\bar{\Omega}; L^q_{\text{per}}(Y_x))$ or $L^p(\Omega; C_{\text{per}}(Y_x))$.

The locally-periodic approximation reflects the microscopic properties of Ω , where in each Ω_n^ε the microstructure is represented by a ‘unit cell’ $Y_{x_n^\varepsilon} = D_{x_n^\varepsilon} Y$ for arbitrary fixed $x_n^\varepsilon \in \Omega_n^\varepsilon$, see Figs. 1 and 2.

In the context of admissible test functions in weak formulations of partial differential equations, we define a regular approximation of $\mathcal{L}^\varepsilon \psi$ by

$$(\mathcal{L}_\rho^\varepsilon \psi)(x) = \sum_{n=1}^{N_\varepsilon} \tilde{\psi}\left(x, \frac{D_{x_n^\varepsilon}^{-1}x}{\varepsilon}\right) \phi_{\Omega_n^\varepsilon}(x) \quad \text{for } x \in \Omega,$$

where $\phi_{\Omega_n^\varepsilon}$ are approximations of $\chi_{\Omega_n^\varepsilon}$ such that $\phi_{\Omega_n^\varepsilon} \in C_0^\infty(\Omega_n^\varepsilon)$ and

$$(2.2) \quad \sum_{n=1}^{N_\varepsilon} |\phi_{\Omega_n^\varepsilon} - \chi_{\Omega_n^\varepsilon}| \rightarrow 0 \text{ in } L^2(\Omega), \quad \|\nabla^m \phi_{\Omega_n^\varepsilon}\|_{L^\infty(\mathbb{R}^d)} \leq C \varepsilon^{-\rho m} \text{ for } 0 < r < \rho < 1,$$

see e.g. [12, 14, 49]. In the definition of the l-p unfolding operator we shall use subdomains of Ω_n^ε given by unit cells $\varepsilon Y_{x_n^\varepsilon}$ that are completely included in $\Omega_n^\varepsilon \cap \Omega$, see Fig. 2.

$$(2.3) \quad \hat{\Omega}^\varepsilon = \bigcup_{n=1}^{N_\varepsilon} \hat{\Omega}_n^\varepsilon, \quad \text{with } \hat{\Omega}_n^\varepsilon = \text{Int}\left(\bigcup_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon D_{x_n^\varepsilon}(\bar{Y} + \xi)\right) \quad \text{and} \quad \Lambda^\varepsilon = \bigcup_{n=1}^{N_\varepsilon} \Lambda_n^\varepsilon \cap \Omega,$$

where $\Lambda_n^\varepsilon = \Omega_n^\varepsilon \setminus \hat{\Omega}_n^\varepsilon$ and $\hat{\Xi}_n^\varepsilon = \{\xi \in \Xi_n^\varepsilon : \varepsilon D_{x_n^\varepsilon}(Y + \xi) \subset (\Omega_n^\varepsilon \cap \Omega)\}$.

FIG. 2. Schematic representation of subdomains Ω_n^ε and $\hat{\Omega}_n^\varepsilon$.

As it is known from the periodic case, the unfolding operator provides a powerful technique for the multiscale analysis of problems posed in perforated domains and nonlinear equations defined on oscillating surfaces of microstructures. Thus, the main emphasis of this work will be on the development of the unfolding method for domains with locally-periodic perforations. Therefore, next we introduce perforated domains with locally-periodic changes in the distribution and in the shape of perforations.

We consider $Y_0 \subset Y$ with a Lipschitz boundary $\Gamma = \partial Y_0$ and a matrix K with $K, K^{-1} \in \text{Lip}(\mathbb{R}^d; \mathbb{R}^{d \times d})$, where $0 < K_1 \leq |\det K(x)| \leq K_2 < \infty$, $K_x Y_0 \subset Y$, and $Y^* = Y \setminus \bar{Y}_0$ and $\tilde{Y}_{K_x}^* = Y \setminus K_x \bar{Y}_0$ are connected, for all $x \in \bar{\Omega}$. Define $Y_{x,K}^* = D_x \tilde{Y}_{K_x}^*$ with the boundary $\Gamma_x = D_x K_x \Gamma$, where $K_x = K(x)$ and $D_x = D(x)$. Then, a domain with locally-periodic perforations is defined as

$$\Omega_{\varepsilon,K}^* = \text{Int} \left(\bigcup_{n=1}^{N_\varepsilon} \Omega_{n,K}^{*,\varepsilon} \right) \cap \Omega, \quad \text{where} \quad \Omega_{n,K}^{*,\varepsilon} = \bigcup_{\xi \in \Xi_n^{*,\varepsilon}} \varepsilon D_{x_n^\varepsilon} (\overline{\tilde{Y}_{K_{x_n^\varepsilon}}^* + \xi}) \cup \overline{\Lambda_n^{*,\varepsilon}}$$

Here $\Lambda_n^{*,\varepsilon} = \Omega_n^\varepsilon \setminus \bigcup_{\xi \in \Xi_n^{*,\varepsilon}} \varepsilon D_{x_n^\varepsilon} (\bar{Y} + \xi)$, with $\Xi_n^{*,\varepsilon} = \{\xi \in \Xi_n^\varepsilon : \varepsilon D_{x_n^\varepsilon} (Y + \xi) \subset \Omega_n^\varepsilon\}$, $\tilde{Y}_{K_{x_n^\varepsilon}}^* = Y \setminus K_{x_n^\varepsilon} \bar{Y}_0$ and $K_{x_n^\varepsilon} = K(x_n^\varepsilon)$ for $n = 1, \dots, N_\varepsilon$. The boundaries of the locally-periodic microstructure of $\Omega_{\varepsilon,K}^*$ are denoted by

$$\Gamma^\varepsilon = \bigcup_{n=1}^{N_\varepsilon} \Gamma_n^\varepsilon \cap \Omega, \quad \text{where} \quad \Gamma_n^\varepsilon = \bigcup_{\xi \in \Xi_n^{*,\varepsilon}} \varepsilon D_{x_n^\varepsilon} (\tilde{\Gamma}_{K_{x_n^\varepsilon}} + \xi) \quad \text{with} \quad x_n^\varepsilon \in \hat{\Omega}_n^\varepsilon,$$

and $\tilde{\Gamma}_{K_{x_n^\varepsilon}} = K_{x_n^\varepsilon} \Gamma$. Notice that changes in the microstructure of $\Omega_{\varepsilon,K}^*$ are defined by changes in the periodicity given by $D(x)$ and additional changes in the shape of perforations described by $K(x)$ for $x \in \Omega$.

Along with plywood-like structures (see Fig. 1), examples of locally-periodic microstructures are e.g. concrete materials with space-dependent perforations, plant and epithelial tissues, see Fig. 3. In the definition of microstructure of concrete materials with space-dependent perforations we have e.g. $D(x) = \mathbf{I}$ and $K(x) = \rho(x)\mathbf{I}$ for such $0 < \rho_1 \leq \rho(x) \leq \rho_2 < \infty$ that $K(x)\bar{Y}_0 \subset Y$, see e.g. [17, 45] and Fig. 2. For plant or epithelial tissues additionally we have space-dependent deformations of cells given by $D(x) \neq \mathbf{I}$, where \mathbf{I} denotes the identity matrix.

Using the mathematical definition of general locally-periodic microstructures, next we introduce the definition of the locally-periodic (l-p) unfolding operator, map-

FIG. 3. Examples of locally periodic microstructures with local changes in the shape and the periodicity of a microstructure. We observe changes in shape and size of cells in an epithelial tissue due to maturation, <http://www.eastcentral.edu>, and changes in the size of plant cells in a wood tissue, Schoch, Heller, Schweingruber, Kienast, 2004, [51].

ping functions defined on ε -dependent domains to functions depending on two variables (i.e. a microscopic variable and a macroscopic variable), but defined on fixed domains.

3. Definitions of l-p unfolding operator and l-p two-scale convergence on oscillating surfaces. The main idea of the two-scale convergence is to consider test functions which comprise the information about the microstructure and the microscopic properties of a composite material and of model equations. The same idea is used in the definition of l-t-s by considering a l-p approximation of $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$ (reflecting the locally-periodic properties of microscopic problems) as a test function.

DEFINITION 3.1. [49] Let $u^\varepsilon \in L^p(\Omega)$ for all $\varepsilon > 0$ and $p \in (1, \infty)$. We say the sequence $\{u^\varepsilon\}$ converges l-t-s to $u \in L^p(\Omega; L^p(Y_x))$ as $\varepsilon \rightarrow 0$ if $\|u^\varepsilon\|_{L^p(\Omega)} \leq C$ and for any $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega} u^\varepsilon(x) \mathcal{L}^\varepsilon \psi(x) dx = \int_{\Omega} \int_{Y_x} u(x, y) \psi(x, y) dy dx,$$

where \mathcal{L}^ε is the l-p approximation of ψ , defined in (2.1), and $1/p + 1/q = 1$.

Remark. Notice that the definition of l-t-s and convergence results presented in [49] for $p = 2$ are directly generalized to $p \in (1, \infty)$.

Motivated by the notion of the periodic unfolding operator and l-t-s convergence we define the l-p unfolding operator in the following way.

DEFINITION 3.2. For any Lebesgue-measurable on Ω function ψ the locally-periodic (l-p) unfolding operator $\mathcal{T}_{\mathcal{L}}^\varepsilon$ is defined as

$$\mathcal{T}_{\mathcal{L}}^\varepsilon(\psi)(x, y) = \sum_{n=1}^{N_\varepsilon} \psi(\varepsilon D_{x_n^\varepsilon} [D_{x_n^\varepsilon}^{-1} x / \varepsilon]_Y + \varepsilon D_{x_n^\varepsilon} y) \chi_{\hat{\Omega}_n^\varepsilon}(x) \quad \text{for } x \in \Omega \text{ and } y \in Y.$$

The definition implies that $\mathcal{T}_{\mathcal{L}}^\varepsilon(\psi)$ is Lebesgue-measurable on $\Omega \times Y$ and is zero for $x \in \Lambda^\varepsilon$.

For perforated domains with local changes in the distribution of perforations, but without additional changes in the shape of perforations, i.e. $K = \mathbf{I}$ and

$$\Omega_\varepsilon^* = \text{Int} \left(\bigcup_{n=1}^{N_\varepsilon} \Omega_n^{*,\varepsilon} \right) \cap \Omega, \quad \text{where} \quad \Omega_n^{*,\varepsilon} = \bigcup_{\xi \in \Xi_n^{*,\varepsilon}} \varepsilon D_{x_n^\varepsilon} (\bar{Y}^* + \xi) \cup \bar{\Lambda}_n^{*,\varepsilon},$$

and $Y^* = Y \setminus \bar{Y}_0$, we define the l-p unfolding operator in the following way:

DEFINITION 3.3. For any Lebesgue-measurable on Ω_ε^* function ψ the l-p unfolding operator $\mathcal{T}_\mathcal{L}^{*,\varepsilon}$ is defined as

$$\mathcal{T}_\mathcal{L}^{*,\varepsilon}(\psi)(x, y) = \sum_{n=1}^{N_\varepsilon} \psi(\varepsilon D_{x_n^\varepsilon} [D_{x_n^\varepsilon}^{-1} x / \varepsilon]_Y + \varepsilon D_{x_n^\varepsilon} y) \chi_{\hat{\Omega}_n^\varepsilon}(x) \quad \text{for } x \in \Omega \text{ and } y \in Y^*.$$

The definition implies that $\mathcal{T}_\mathcal{L}^{*,\varepsilon}(\psi)$ is Lebesgue-measurable on $\Omega \times Y^*$ and is zero for $x \in \Lambda^\varepsilon$.

In mathematical models posed in perforated domains we often have some processes defined on the surfaces of the microstructure (e.g. non-homogeneous Neumann conditions or equations defined on the boundaries of the microstructure). Therefore it is important to have a notion of a convergence for sequences defined on oscillating surfaces of locally-periodic microstructures. Applying the same idea as in the definition of l-t-s convergence for sequences in $L^p(\Omega)$ (i.e. considering l-p approximations of functions with space-dependent periodicity as test functions) we define the l-t-s on surfaces of locally-periodic microstructures.

DEFINITION 3.4. A sequence $\{u^\varepsilon\} \subset L^p(\Gamma^\varepsilon)$, with $p \in (1, \infty)$, is said to converge locally-periodic two-scale (l-t-s) to $u \in L^p(\Omega; L^p(\Gamma_x))$ if $\varepsilon \|u^\varepsilon\|_{L^p(\Gamma^\varepsilon)}^p \leq C$ and for any $\psi \in C(\bar{\Omega}; C_{per}(Y_x))$

$$\lim_{\varepsilon \rightarrow 0} \varepsilon \int_{\Gamma^\varepsilon} u^\varepsilon(x) \mathcal{L}^\varepsilon \psi(x) d\sigma_x = \int_\Omega \frac{1}{|Y_x|} \int_{\Gamma_x} u(x, y) \psi(x, y) d\sigma_y dx,$$

where \mathcal{L}^ε is the l-p approximation of ψ defined in (2.1).

Often, to show the strong convergence of a sequence defined on oscillating boundaries of a microstructure, we need to map it to a sequence defined on a fixed domain. This can be achieved by using the boundary unfolding operator.

DEFINITION 3.5. For any Lebesgue-measurable on Γ^ε function ψ the l-p boundary unfolding operator $\mathcal{T}_\mathcal{L}^{b,\varepsilon}$ is defined as

$$\mathcal{T}_\mathcal{L}^{b,\varepsilon}(\psi)(x, y) = \sum_{n=1}^{N_\varepsilon} \psi(\varepsilon D_{x_n^\varepsilon} [D_{x_n^\varepsilon}^{-1} x / \varepsilon]_Y + \varepsilon D_{x_n^\varepsilon} K_{x_n^\varepsilon} y) \chi_{\hat{\Omega}_n^\varepsilon}(x) \quad \text{for } x \in \Omega \text{ and } y \in \Gamma.$$

The definition implies that $\mathcal{T}_\mathcal{L}^{b,\varepsilon}(\psi)$ is Lebesgue-measurable on $\Omega \times \Gamma$ and is zero for $x \in \Lambda^\varepsilon$. The l-p boundary unfolding operator is a generalization of the periodic boundary unfolding operator, see e.g. [21, 22, 24, 46]. Similar to the periodic unfolding operator, the l-p unfolding operator maps functions defined in domains depending on ε (on Ω_ε^* or Γ^ε) to functions defined on fixed domains ($\Omega \times Y^*$ or $\Omega \times \Gamma$). The locally-periodic microstructures of domains are reflected in the definition of the l-p unfolding operator.

4. Main convergence results for the l-p unfolding operator and l-t-s convergence on oscillating surfaces. In this section we summarize the main results of the paper. Similar to the periodic case [21, 22], we obtain compactness results for l-t-s convergence on oscillating boundaries, for the l-p unfolding operator and for the l-p boundary unfolding operator. We prove convergence results for sequences bounded in $L^p(\Gamma^\varepsilon)$, $H^1(\Omega)$, and $H^1(\Omega_\varepsilon^*)$, respectively. The properties of the transformation matrices D and K , assumed in Section 3, are used to prove the convergence results stated in this section.

THEOREM 4.1. For a sequence $\{w^\varepsilon\} \subset L^p(\Omega)$, with $p \in (1, \infty)$, satisfying

$$\|w^\varepsilon\|_{L^p(\Omega)} + \varepsilon \|\nabla w^\varepsilon\|_{L^p(\Omega)} \leq C$$

there exist a subsequence (denoted again by $\{w^\varepsilon\}$) and $w \in L^p(\Omega; W_{per}^{1,p}(Y_x))$ such that

$$\begin{aligned} \mathcal{T}_{\mathcal{L}}^\varepsilon(w^\varepsilon) &\rightharpoonup w(\cdot, D_x \cdot) && \text{weakly in } L^p(\Omega; W^{1,p}(Y)), \\ \varepsilon \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla w^\varepsilon) &\rightharpoonup D_x^{-T} \nabla_y w(\cdot, D_x \cdot) && \text{weakly in } L^p(\Omega \times Y). \end{aligned}$$

For a uniformly bounded sequence in $W^{1,p}(\Omega)$, in addition we obtain the weak convergence of the unfolded sequence of derivatives, important for the homogenization of equations comprising elliptic operators of second order.

THEOREM 4.2. For a sequence $\{w^\varepsilon\} \subset W^{1,p}(\Omega)$, with $p \in (1, \infty)$, that converges weakly to w in $W^{1,p}(\Omega)$, there exist a subsequence (denoted again by $\{w^\varepsilon\}$) and a function $w_1 \in L^p(\Omega; W_{per}^{1,p}(Y_x))$ such that

$$\begin{aligned} \mathcal{T}_{\mathcal{L}}^\varepsilon(w^\varepsilon) &\rightharpoonup w && \text{weakly in } L^p(\Omega; W^{1,p}(Y)), \\ \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla w^\varepsilon)(\cdot, \cdot) &\rightharpoonup \nabla_x w(\cdot) + D_x^{-T} \nabla_y w_1(\cdot, D_x \cdot) && \text{weakly in } L^p(\Omega \times Y). \end{aligned}$$

Two of the main advantages of the unfolding operator are that it helps to overcome one of the difficulties of perforated domains which is the use of extension operators and it allows us to prove strong convergence for sequences defined on boundaries of microstructures. Thus next we formulate convergence results for the l-p unfolding operator in perforated domains and the l-p boundary unfolding operator.

THEOREM 4.3. For a sequence $\{w^\varepsilon\} \subset W^{1,p}(\Omega_\varepsilon^*)$, where $p \in (1, \infty)$, satisfying

$$(4.1) \quad \|w^\varepsilon\|_{L^p(\Omega_\varepsilon^*)} + \varepsilon \|\nabla w^\varepsilon\|_{L^p(\Omega_\varepsilon^*)} \leq C$$

there exist a subsequence (denoted again by $\{w^\varepsilon\}$) and $w \in L^p(\Omega; W_{per}^{1,p}(Y_x^*))$ such that

$$(4.2) \quad \begin{aligned} \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon) &\rightharpoonup w(\cdot, D_x \cdot) && \text{weakly in } L^p(\Omega; W^{1,p}(Y^*)), \\ \varepsilon \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\nabla w^\varepsilon) &\rightharpoonup D_x^{-T} \nabla_y w(\cdot, D_x \cdot) && \text{weakly in } L^p(\Omega \times Y^*). \end{aligned}$$

In the case w^ε is bounded in $W^p(\Omega_\varepsilon^*)$ uniformly with respect to ε , we obtain weak convergence of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\nabla w^\varepsilon)$ in $L^p(\Omega \times Y^*)$ and local strong convergence of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$.

THEOREM 4.4. For a sequence $\{w^\varepsilon\} \subset W^{1,p}(\Omega_\varepsilon^*)$, where $p \in (1, \infty)$, satisfying

$$\|w^\varepsilon\|_{W^{1,p}(\Omega_\varepsilon^*)} \leq C$$

there exist a subsequence (denoted again by $\{w^\varepsilon\}$) and functions $w \in W^{1,p}(\Omega)$ and $w_1 \in L^p(\Omega; W_{per}^{1,p}(Y_x^*))$ such that

$$\begin{aligned} \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon) &\rightharpoonup w && \text{weakly in } L^p(\Omega; W^{1,p}(Y^*)), \\ \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\nabla w^\varepsilon) &\rightharpoonup \nabla w + D_x^{-T} \nabla_y w_1(\cdot, D_x \cdot) && \text{weakly in } L^p(\Omega \times Y^*), \\ \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon) &\rightarrow w && \text{strongly in } L_{loc}^p(\Omega; W^{1,p}(Y^*)). \end{aligned}$$

Notice that the weak limit of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\nabla w^\varepsilon)$ reflects the locally-periodic microstructure of Ω_ε^* and depends on the transformation matrix D .

For l-t-s convergence on oscillating surfaces of microstructures we have following compactness result.

THEOREM 4.5. *For a sequence $\{w^\varepsilon\} \subset L^p(\Gamma^\varepsilon)$, with $p \in (1, \infty)$, satisfying*

$$\varepsilon \|w^\varepsilon\|_{L^p(\Gamma^\varepsilon)}^p \leq C$$

there exist a subsequence (denoted again by $\{w^\varepsilon\}$) and $w \in L^p(\Omega; L^p(\Gamma_x))$ such that

$$w^\varepsilon \rightarrow w \quad \text{locally periodic two-scale (l-t-s).}$$

Similar to the periodic case [21, 22], we show the relation between the l-t-s convergence on oscillating surfaces and the weak convergence of a sequence obtained by applying the l-p boundary unfolding operator.

THEOREM 4.6. *Let $\{w^\varepsilon\} \subset L^p(\Gamma^\varepsilon)$ with $\varepsilon \|w^\varepsilon\|_{L^p(\Gamma^\varepsilon)}^p \leq C$, where $p \in (1, \infty)$. The following assertions are equivalent*

- (i) $w^\varepsilon \rightarrow w$ *l-t-s, $w \in L^p(\Omega; L^p(\Gamma_x))$.*
- (ii) $\mathcal{T}_\mathcal{L}^{b,\varepsilon}(w^\varepsilon) \rightharpoonup w(\cdot, D_x K_x \cdot)$ *weakly in $L^p(\Omega \times \Gamma)$.*

Theorems 4.5 and 4.6 imply that for $\{w^\varepsilon\} \subset L^p(\Gamma^\varepsilon)$ with $\varepsilon \|w^\varepsilon\|_{L^p(\Gamma^\varepsilon)}^p \leq C$ we have the weak convergence of $\{\mathcal{T}_\mathcal{L}^{b,\varepsilon}(w^\varepsilon)\}$ in $L^p(\Omega \times \Gamma)$, where $p \in (1, \infty)$.

The definition of the l-p boundary unfolding operator and the relation between the l-t-s convergence of sequences defined on l-p oscillating boundaries and the l-p boundary unfolding operator allow us to obtain homogenization results for equations posed on the boundaries of locally-periodic microstructures.

5. The l-p unfolding operator: Proofs of convergence results. First we prove some properties of the l-p unfolding operator. Similar to the periodic case, we obtain that the l-p unfolding operator is linear and preserves strong convergence.

LEMMA 5.1. (i) *For $\phi \in L^p(\Omega)$, with $1 \leq p < \infty$, holds*

$$(5.1) \quad \int_{\Omega \times Y} |\mathcal{T}_\mathcal{L}^\varepsilon(\phi)(x, y)|^p dy dx \leq |Y| \int_{\Omega} |\phi(x)|^p dx.$$

(ii) $\mathcal{T}_\mathcal{L}^\varepsilon : L^p(\Omega) \rightarrow L^p(\Omega \times Y)$ *is a linear continuous operator, where $1 \leq p < \infty$.*

(iii) *For $\phi \in L^p(\Omega)$, with $1 \leq p < \infty$, we have strong convergence*

$$(5.2) \quad \mathcal{T}_\mathcal{L}^\varepsilon(\phi) \rightarrow \phi \quad \text{in} \quad L^p(\Omega \times Y).$$

(iv) *If $\phi^\varepsilon \rightarrow \phi$ in $L^p(\Omega)$, with $1 \leq p < \infty$, then $\mathcal{T}_\mathcal{L}^\varepsilon(\phi^\varepsilon) \rightarrow \phi$ in $L^p(\Omega \times Y)$.*

Proof. Using the definition of the l-p unfolding operator we obtain

$$(5.3) \quad \begin{aligned} \int_{\Omega \times Y} |\mathcal{T}_\mathcal{L}^\varepsilon(\phi)(x, y)|^p dy dx &= \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon^d |D_{x_n^\varepsilon} Y| \int_Y |\phi(D_{x_n^\varepsilon}(\varepsilon\xi + \varepsilon y))|^p dy \\ &= \sum_{n=1}^{N_\varepsilon} |Y| \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \int_{\varepsilon D_{x_n^\varepsilon}(\xi + Y)} |\phi(x)|^p dx = \sum_{n=1}^{N_\varepsilon} |Y| \int_{\hat{\Omega}_n^\varepsilon} |\phi(x)|^p dx. \end{aligned}$$

Then estimate (5.1) follows from the properties of the covering of Ω by $\{\hat{\Omega}_n^\varepsilon\}_{n=1}^{N_\varepsilon}$.

The result in (ii) is ensured by the definition of the l-p unfolding operator and inequality (5.1).

(iii) Using the fact that $\phi \in L^p(\Omega)$ and $|\Lambda^\varepsilon| \rightarrow 0$ as $\varepsilon \rightarrow 0$ (ensured by the properties of the covering of Ω by $\{\Omega_n^\varepsilon\}_{n=1}^{N_\varepsilon}$) and applying Lebesgue's Dominated Convergence Theorem, see e.g. [29], we obtain $\int_{\Lambda^\varepsilon} |\phi(x)|^p dx \rightarrow 0$ as $\varepsilon \rightarrow 0$.

Considering the approximation of L^p -functions by continuous functions, using the definition of $\mathcal{T}_\mathcal{L}^\varepsilon$ and equality (5.3), and taking the limit as $\varepsilon \rightarrow 0$ in the equality (5.3) imply the convergence stated in (iii).

(iv) The linearity of the l-p unfolding operator along with (5.1) and (5.2) yield

$$\|\mathcal{T}_\mathcal{L}^\varepsilon(\phi^\varepsilon) - \phi\|_{L^p(\Omega \times Y)} \leq |Y|^{\frac{1}{p}} \|\phi^\varepsilon - \phi\|_{L^p(\Omega)} + \|\mathcal{T}_\mathcal{L}^\varepsilon(\phi) - \phi\|_{L^p(\Omega \times Y)} \rightarrow 0 \text{ as } \varepsilon \rightarrow 0.$$

□

Similar to l-t-s convergence, the average of the weak limit of the unfolded sequence with respect to microscopic variables is equal to the weak limit of the original sequence.

LEMMA 5.2. *For $\{w^\varepsilon\}$ bounded in $L^p(\Omega)$, with $p \in (1, \infty)$, we have that $\{\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)\}$ is bounded in $L^p(\Omega \times Y)$ and if*

$$\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon) \rightharpoonup \tilde{w} \quad \text{weakly in } L^p(\Omega \times Y),$$

then

$$w^\varepsilon \rightharpoonup \int_Y \tilde{w} dy \quad \text{weakly in } L^p(\Omega).$$

Proof. The boundedness of $\{\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)\}$ in $L^p(\Omega \times Y)$ follows directly from the boundedness of $\{w^\varepsilon\}$ in $L^p(\Omega)$ and the estimate (5.1). For $\psi \in L^q(\Omega)$, $1/p + 1/q = 1$, using the definition of $\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)$ we have

$$\int_\Omega w^\varepsilon \psi dx = \frac{1}{|Y|} \int_{\Omega \times Y} \mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon) \mathcal{T}_\mathcal{L}^\varepsilon(\psi) dy dx + \mathcal{A}_\varepsilon, \quad \text{where } \mathcal{A}_\varepsilon = \int_{\Lambda^\varepsilon} w^\varepsilon \psi dx.$$

For $\{w^\varepsilon\}$ bounded in $L^p(\Omega)$ and $\psi \in L^q(\Omega)$, using the properties of the covering of Ω and the definition of $\hat{\Omega}_n^\varepsilon$ and Λ^ε , where $1 \leq n \leq N_\varepsilon$, we obtain $\mathcal{A}_\varepsilon \rightarrow 0$ as $\varepsilon \rightarrow 0$. Then, the weak convergence of $\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)$ and the strong convergence of $\mathcal{T}_\mathcal{L}^\varepsilon(\psi)$, shown in Lemma 5.1, imply

$$\lim_{\varepsilon \rightarrow 0} \int_\Omega w^\varepsilon(x) \psi(x) dx = \frac{1}{|Y|} \int_\Omega \int_Y \tilde{w}(x, y) \psi(x) dy dx$$

for any $\psi \in L^q(\Omega)$. □

For the periodic unfolding operator we have that $\mathcal{T}^\varepsilon(\psi(\cdot, \cdot/\varepsilon)) \rightarrow \psi$ in $L^q(\Omega \times Y)$ for $\psi \in L^q(\Omega, C_{\text{per}}(Y))$. A similar result holds for the l-p unfolding operator and $\psi \in L^q(\Omega, C_{\text{per}}(Y_x))$, but with $\psi(\cdot, \cdot/\varepsilon)$ replaced by the l-p approximation $\mathcal{L}^\varepsilon \psi(\cdot)$.

LEMMA 5.3. (i) *For $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$, with $q \in [1, \infty)$, we have*

$$\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}^\varepsilon \psi) \rightarrow \psi(\cdot, D_x \cdot) \quad \text{strongly in } L^q(\Omega \times Y).$$

(ii) *For $\psi \in C(\bar{\Omega}; L^q_{\text{per}}(Y_x))$, with $q \in [1, \infty)$, we have*

$$\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}_0^\varepsilon \psi) \rightarrow \psi(\cdot, D_x \cdot) \quad \text{strongly in } L^q(\Omega \times Y).$$

Proof. (i) For $\psi \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ using the definition of $\mathcal{L}^\varepsilon \psi$ and $\mathcal{T}_\mathcal{L}^\varepsilon$ we obtain

$$\int_{\Omega \times Y} |\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}^\varepsilon \psi)|^q dy dx = \sum_{n=1}^{N_\varepsilon} \int_{\hat{\Omega}_n^\varepsilon \times Y} \left| \tilde{\psi} \left(\varepsilon D_{x_n^\varepsilon} \left[\frac{D_{x_n^\varepsilon}^{-1} x}{\varepsilon} \right]_Y + \varepsilon D_{x_n^\varepsilon} y, y \right) \right|^q dy dx,$$

where $q \in [1, \infty)$ and $\tilde{\psi} \in C(\bar{\Omega}; C_{\text{per}}(Y))$ such that $\psi(x, y) = \tilde{\psi}(x, D_x^{-1} y)$ for $x \in \Omega$, $y \in Y_x$. Then, using the properties of the covering of Ω_n^ε by $\varepsilon Y_{x_n^\varepsilon}^\xi = \varepsilon D_{x_n^\varepsilon}(Y + \xi)$, with $\xi \in \Xi_n^\varepsilon$, and considering fixed points $y_\xi \in Y + \xi$ for $\xi \in \hat{\Xi}_n^\varepsilon$ we obtain

$$\int_{\Omega \times Y} |\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}^\varepsilon \psi)|^q dy dx = \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon^d |Y_{x_n^\varepsilon}| \int_Y |\tilde{\psi}(\varepsilon D_{x_n^\varepsilon}(\xi + y_\xi), y)|^q dy + \delta(\varepsilon),$$

where, due to the continuity of ψ and the properties of the covering of Ω by $\{\Omega_n^\varepsilon\}_{n=1}^{N_\varepsilon}$,

$$\delta(\varepsilon) = \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon^d |Y_{x_n^\varepsilon}| \int_Y \left(|\tilde{\psi}(\varepsilon D_{x_n^\varepsilon}(\xi + y_\xi), y)|^q - |\tilde{\psi}(\varepsilon D_{x_n^\varepsilon}(\xi + y), y)|^q \right) dy \rightarrow 0$$

as $\varepsilon \rightarrow 0$. Then, using the continuity of ψ and D together with the relation between ψ and $\tilde{\psi}$ we obtain

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega \times Y} |\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}^\varepsilon \psi)|^q dy dx = \int_{\Omega \times Y} |\tilde{\psi}(x, y)|^q dy dx = \int_{\Omega \times Y} |\psi(x, D_x y)|^q dy dx.$$

The continuity of ψ with respect to x ensures the point-wise convergence of $\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}^\varepsilon \psi)(x, y)$ to $\psi(x, D_x y)$ a.e. in $\Omega \times Y$.

Considering an approximation of $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$ by $\psi_m \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ and the convergence

$$\lim_{m \rightarrow \infty} \lim_{\varepsilon \rightarrow 0} \int_{\Omega} (|\mathcal{L}^\varepsilon \psi_m(x)|^q - |\mathcal{L}^\varepsilon \psi(x)|^q) dx = 0,$$

see [49, Lemma 3.4] for the proof, yields $\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}^\varepsilon \psi)(\cdot, \cdot) \rightarrow \psi(\cdot, D_x \cdot)$ in $L^q(\Omega \times Y)$ for $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$.

(ii) For $\psi \in C(\bar{\Omega}; L_{\text{per}}^q(Y_x))$, we can prove the strong convergence only of $\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}_0^\varepsilon \psi)$. Consider

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega \times Y} |\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}_0^\varepsilon \psi)(x, y)|^q dy dx = |Y| \lim_{\varepsilon \rightarrow 0} \left[\int_{\Omega} |\mathcal{L}_0^\varepsilon \psi(x)|^q dx - \int_{\Lambda^\varepsilon} |\mathcal{L}_0^\varepsilon \psi(x)|^q dx \right].$$

Then, using Lemma 3.4 in [49] along with the regularity of ψ and the properties of Λ^ε we obtain

$$|Y| \lim_{\varepsilon \rightarrow 0} \int_{\Omega} |\mathcal{L}_0^\varepsilon \psi(x)|^q dx = \int_{\Omega \times Y} |\psi(x, D_x y)|^q dy dx, \quad \lim_{\varepsilon \rightarrow 0} \int_{\Lambda^\varepsilon} |\mathcal{L}_0^\varepsilon \psi(x)|^q dx = 0.$$

The continuity of ψ with respect to $x \in \Omega$ implies $\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}_0^\varepsilon \psi)(x, y) \rightarrow \psi(x, D_x y)$ point-wise a.e. in $\Omega \times Y$. \square

Remark. Notice that for $\psi \in C(\bar{\Omega}; L_{\text{per}}^q(Y_x))$ we have the strong convergence only of $\mathcal{T}_\mathcal{L}^\varepsilon(\mathcal{L}_0^\varepsilon \psi)$. However, this convergence result is sufficient for the derivation of

homogenization results, since the microscopic properties of the considered processes or domains can be represented by coefficients in the form $B\mathcal{L}_0^\varepsilon A$, with some given functions $B \in L^\infty(\Omega)$ and $A \in C(\bar{\Omega}; L^q_{\text{per}}(Y_x))$.

The strong convergence of $\mathcal{T}_\varepsilon^\varepsilon(\mathcal{L}^\varepsilon\psi)$ for $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$ is now used to show the equivalence between the weak convergence of the l-p unfolded sequence and l-t-s convergence of the original sequence. Notice that $L^q(\Omega; C_{\text{per}}(Y_x))$ represents the set of test functions admissible in the definition of the l-t-s convergence.

LEMMA 5.4. *Let $\{w^\varepsilon\}$ be a bounded sequence in $L^p(\Omega)$, where $p \in (1, \infty)$. Then the following assertions are equivalent*

- (i) $w^\varepsilon \rightarrow w$ l-t-s, $w \in L^p(\Omega; L^p(Y_x))$,
- (ii) $\mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon)(\cdot, \cdot) \rightharpoonup w(\cdot, D_x \cdot)$ weakly in $L^p(\Omega \times Y)$.

Proof. [(ii) \Rightarrow (i)] Since $\{w^\varepsilon\}$ is bounded in $L^p(\Omega)$, there exists (up to a subsequence) a l-t-s limit of w^ε as $\varepsilon \rightarrow 0$. For an arbitrary $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$ the weak convergence of $\mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon)$, and the strong convergence of $\mathcal{T}_\varepsilon^\varepsilon(\mathcal{L}^\varepsilon(\psi))$ ensure

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_{\Omega} w^\varepsilon \mathcal{L}^\varepsilon(\psi) dx &= \lim_{\varepsilon \rightarrow 0} \left[\int_{\Omega} \int_Y \mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon) \mathcal{T}_\varepsilon^\varepsilon(\mathcal{L}^\varepsilon(\psi)) dy dx + \int_{\Lambda^\varepsilon} w^\varepsilon \mathcal{L}^\varepsilon(\psi) dx \right] \\ &= \int_{\Omega} \int_Y \hat{w}(x, y) \psi(x, D_x y) dy dx = \int_{\Omega} \int_{Y_x} w \psi dy dx, \end{aligned}$$

where $\hat{w}(x, y) = w(x, D(x)y)$ for a.a. $x \in \Omega$, $y \in Y$. Thus the whole sequence w^ε converges l-t-s to w .

[(i) \Rightarrow (ii)] On the other hand, the boundedness of $\{w^\varepsilon\}$ in $L^p(\Omega)$ implies the boundedness of $\{\mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon)\}$ and (up to a subsequence) the weak convergence of $\mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon)$ in $L^p(\Omega \times Y)$. If $w^\varepsilon \rightarrow w$ l-t-s, then

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_{\Omega} \int_Y \mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon) \mathcal{T}_\varepsilon^\varepsilon(\mathcal{L}^\varepsilon(\psi)) dy dx &= \lim_{\varepsilon \rightarrow 0} \left[\int_{\Omega} w^\varepsilon \mathcal{L}^\varepsilon(\psi) dx - \int_{\Lambda^\varepsilon} w^\varepsilon \mathcal{L}^\varepsilon(\psi) dx \right] \\ &= \int_{\Omega} \int_{Y_x} w \psi dy dx \end{aligned}$$

for $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$. Since $\mathcal{T}_\varepsilon^\varepsilon(\mathcal{L}^\varepsilon(\psi))(\cdot, \cdot) \rightarrow \psi(\cdot, D_x \cdot)$ in $L^q(\Omega \times Y)$, we obtain the weak convergence of the whole sequence $\mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon)$ to $w(\cdot, D_x \cdot)$ in $L^p(\Omega \times Y)$. Notice that the boundedness of $\{w^\varepsilon\}$ in $L^p(\Omega)$ and the fact that $|\Lambda^\varepsilon| \rightarrow 0$ as $\varepsilon \rightarrow 0$ imply

$$\int_{\Lambda^\varepsilon} |w^\varepsilon \mathcal{L}^\varepsilon(\psi)| dx \leq C \left(\int_{\Lambda^\varepsilon} \sup_{y \in Y} |\psi(x, D_x y)|^q dx \right)^{1/q} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0$$

for $\psi \in L^q(\Omega; C_{\text{per}}(Y_x))$ and $1/p + 1/q = 1$. \square

Next, we prove the main convergence results for the l-p unfolding operator, i.e. convergence results for $\{\mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon)\}$, $\{\varepsilon \mathcal{T}_\varepsilon^\varepsilon(\nabla w^\varepsilon)\}$ and $\{\mathcal{T}_\varepsilon^\varepsilon(\nabla w^\varepsilon)\}$.

The definition of the l-p unfolding operator yields that for $w \in W^{1,p}(\Omega)$

$$(5.4) \quad \nabla_y \mathcal{T}_\varepsilon^\varepsilon(w) = \varepsilon \sum_{n=1}^{N_\varepsilon} D_{x_n^T}^T \mathcal{T}_\varepsilon^\varepsilon(\nabla w) \chi_{\Omega_n^\varepsilon}.$$

Due to the regularity of D , the uniform boundedness of $\varepsilon \nabla w^\varepsilon$ implies the uniform boundedness of $\nabla_y \mathcal{T}_\varepsilon^\varepsilon(w^\varepsilon)$. Thus, assuming the boundedness of $\{\varepsilon \nabla w^\varepsilon\}$ we obtain convergence of the derivatives with respect to the microscopic variables, but have no information about the macroscopic derivatives.

Proof. [**Proof of Theorem 4.1**] The assumptions on $\{w^\varepsilon\}$ together with inequality (5.1), equality (5.4), and regularity of D ensure that $\{\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)\}$ is bounded in $L^p(\Omega; W^{1,p}(Y))$. Thus, there exists a subsequence, denoted again by $\{w^\varepsilon\}$, and a function $\tilde{w} \in L^p(\Omega; W^{1,p}(Y))$, such that $\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon) \rightharpoonup \tilde{w}$ in $L^p(\Omega; W^{1,p}(Y))$. We define $w(x, y) = \tilde{w}(x, D_x^{-1}y)$ for a.a. $x \in \Omega$, $y \in Y_x$. Due to the regularity of D , we have $w \in L^p(\Omega; W^{1,p}(Y_x))$. For $\phi \in C_0^\infty(\Omega \times Y)$, using the convergence of $\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)$, we have

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_{\Omega \times Y} \varepsilon \mathcal{T}_\mathcal{L}^\varepsilon(\nabla w^\varepsilon) \phi \, dy \, dx &= - \lim_{\varepsilon \rightarrow 0} \int_{\Omega \times Y} \mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon) \sum_{n=1}^{N_\varepsilon} \operatorname{div}_y(D_{x_n}^{-1} \phi(x, y)) \, dy \, \chi_{\Omega_n^\varepsilon} \, dx \\ &= - \int_{\Omega \times Y} w(x, D_x y) \operatorname{div}_y(D_x^{-1} \phi(x, y)) \, dy \, dx = \int_{\Omega \times Y} D_x^{-T} \nabla_y w(x, D_x y) \phi(x, y) \, dy \, dx. \end{aligned}$$

Hence, $\varepsilon \mathcal{T}_\mathcal{L}^\varepsilon(\nabla w^\varepsilon)(\cdot, \cdot) \rightharpoonup D_x^{-T} \nabla_y w(\cdot, D_x \cdot)$ in $L^p(\Omega \times Y)$ as $\varepsilon \rightarrow 0$. To show the Y_x -periodicity of w , i.e. Y -periodicity of \tilde{w} , we show first the periodicity in e_d -direction. Then considering similar calculations in each e_j -direction, with $j = 1, \dots, d-1$ and $\{e_j\}_{j=1, \dots, d}$ being the canonical basis of \mathbb{R}^d , we obtain the Y_x -periodicity of w . For $\psi \in C_0^\infty(\Omega; C^\infty(Y'))$ we consider

$$I = \int_{\Omega \times Y'} [\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)(x, (y', 1)) - \mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)(x, (y', 0))] \psi(x, y') \, dy' \, dx,$$

where $Y' = (0, 1)^{d-1}$. We define

$$\tilde{\Omega}_n^{\varepsilon, d} = \operatorname{Int} \left(\bigcup_{\xi \in \tilde{\Xi}_n^{\varepsilon, d}} \varepsilon D_{x_n}(\bar{Y} + \xi) \right), \quad \tilde{\Lambda}_{n, j}^\varepsilon = \operatorname{Int} \left(\bigcup_{\xi \in \tilde{\Xi}_{n, j}^{\varepsilon, d}} \varepsilon D_{x_n}(\bar{Y} + \xi) \right) \quad \text{for } j = 1, 2,$$

where $\tilde{\Xi}_n^{\varepsilon, d} = \{\xi \in \hat{\Xi}_n^\varepsilon : \varepsilon D_{x_n}(Y + e_d + \xi) \subset \hat{\Omega}_n^\varepsilon \text{ and } \varepsilon D_{x_n}(Y - e_d + \xi) \subset \hat{\Omega}_n^\varepsilon\}$, and $\tilde{\Xi}_n^{\varepsilon, d} = \hat{\Xi}_n^\varepsilon \setminus \tilde{\Xi}_n^{\varepsilon, d}$, with $\tilde{\Xi}_n^{\varepsilon, d} = \tilde{\Xi}_{n, 1}^{\varepsilon, d} \cup \tilde{\Xi}_{n, 2}^{\varepsilon, d}$, where $\tilde{\Xi}_{n, 1}^{\varepsilon, d}$ corresponds to upper cells in the e_d direction and $\tilde{\Xi}_{n, 2}^{\varepsilon, d}$ corresponds to lower cells in the e_d direction in $\hat{\Omega}_n^\varepsilon \setminus \tilde{\Omega}_n^{\varepsilon, d}$. Then using the definition of $\mathcal{T}_\mathcal{L}^\varepsilon$ we can write

$$\begin{aligned} I &= \sum_{n=1}^{N_\varepsilon} \int_{\tilde{\Omega}_n^{\varepsilon, d} \times Y'} \mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)(x, y^0) [\psi(x - \varepsilon D_{x_n} e_d, y') - \psi(x, y')] \, dy' \, dx \\ &+ \sum_{n=1}^{N_\varepsilon} \left[\int_{\tilde{\Lambda}_{n, 1}^\varepsilon \times Y'} \mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)(x, y^1) \psi(x, y') \, dy' \, dx - \int_{\tilde{\Lambda}_{n, 2}^\varepsilon \times Y'} \mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)(x, y^0) \psi(x, y') \, dy' \, dx \right], \end{aligned}$$

where $y^1 = (y', 1)$ and $y^0 = (y', 0)$. Using the continuity of ψ , the boundedness of the trace of $\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)$ in $L^p(\Omega \times Y')$, ensured by the assumptions on w^ε , and the fact that $\sum_{n=1}^{N_\varepsilon} |\tilde{\Lambda}_{n, j}^\varepsilon| \leq C \varepsilon^{1-r} \rightarrow 0$ as $\varepsilon \rightarrow 0$, with $0 < r < 1$ and $j = 1, 2$, we obtain that $I \rightarrow 0$ as $\varepsilon \rightarrow 0$. Similar calculations for e_j , with $j = 1, \dots, d-1$, and the convergence of the trace of $\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)$ in $L^p(\Omega \times Y')$, ensured by the weak convergence of $\mathcal{T}_\mathcal{L}^\varepsilon(w^\varepsilon)$ in $L^p(\Omega; W^{1,p}(Y))$, imply the Y_x -periodicity of w . \square

If $\|\nabla w^\varepsilon\|_{L^p(\Omega)}$ is bounded uniformly in ε , we have the weak convergence of w^ε in $W^{1,p}(\Omega)$ and of $\mathcal{T}_\mathcal{L}^\varepsilon(\nabla w^\varepsilon)$ in $L^p(\Omega \times Y)$. Hence we have information about the macroscopic and microscopic gradients of limit functions. The proof of the convergence results for $\mathcal{T}_\mathcal{L}^\varepsilon(\nabla w^\varepsilon)$ makes use of the Poincaré inequality for an auxiliary sequence. For this purpose we define a local average operator $\mathcal{M}_\mathcal{L}^\varepsilon$, i.e. an average of the unfolded function with respect to the microscopic variables.

DEFINITION 5.5. *The local average operator $\mathcal{M}_{\mathcal{L}}^\varepsilon : L^p(\Omega) \rightarrow L^p(\Omega)$, $p \in [1, \infty]$, is defined as*

$$(5.5) \quad \mathcal{M}_{\mathcal{L}}^\varepsilon(\psi)(x) = \int_Y \mathcal{T}_{\mathcal{L}}^\varepsilon(\psi)(x, y) dy = \sum_{n=1}^{N_\varepsilon} \int_Y \psi(\varepsilon D_{x_n^\varepsilon}([D_{x_n^\varepsilon}^{-1}x/\varepsilon] + y)) dy \chi_{\hat{\Omega}_n^\varepsilon}(x).$$

Proof. [Proof of Theorem 4.2] The proof of the convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla w^\varepsilon)$ follows similar ideas as in the case of the periodic unfolding operator. However, the proof of the periodicity of the corrector w_1 involves new ideas and technical details.

The convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(w^\varepsilon)$ follows from Lemma 5.2 and the fact that due to the assumption on $\{w^\varepsilon\}$ and regularity of D we have

$$\|\nabla_y \mathcal{T}_{\mathcal{L}}^\varepsilon(w^\varepsilon)\|_{L^p(\Omega \times Y)} \leq C\varepsilon \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0.$$

To show the convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla w^\varepsilon)$ we consider a function $V^\varepsilon : \Omega \times Y \rightarrow \mathbb{R}$ defined as

$$(5.6) \quad V^\varepsilon = \varepsilon^{-1} (\mathcal{T}_{\mathcal{L}}^\varepsilon(w^\varepsilon) - \mathcal{M}_{\mathcal{L}}^\varepsilon(w^\varepsilon)).$$

Then, the definition of $\mathcal{T}_{\mathcal{L}}^\varepsilon$ and $\mathcal{M}_{\mathcal{L}}^\varepsilon$ implies

$$\nabla_y V^\varepsilon = \frac{1}{\varepsilon} \nabla_y \mathcal{T}_{\mathcal{L}}^\varepsilon(w^\varepsilon) = \sum_{n=1}^{N_\varepsilon} D_{x_n^\varepsilon}^T \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla w^\varepsilon) \chi_{\Omega_n^\varepsilon}.$$

The boundedness of $\{w^\varepsilon\}$ in $W^{1,p}(\Omega)$ together with (5.1) and regularity assumptions on D imply that the sequence $\{\nabla_y V^\varepsilon\}$ is bounded in $L^p(\Omega \times Y)$. Considering

$$\int_Y V^\varepsilon dy = 0 \quad \text{and} \quad \int_Y y_c^\varepsilon \cdot \nabla w dy = 0 \quad \text{with} \quad y_c^\varepsilon = \sum_{n=1}^{N_\varepsilon} D_{x_n^\varepsilon} y_c \chi_{\Omega_n^\varepsilon},$$

where $y_c = (y_1 - \frac{1}{2}, \dots, y_d - \frac{1}{2})$ for $y \in Y$, and applying the Poincaré inequality to $V^\varepsilon - y_c^\varepsilon \cdot \nabla w$ yields

$$\|V^\varepsilon - y_c^\varepsilon \cdot \nabla w\|_{L^p(\Omega \times Y)} \leq C_1 \|\nabla_y V^\varepsilon - D_{x_n^\varepsilon}^T \nabla w\|_{L^p(\Omega \times Y)} \leq C_2.$$

Thus, there exists a subsequence (denoted again by $\{V^\varepsilon - y_c^\varepsilon \cdot \nabla w\}$) and $\tilde{w}_1 \in L^p(\Omega; W^{1,p}(Y))$ such that

$$(5.7) \quad V^\varepsilon - y_c^\varepsilon \cdot \nabla w \rightharpoonup \tilde{w}_1 \quad \text{weakly in} \quad L^p(\Omega; W^{1,p}(Y)).$$

For $\phi \in W^{1,p}(\Omega)$ we have the following relation

$$\mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla \phi)(x, y) = \varepsilon^{-1} \sum_{n=1}^{N_\varepsilon} D_{x_n^\varepsilon}^{-T} \nabla_y \mathcal{T}_{\mathcal{L}}^\varepsilon(\phi)(x, y) \chi_{\Omega_n^\varepsilon}(x).$$

Then the convergence in (5.7) and the continuity of D yield

$$(5.8) \quad \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla w^\varepsilon) = \sum_{n=1}^{N_\varepsilon} D_{x_n^\varepsilon}^{-T} \nabla_y V^\varepsilon \chi_{\Omega_n^\varepsilon} \rightharpoonup \nabla w + D_x^{-T} \nabla_y \tilde{w}_1 \quad \text{weakly in} \quad L^p(\Omega \times Y).$$

We show now that $\tilde{w}_1(x, y)$ is Y -periodic. Then the function $w_1(x, y) = \tilde{w}_1(x, D_x^{-1}y)$ for a.a. $x \in \Omega$, $y \in Y_x$ will be Y_x -periodic. For $\psi \in C_0^\infty(\Omega; C_0^\infty(Y'))$ we consider

$$\int_{\Omega} \int_{Y'} [V^\varepsilon(x, y^1) - V^\varepsilon(x, y^0)] \psi(x, y') dy' dx = \sum_{n=1}^{N_\varepsilon} (\mathcal{I}_{1,n} + \mathcal{I}_{2,n})$$

with

$$\begin{aligned} \mathcal{I}_{1,n} &= \int_{\tilde{\Omega}_n^{\varepsilon,d}} \int_{Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^0) \frac{1}{\varepsilon} [\psi(x - \varepsilon D_{x_n^\varepsilon} e_d, y') - \psi(x, y')] dy' dx, \\ \mathcal{I}_{2,n} &= \frac{1}{\varepsilon} \left[\int_{\tilde{\Lambda}_{n,1}^\varepsilon \times Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^1) \psi(x, y') dy' dx - \int_{\tilde{\Lambda}_{n,2}^\varepsilon \times Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^0) \psi(x, y') dy' dx \right] \\ &= \mathcal{I}_{2,n}^u - \mathcal{I}_{2,n}^l, \end{aligned}$$

where y^1 , y^0 , $\tilde{\Omega}_n^{\varepsilon,d}$, and $\tilde{\Lambda}_{n,j}^\varepsilon$, with $j = 1, 2$, are defined in the proof of Theorem 4.1. Then Lemma 5.1 and the strong convergence of $\{w^\varepsilon\}$ in $L^p(\Omega)$, ensured by the boundedness of $\{w^\varepsilon\}$ in $W^{1,p}(\Omega)$, imply the strong convergence of $\{\mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)\}$ to w in $L^p(\Omega \times Y)$. The boundedness of $\{\nabla_y \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)\}$ (ensured by the boundedness of $\{\nabla w^\varepsilon\}$) yields the weak convergence of $\{\mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)\}$ in $L^p(\Omega; W^{1,p}(Y))$ to the same w . Applying the trace theorem in $W^{1,p}(Y)$ we obtain that the trace of $\mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)$ on $\Omega \times Y'$ converges weakly to w in $L^p(\Omega \times Y')$ as $\varepsilon \rightarrow 0$. This together with the regularity of ψ and D gives

$$\lim_{\varepsilon \rightarrow 0} \sum_{n=1}^{N_\varepsilon} \mathcal{I}_{1,n} = - \int_{\Omega} \int_{Y'} w(x) D_d(x) \cdot \nabla_x \psi(x, y') dy' dx,$$

where $D_d(x) = (D_{1d}(x), \dots, D_{dd}(x))^T$. As next we consider the integrals over the upper (in e_d direction) cells \mathcal{I}_{2,n_1}^u and over the lower cells \mathcal{I}_{2,n_2}^l in neighboring $\Omega_{n_1}^\varepsilon$ and $\Omega_{n_2}^\varepsilon$, i.e. for such $1 \leq n_1, n_2 \leq N_\varepsilon$ that $\Theta_{n_1,2} = (\partial\Omega_{n_1}^\varepsilon \cap \partial\Omega_{n_2}^\varepsilon) \cap \{x_d = \text{const}\} \neq \emptyset$, $\dim(\Theta_{n_1,2}) = d - 1$, and $x_{n_1,d}^\varepsilon < x_{n_2,d}^\varepsilon$, and write

$$\begin{aligned} \mathcal{I}_{2,n_1}^u + \mathcal{I}_{2,n_2}^l &= \frac{1}{\varepsilon} \left[\int_{\tilde{\Lambda}_{n_1,1}^\varepsilon \times Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^0) \psi dy' dx - \int_{\tilde{\Lambda}_{n_2,2}^\varepsilon \times Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^0) \psi dy' dx \right] \\ &\quad + \int_{\tilde{\Lambda}_{n_1,1}^\varepsilon} \frac{1}{\varepsilon} \left[\int_{Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^1) \psi dy' - \int_{Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^0) \psi dy' \right] dx = \mathcal{I}_{2,n}^{1,2} + \mathcal{I}_{2,n}^1. \end{aligned}$$

The second integral $\mathcal{I}_{2,n}^1$ can be rewritten as

$$\mathcal{I}_{2,n}^1 = \frac{1}{\varepsilon} \int_{\tilde{\Lambda}_{n_1,1}^\varepsilon \times Y} \partial_{y_d} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y) \psi(x, y') dy dx = \int_{\tilde{\Lambda}_{n_1,1}^\varepsilon \times Y} D_d(x) \cdot \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(\nabla w^\varepsilon) \psi dy dx.$$

Using the boundedness of $\{\nabla w^\varepsilon\}$ in $L^p(\Omega)$, $p \in (1, \infty)$, and $\sum_{n_1=1}^{N_\varepsilon} |\tilde{\Lambda}_{n_1,1}^\varepsilon| \leq C\varepsilon^{1-r}$, we conclude that $\sum_{n=1}^{N_\varepsilon} \mathcal{I}_{2,n}^1 \rightarrow 0$ as $\varepsilon \rightarrow 0$ and $r < 1$. In $\mathcal{I}_{2,n}^{1,2}$ we shall distinguish between variations in x_1, \dots, x_{d-1} -directions and in x_d -direction. For an arbitrary fixed $x_{n_1,2}^\varepsilon \in \Theta_{n_1,2}$ we define $\hat{D}_{x_{n_1,2}^\varepsilon}^j = (D_1(x_{n_1,2}^\varepsilon), \dots, D_{d-1}(x_{n_1,2}^\varepsilon), D_d(x_{n_1,2}^\varepsilon))$, where $D_i(x) = (D_{1i}(x), \dots, D_{di}(x))^T$ with $i = 1, \dots, d$, and $j = 1, 2$. We introduce

$$\hat{\Lambda}_{n_j}^\varepsilon = \bigcup_{\xi \in \tilde{\Xi}_{n_1,2}^{\varepsilon,j}} \varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^j (Y + \xi) \quad \text{for } j = 1, 2,$$

where

$$\begin{aligned}\tilde{\Xi}_{n_1,2}^{\varepsilon,1} &= \left\{ \xi \in \mathbb{Z}^d : \varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^1 (\bar{Y} + \xi + e_d) \cap \Theta_{n_1,2} \neq \emptyset \text{ and } \varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^1 (Y + \xi) \subset \Omega_{n_1}^\varepsilon \right\}, \\ \tilde{\Xi}_{n_1,2}^{\varepsilon,2} &= \left\{ \xi \in \mathbb{Z}^d : \varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^2 (\bar{Y} + \xi - e_d) \cap \Theta_{n_1,2} \neq \emptyset \text{ and } \varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^2 (Y + \xi) \subset \Omega_{n_2}^\varepsilon \right\}.\end{aligned}$$

Then each of the integrals in $\mathcal{I}_{2,n}^{1,2}$ we rewrite as

$$\begin{aligned}\frac{1}{\varepsilon} \int_{\tilde{\Lambda}_{n_j,j}^\varepsilon} \int_{Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^0) \psi dy' dx &= \frac{1}{\varepsilon} \int_{\hat{\Lambda}_{n_j}^\varepsilon} \int_{Y'} w^\varepsilon(\varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^j ([x_{D,n}^j/\varepsilon] + y^0)) \psi dy' dx \\ + \frac{1}{\varepsilon} \left[\int_{\tilde{\Lambda}_{n_j,j}^\varepsilon} \int_{Y'} \mathcal{T}_{\tilde{\mathcal{L}}}^\varepsilon(w^\varepsilon)(x, y^0) \psi dy' dx - \int_{\hat{\Lambda}_{n_j}^\varepsilon} \int_{Y'} w^\varepsilon(\varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^j ([x_{D,n}^j/\varepsilon] + y^0)) \psi dy' dx \right] \\ &= J_{j,n}^1 + J_{j,n}^2,\end{aligned}$$

where $x_{D,n}^j = (\hat{D}_{x_{n_1,2}^\varepsilon}^j)^{-1}x$ and $j = 1, 2$. Using the definition of $\hat{\Lambda}_{n_j}^\varepsilon$, for $j = 1, 2$, and the fact that $|\tilde{\Xi}_{n_1,2}^{\varepsilon,1}| = |\tilde{\Xi}_{n_1,2}^{\varepsilon,2}| = I_{n_1,2}^\varepsilon$ yields

$$\begin{aligned}J_{1,n}^1 - J_{2,n}^1 &= \varepsilon^n \sum_{i=1}^{I_{n_1,2}^\varepsilon} \int_Y \int_{Y'} \frac{1}{\varepsilon} \left[w^\varepsilon(\varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^1 (\xi_i^1 + y^0)) \psi(\varepsilon \tilde{y}_{n_1,\xi}^i, y') \right. \\ &\quad \left. - w^\varepsilon(\varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^2 (\xi_i^2 + y^0)) \psi(\varepsilon \tilde{y}_{n_2,\xi}^i, y') \right] |\hat{D}_{x_{n_1,2}^\varepsilon}^1| dy' d\tilde{y} \\ &- \varepsilon^{r-1} \sum_{\xi \in \tilde{\Xi}_{n_1,2}^{\varepsilon,2}} \int_{\varepsilon(Y+\xi)} \int_{Y'} w^\varepsilon(\varepsilon \hat{D}_{x_{n_1,2}^\varepsilon}^2 (\xi + y^0)) \psi dy' \frac{1}{\varepsilon^r} \left[d(\hat{D}_{x_{n_1,2}^\varepsilon}^2 \tilde{x}) - d(\hat{D}_{x_{n_1,2}^\varepsilon}^1 \tilde{x}) \right],\end{aligned}$$

where $\tilde{y}_{n_j,\xi}^i = \hat{D}_{x_{n_1,2}^\varepsilon}^j(\tilde{y} + \xi_i^j)$ for $j = 1, 2$. The first integral in the last equality can be estimated by

$$C\varepsilon^{rd+(1-r)} \|w^\varepsilon\|_{W^{1,p}(\Omega)} \|\psi\|_{C_0^1(\Omega; C_0^1(Y'))}.$$

In the second integral we have a discrete derivative of an integral over an evolving domain, which converges to the divergence of the velocity vector D_d as $\varepsilon \rightarrow 0$. Then, using the fact that $|N_\varepsilon| \leq C\varepsilon^{-dr}$ and $x_{n_1,d}^\varepsilon < x_{n_2,d}^\varepsilon$ together with the regularity of D and the definition of $\hat{D}_{x_{n_1,2}^\varepsilon}^j$, where $j = 1, 2$, yields

$$\sum_{n=1}^{N_\varepsilon} (J_{1,n}^1 - J_{2,n}^1) \rightarrow - \int_\Omega \int_{Y'} w(x) \psi(x, y') \operatorname{div} D_d(x) dy' dx \quad \text{as } \varepsilon \rightarrow 0.$$

For $J_{1,n}^2 - J_{2,n}^2$ using the definition of $\tilde{\Lambda}_{n_j,j}^\varepsilon$ and $\hat{\Lambda}_{n_j}^\varepsilon$, the regularity of D and ψ , the boundedness of $\{w^\varepsilon\}$ in $W^{1,p}(\Omega)$, along with the properties of the covering of Ω by $\{\Omega_n^\varepsilon\}_{n=1}^{N_\varepsilon}$ we obtain

$$\sum_{n=1}^{N_\varepsilon} |J_{1,n}^2 - J_{2,n}^2| \leq C\varepsilon^{1-r} \sum_{k=1}^{d-1} \|\operatorname{div} D_k\|_{L^\infty(\Omega)} \|w^\varepsilon\|_{W^{1,p}(\Omega)} \|\psi\|_{C_0^\infty(\Omega; C_0(Y'))} \rightarrow 0$$

as $\varepsilon \rightarrow 0$ for $r \in (0, 1)$. Combining the obtained results we conclude that

$$\sum_{n=1}^{N_\varepsilon} (\mathcal{I}_{1,n} + \mathcal{I}_{2,n}) \rightarrow - \int_{\Omega \times Y'} [w(x) D_d(x) \cdot \nabla_x \psi(x, y') + w(x) \psi(x, y') \operatorname{div} D_d(x)] dy' dx$$

as $\varepsilon \rightarrow 0$. The definition of $y_c^\varepsilon \cdot \nabla w$ implies

$$(y_c^\varepsilon \cdot \nabla w(x))(y', 1) - (y_c^\varepsilon \cdot \nabla w(x))(y', 0) = \sum_{n=1}^{N_\varepsilon} D_d(x_n^\varepsilon) \cdot \nabla w(x) \chi_{\Omega_n^\varepsilon}(x)$$

for $y' \in Y'$ and $x \in \Omega$. Taking the limit as $\varepsilon \rightarrow 0$ yields

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_{\Omega \times Y'} [(y_c^\varepsilon \cdot \nabla w)(y^1) - (y_c^\varepsilon \cdot \nabla w)(y^0)] \psi dy' dx &= \int_{\Omega \times Y'} D_d(x) \cdot \nabla w \psi dy' dx \\ &= - \int_{\Omega \times Y'} w(x) [D_d(x) \cdot \nabla \psi(x, y') + \operatorname{div} D_d(x) \psi(x, y')] dy' dx. \end{aligned}$$

Then using the convergence of $V^\varepsilon - y_c^\varepsilon \cdot \nabla w$ to \tilde{w}_1 in $L^p(\Omega; W^{1,p}(Y))$ we obtain

$$\begin{aligned} \int_{\Omega} \int_{Y'} [\tilde{w}_1(x, (y', 1)) - \tilde{w}_1(x, (y', 0))] \psi(x, y') dy' dx &= \lim_{\varepsilon \rightarrow 0} \int_{\Omega} \int_{Y'} [V^\varepsilon(x, (y', 1)) \\ - (y_c^\varepsilon \cdot \nabla w)(x, (y', 1)) - V^\varepsilon(x, (y', 0)) + (y_c^\varepsilon \cdot \nabla w)(x, (y', 0))] \psi(x, y') dy' dx &= 0. \end{aligned}$$

Carrying out similar calculations for y_j with $j = 1, \dots, d-1$ yields the Y -periodicity of \tilde{w}_1 and, hence, Y_x -periodicity of w_1 , defined by $w_1(x, y) = \tilde{w}_1(x, D_x^{-1}y)$ for $x \in \Omega$ and $y \in D_x Y$. \square

6. Micro-macro decomposition: The interpolation operator $\mathcal{Q}_\mathcal{L}^\varepsilon$. Similar to the periodic case [20, 22], in the context of convergence results for the unfolding method in perforated domains as well as for the derivation of error estimates, [28, 31, 32, 33, 47], it is important to consider micro-macro decomposition of a function in $W^{1,p}$ and to introduce an interpolation operator $\mathcal{Q}_\mathcal{L}^\varepsilon$. For any $\varphi \in W^{1,p}(\Omega)$ we consider the splitting $\varphi = \mathcal{Q}_\mathcal{L}^\varepsilon(\varphi) + \mathcal{R}_\mathcal{L}^\varepsilon(\varphi)$ and show that $\mathcal{Q}_\mathcal{L}^\varepsilon(\varphi)$ has a similar behavior as φ , whereas $\mathcal{R}_\mathcal{L}^\varepsilon(\varphi)$ is of order ε .

We consider a continuous extension operator $\mathcal{P} : W^{1,p}(\Omega) \rightarrow W^{1,p}(\mathbb{R}^d)$ satisfying

$$\|\mathcal{P}(\varphi)\|_{W^{1,p}(\mathbb{R}^d)} \leq C \|\varphi\|_{W^{1,p}(\Omega)} \quad \text{for all } \varphi \in W^{1,p}(\Omega),$$

where the constant C depends only on p and Ω , see e.g. [29]. In the following we use the same notation for a function in $W^{1,p}(\Omega)$ and its continuous extension into \mathbb{R}^d .

We consider $\mathcal{Y} = \operatorname{Int}(\bigcup_{k \in \{0,1\}^d} (\bar{Y} + k))$ and define

$$\begin{aligned} \Omega_{\mathcal{Y}}^\varepsilon &= \operatorname{Int}\left(\bigcup_{n=1}^{N_\varepsilon} \bar{\Omega}_{n,\mathcal{Y}}^\varepsilon\right), \quad \text{with } \Omega_{n,\mathcal{Y}}^\varepsilon = \operatorname{Int}\left(\bigcup_{\xi \in \Xi_{n,\mathcal{Y}}^\varepsilon} \varepsilon D_{x_n^\varepsilon}(\bar{Y} + \xi)\right), \\ \Lambda_{\mathcal{Y}}^\varepsilon &= \Omega \setminus \Omega_{\mathcal{Y}}^\varepsilon, \quad \tilde{\Omega}_b^\varepsilon = \operatorname{Int}\left(\bigcup_{n=1}^{N_\varepsilon} \bigcup_{\xi \in \Xi_{n,b}^\varepsilon} \varepsilon D_{x_n^\varepsilon}(\bar{Y} + \xi)\right) \cap \Omega, \end{aligned}$$

where $\Xi_{n,\mathcal{Y}}^\varepsilon = \{\xi \in \Xi_n^\varepsilon : \varepsilon D_{x_n^\varepsilon}(\mathcal{Y} + \xi) \subset (\Omega_n^\varepsilon \cap \Omega)\}$ and $\Xi_{n,b}^\varepsilon = \{\xi \in \Xi_n^\varepsilon : \varepsilon D_{x_n^\varepsilon}(\bar{Y} + \xi) \cap \partial\Omega \neq \emptyset \text{ or } \varepsilon D_{x_n^\varepsilon}(\mathcal{Y} + \xi) \cap \partial\Omega \neq \emptyset\}$.

In order to define an interpolation between two neighboring Ω_n^ε and Ω_m^ε we introduce $\mathcal{Y}^- = \operatorname{Int}(\bigcup_{k \in \{0,1\}^d} (\bar{Y} - k))$.

For $1 \leq n \leq N_\varepsilon$ and $m \in Z_n = \{1 \leq m \leq N_\varepsilon : \partial\Omega_n^\varepsilon \cap \partial\Omega_m^\varepsilon \neq \emptyset\}$ we shall consider unit cells near the corresponding neighboring parts of the boundaries $\partial\Omega_n^\varepsilon$ and $\partial\Omega_m^\varepsilon$, respectively. For $\xi_n \in \Xi_n^\varepsilon$, where $\Xi_n^\varepsilon = \{\xi \in \Xi_n^\varepsilon : \varepsilon D_{x_n^\varepsilon}(\bar{Y} + \xi) \cap \partial\Omega_n^\varepsilon \neq \emptyset\}$, we consider

$$\tilde{\Xi}_{n,m}^\varepsilon = \{\xi_m \in \Xi_m^\varepsilon : \varepsilon D_{x_n^\varepsilon}(\mathcal{Y} + \xi_n) \cap \varepsilon D_{x_m^\varepsilon}(\mathcal{Y}^- + \xi_m) \neq \emptyset\}$$

FIG. 4. Schematic diagram of the covering of Ω by Ω_n^ε , of $D_{x_n^\varepsilon}\mathcal{Y}$ and $D_{x_n^\varepsilon}\mathcal{Y}^-$, and of the interpolation points ξ_n and ξ_m for $\mathcal{Q}_{\mathcal{L}}^\varepsilon$ and $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}$.

and

$$\hat{K}_n = \{k \in \{0, 1\}^d : \xi_n + k \in \bar{\Xi}_n^\varepsilon\}, \quad \hat{K}_m^- = \{k \in \{0, 1\}^d : \xi_m - k \in \bar{\Xi}_m^\varepsilon\}.$$

One of the important part in the definition of $\mathcal{Q}_{\mathcal{L}}^\varepsilon$ is to define an interpolation between neighboring Ω_n^ε and Ω_m^ε . For two neighboring Ω_n^ε and Ω_m^ε we consider triangular interpolations between such vertices of $\varepsilon D_{x_n^\varepsilon}(Y + \xi_n)$ and $\varepsilon D_{x_m^\varepsilon}(Y + \xi_m)$ that are lying on $\partial\Omega_{n,\mathcal{Y}}^\varepsilon$ and $\partial\Omega_{m,\mathcal{Y}}^\varepsilon$, respectively.

DEFINITION 6.1. *The operator $\mathcal{Q}_{\mathcal{L}}^\varepsilon : L^p(\Omega) \rightarrow W^{1,\infty}(\Omega)$ is defined by*

$$(6.1) \quad \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi) = \int_Y \varphi(D_{x_n^\varepsilon}(\varepsilon\xi + \varepsilon y)) dy \quad \text{for } \xi \in \bar{\Xi}_n^\varepsilon \text{ and } 1 \leq n \leq N_\varepsilon,$$

and for $x \in \Omega_{n,\mathcal{Y}}^\varepsilon \cup \tilde{\Omega}_b^\varepsilon$ we define $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(x)$ as the Q_1 -interpolant of $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi)$ at the vertices of $\varepsilon[D_{x_n^\varepsilon}^{-1}x/\varepsilon]_Y + \varepsilon Y$, where $1 \leq n \leq N_\varepsilon$.

For $x \in \Lambda_{\mathcal{Y}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon$ we define $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(x)$ as a triangular Q_1 -interpolant of the values of $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi)$ at $\xi_n + k_n$ and ξ_m such that $\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon$ for $m \in Z_n$ and $k_n \in \hat{K}_n$, where $1 \leq n \leq N_\varepsilon$.

The vertices of $\varepsilon D_{x_n^\varepsilon}(Y + \xi_n + k_n)$ and $\varepsilon D_{x_m^\varepsilon}(Y + \xi_m)$ for $\xi_n \in \bar{\Xi}_n^\varepsilon$, $\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon$ and $k_n \in \hat{K}_n$, in the definition of $\mathcal{Q}_{\mathcal{L}}^\varepsilon$, belong to $\partial\Omega_{n,\mathcal{Y}}^\varepsilon$ and $\partial\Omega_{m,\mathcal{Y}}^\varepsilon$, see Figure 4.

For $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)$ and $\mathcal{R}_{\mathcal{L}}^\varepsilon(\varphi) = \varphi - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)$ we have the following estimates.

LEMMA 6.2. *For every $\varphi \in W^{1,p}(\Omega)$, where $1 \leq p < \infty$, we have*

$$(6.2) \quad \begin{aligned} \|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Omega)} &\leq C\|\varphi\|_{L^p(\Omega)}, & \|\mathcal{R}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Omega)} &\leq C\varepsilon\|\nabla\varphi\|_{L^p(\Omega)}, \\ \|\nabla\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Omega)} + \|\nabla\mathcal{R}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Omega)} &\leq C\|\nabla\varphi\|_{L^p(\Omega)}, \end{aligned}$$

where the constant C is independent of ε and depends only on Y , D , and $d = \dim(\Omega)$.

Proof. Similar to the periodic case [20], we use the fact that the space of Q_1 -interpolants is a finite-dimensional space of dimension 2^d and all norms are equivalent. Then for $\xi \in \bar{\Xi}_{n,\mathcal{Y}}^\varepsilon \cup \Xi_{n,b}^\varepsilon$, where $n = 1, \dots, N_\varepsilon$, we obtain

$$(6.3) \quad \|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))}^p \leq C_1 \varepsilon^d \sum_{k \in \{0,1\}^d} |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi + \varepsilon k)|^p.$$

For $\xi_n \in \tilde{\Xi}_n^\varepsilon$ and triangular elements $\omega_{\xi_n, m}^\varepsilon$ between $\Omega_{n, \mathcal{Y}}^\varepsilon$ and $\Omega_{m, \mathcal{Y}}^\varepsilon$, with $m \in Z_n$, holds

$$\|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\omega_{\xi_n, m}^\varepsilon)}^p \leq C_2 \varepsilon^d \sum_{k \in \hat{K}_n} \sum_{\substack{m \in Z_n \\ \xi_m \in \tilde{\Xi}_{n, m}^\varepsilon}} \left[|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon \xi_n + \varepsilon k)|^p + |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon \xi_m)|^p \right],$$

where $|Z_n| \leq 2^d$ and $|\tilde{\Xi}_{n, m}^\varepsilon| \leq 2^{2(d-1)}$ for every $n = 1, \dots, N_\varepsilon$. Thus for $\Lambda_{\mathcal{Y}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon$ holds

$$(6.4) \quad \begin{aligned} & \|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Lambda_{\mathcal{Y}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon)}^p \\ & \leq C_3 \varepsilon^d \sum_{n=1}^{N_\varepsilon} \sum_{\xi_n \in \tilde{\Xi}_n^\varepsilon, k \in \hat{K}_n} \sum_{m \in Z_n, \xi_m \in \tilde{\Xi}_{n, m}^\varepsilon} \left[|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon \xi_n + \varepsilon k)|^p + |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon \xi_m)|^p \right]. \end{aligned}$$

From the definition of $\mathcal{Q}_{\mathcal{L}}^\varepsilon$ it follows that

$$|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon \xi)|^p \leq \int_Y |\varphi(\varepsilon D_{x_n^\varepsilon}(\xi + y))|^p dy = \frac{1}{\varepsilon^d |D_{x_n^\varepsilon} Y|} \int_{\varepsilon D_{x_n^\varepsilon}(\xi + Y)} |\varphi(x)|^p dx$$

for $\xi \in \Xi_n^\varepsilon$ and $n = 1, \dots, N_\varepsilon$. Then using (6.3) and (6.4) implies

$$(6.5) \quad \|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi + Y))}^p \leq C_4 \sum_{k \in \{0, 1\}^d} \int_{\varepsilon D_{x_n^\varepsilon}(\xi + k + Y)} |\varphi(x)|^p dx$$

for $\xi \in \Xi_{n, \mathcal{Y}}^\varepsilon \cup \Xi_{n, b}^\varepsilon$ and $n = 1, \dots, N_\varepsilon$, and in $\Lambda_{\mathcal{Y}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon$ we have

$$(6.6) \quad \|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Lambda_{\mathcal{Y}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon)}^p \leq C_5 \sum_{n=1}^{N_\varepsilon} \sum_{m \in Z_n} \sum_{j=n, m} \sum_{\xi \in \tilde{\Xi}_j^\varepsilon} \int_{\varepsilon D_{x_j^\varepsilon}(\xi + Y)} |\varphi(x)|^p dx.$$

Summing up in (6.5) over $\xi \in \Xi_{n, \mathcal{Y}}^\varepsilon \cup \Xi_{n, b}^\varepsilon$ and $n = 1, \dots, N_\varepsilon$, and adding (6.6) we obtain the estimate for the L^p -norm of $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)$, stated in the Lemma.

From the definition of Q_1 -interpolants we obtain that for $\xi \in \Xi_{n, \mathcal{Y}}^\varepsilon \cup \Xi_{n, b}^\varepsilon$

$$(6.7) \quad \|\nabla \mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi + Y))}^p \leq C \varepsilon^{d-p} \sum_{k \in \{0, 1\}^d} |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon \xi + \varepsilon k) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon \xi)|^p.$$

For the triangular regions $\omega_{\xi_n, m}^\varepsilon$ between neighboring $\Omega_{n, \mathcal{Y}}^\varepsilon$ and $\Omega_{m, \mathcal{Y}}^\varepsilon$ we have

$$\begin{aligned} \|\nabla \mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)\|_{L^p(\omega_{\xi_n, m}^\varepsilon)}^p & \leq C \varepsilon^{d-p} \sum_{\substack{m \in Z_n \\ \xi_m \in \tilde{\Xi}_{n, m}^\varepsilon}} \sum_{k_n \in \hat{K}_n, k_m \in \hat{K}_m} \left[|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon(\xi_n + k_n)) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon \xi_n)|^p \right. \\ & \left. + |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon(\xi_n + k_n)) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon(\xi_m - k_m))|^p + |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon(\xi_m - k_m)) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon \xi_m)|^p \right]. \end{aligned}$$

For $\phi \in W^{1,p}(D_{x_n^\varepsilon} Y)$, using the regularity of D and the Poincaré inequality, we obtain

$$(6.8) \quad \begin{aligned} & \left\| \phi - \int_{D_{x_n^\varepsilon} Y} \phi dy \right\|_{L^p(D_{x_n^\varepsilon} Y)} \leq C \|\nabla_y \phi\|_{L^p(D_{x_n^\varepsilon} Y)}, \\ & \left| \int_{D_{x_n^\varepsilon} Y} \phi dy - \int_{D_{x_n^\varepsilon} \mathcal{Y}} \phi dy \right|^p + \left| \int_{D_{x_n^\varepsilon} \mathcal{Y}} \phi dy - \int_{D_{x_n^\varepsilon} \mathcal{Y}^+} \phi dy \right|^p \leq C \|\nabla_y \phi\|_{L^p(D_{x_n^\varepsilon} \mathcal{Y})}^p, \\ & \left| \int_{D_{x_n^\varepsilon} Y} \phi dy - \int_{D_{x_n^\varepsilon} \mathcal{Y}^-} \phi dy \right|^p + \left| \int_{D_{x_n^\varepsilon} \mathcal{Y}^-} \phi dy - \int_{D_{x_n^\varepsilon} \mathcal{Y}^-} \phi dy \right|^p \leq C \|\nabla_y \phi\|_{L^p(D_{x_n^\varepsilon} \mathcal{Y}^-)}^p, \end{aligned}$$

where $1 \leq n \leq N_\varepsilon$, $k \in \{0, 1\}^d$ and the constant C depends on D and is independent of ε and n . Using a scaling argument we obtain for every $\xi \in \Xi_n^\varepsilon$

$$(6.9) \quad \left\| \phi - \int_{\varepsilon D_{x_n^\varepsilon}(\xi+Y)} \phi dx \right\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))} \leq C\varepsilon \|\nabla \phi\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))}.$$

Hence, for $\xi \in \Xi_{n,Y}^\varepsilon \cup \Xi_{n,b}^\varepsilon$ and $k \in \{0, 1\}^d$ as well as for $\xi_j \in \tilde{\Xi}_j^\varepsilon$, with $j = n, m$ and $k_n \in \hat{K}_n$, $k_m \in \hat{K}_m^-$ we have

$$(6.10) \quad \begin{aligned} |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi + \varepsilon k) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi)|^p &= \left| \int_{Y+k} \varphi(\varepsilon D_{x_n^\varepsilon}(\xi + y)) dy - \int_Y \varphi(\varepsilon D_{x_n^\varepsilon}(\xi + y)) dy \right|^p \\ &\leq C\varepsilon^{p-d} \|\nabla \varphi\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))}^p, \\ |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi_n + \varepsilon k_n) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi_n)|^p &\leq C\varepsilon^{p-d} \|\nabla \varphi\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi_n+Y))}^p, \\ |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi_m - \varepsilon k_m) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi_m)|^p &\leq C\varepsilon^{p-d} \|\nabla \varphi\|_{L^p(\varepsilon D_{x_m^\varepsilon}(\xi_m+Y^-))}^p, \end{aligned}$$

where C depends on D and is independent of ε , n , and m .

For $\xi_n \in \tilde{\Xi}_n^\varepsilon$, $\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon$ and $k_n \in \hat{K}_n$, $k_m \in \hat{K}_m^-$, using the fact $\varepsilon D_{x_m^\varepsilon}(\xi_m + Y^-) \cap \varepsilon D_{x_n^\varepsilon}(\xi_n + Y) \neq \emptyset$, and applying the inequalities (6.8) with a connected domain

$$\tilde{\mathcal{Y}}_{\xi_n} = \bigcup_{m \in Z_n, \xi_m \in \tilde{\Xi}_{n,m}^\varepsilon} \bigcup_{k \in \{0,1\}^d} D_{x_m^\varepsilon}(\xi_m + Y^- + k) \cup D_{x_n^\varepsilon}(\xi_n + Y - k),$$

instead of \mathcal{Y} and \mathcal{Y}^- , together with a scaling argument, yield

$$(6.11) \quad \begin{aligned} |\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi_n + \varepsilon k_n) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi_m - \varepsilon k_m)|^p &\leq \left| \int_{D_{x_n^\varepsilon}(\xi_n+Y+k_n)} \varphi(\varepsilon y) dy - \int_{\tilde{\mathcal{Y}}_{\xi_n}} \varphi(\varepsilon y) dy \right|^p \\ &+ \left| \int_{D_{x_m^\varepsilon}(\xi_m+Y-k_m)} \varphi(\varepsilon y) dy - \int_{\tilde{\mathcal{Y}}_{\xi_n}} \varphi(\varepsilon y) dy \right|^p \leq C\varepsilon^{p-d} \|\nabla \varphi\|_{L^p(\varepsilon \tilde{\mathcal{Y}}_{\xi_n})}^p, \end{aligned}$$

where C depends on D and is independent of ε , n , and m . Thus in $\Lambda_{\tilde{\mathcal{Y}}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon$ we have

$$(6.12) \quad \|\nabla \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Lambda_{\tilde{\mathcal{Y}}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon)}^p \leq C_1 \sum_{n=1}^{N_\varepsilon} \sum_{\xi_n \in \tilde{\Xi}_n^\varepsilon} \sum_{\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon} \|\nabla \varphi\|_{L^p(\varepsilon \tilde{\mathcal{Y}}_{\xi_n})}^p \leq C_2 \|\nabla \varphi\|_{L^p(\Omega)}^p.$$

Applying (6.10) in (6.7), summing up over $\xi \in \Xi_{n,Y}^\varepsilon \cup \Xi_{n,b}^\varepsilon$ and $n = 1, \dots, N_\varepsilon$ and combining with the estimate for $\|\nabla \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Lambda_{\tilde{\mathcal{Y}}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon)}$ in (6.12) we obtain the estimate for $\|\nabla \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Omega)}$ in terms of $\|\nabla \varphi\|_{L^p(\Omega)}$, as stated in the Lemma.

To show the estimates for $\mathcal{R}_{\mathcal{L}}^\varepsilon(\varphi)$ we consider first

$$\begin{aligned} \|\varphi(x) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(x)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))} &\leq \|\varphi(x) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))} \\ &+ \|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(x)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))} \end{aligned}$$

for $\xi \in \Xi_{n,Y}^\varepsilon \cup \Xi_{n,b}^\varepsilon$. Using the definition of $\mathcal{Q}_{\mathcal{L}}^\varepsilon$ and (6.9) we obtain for $\xi \in \Xi_{n,Y}^\varepsilon \cup \Xi_{n,b}^\varepsilon$

$$\|\varphi - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))} \leq C\varepsilon \|\nabla \varphi\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))}.$$

The definition of $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)$ and the properties of Q_1 -interpolants along with (6.10) imply

$$\|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon\xi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))} \leq C\varepsilon \|\nabla \varphi\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi+Y))} \quad \text{for } \xi \in \Xi_{n,Y}^\varepsilon \cup \Xi_{n,b}^\varepsilon.$$

For triangular elements $\omega_{\xi_n, m}^\varepsilon \subset \Lambda_{\mathcal{Y}}^\varepsilon \setminus \tilde{\Omega}_b^\varepsilon$ with $\xi_n \in \tilde{\Xi}_n^\varepsilon$ and $\xi_m \in \tilde{\Xi}_{n, m}^\varepsilon$ we have $\omega_{\xi_n, m}^\varepsilon \subset \varepsilon \tilde{\mathcal{Y}}_{\xi_n}$. Then, the inequalities in (6.8) with $\tilde{\mathcal{Y}}_{\xi_n}$ and a scaling argument imply

$$\|\varphi(x) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon \xi_n)\|_{L^p(\omega_{\xi_n, m}^\varepsilon)} \leq \|\varphi(x) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon \xi_n)\|_{L^p(\varepsilon \tilde{\mathcal{Y}}_{\xi_n})} \leq C\varepsilon \|\nabla \varphi\|_{L^p(\varepsilon \tilde{\mathcal{Y}}_{\xi_n})},$$

whereas (6.10) and (6.11) together with the properties of Q_1 -interpolants ensure

$$\|\mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi) - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)(\varepsilon \xi_n)\|_{L^p(\omega_{\xi_n, m}^\varepsilon)}^p \leq C\varepsilon^p \|\nabla \varphi\|_{L^p(\varepsilon \tilde{\mathcal{Y}}_{\xi_n})}^p.$$

Thus, combining the estimates from above we obtain the following estimate

$$\begin{aligned} \|\mathcal{R}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Omega)} &\leq \sum_{n=1}^{N_\varepsilon} \|\varphi - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\Omega_n^\varepsilon)} \leq \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \Xi_{n, \mathcal{Y}}^\varepsilon \cup \Xi_{n, b}^\varepsilon} \|\varphi - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\xi + Y))} \\ &\quad + \sum_{n=1}^{N_\varepsilon} \sum_{\xi_n \in \tilde{\Xi}_n^\varepsilon, m \in Z_n} \sum_{\xi_m \in \tilde{\Xi}_{n, m}^\varepsilon} \|\varphi - \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)\|_{L^p(\omega_{\xi_n, m}^\varepsilon)} \leq C\varepsilon \|\nabla \varphi\|_{L^p(\Omega)}. \end{aligned}$$

The estimate for $\nabla \mathcal{Q}_{\mathcal{L}}^\varepsilon(\varphi)$ and the definition of $\mathcal{R}_{\mathcal{L}}^\varepsilon(\varphi)$ yield the estimate for $\nabla \mathcal{R}_{\mathcal{L}}^\varepsilon(\varphi)$. \square

To show convergence results for sequences obtained by applying the l-p unfolding operator to sequences of functions defined on locally-periodic perforated domains, we have to introduce the interpolation operator $\mathcal{Q}_{\mathcal{L}}^{*, \varepsilon}$ for functions in $L^p(\Omega_\varepsilon^*)$. We define

$$\hat{\Omega}_\varepsilon^* = \text{Int}\left(\bigcup_{n=1}^{N_\varepsilon} \hat{\Omega}_n^{*, \varepsilon}\right), \quad \Lambda_\varepsilon^* = \Omega_\varepsilon^* \setminus \hat{\Omega}_\varepsilon^*, \quad \text{where } \hat{\Omega}_n^{*, \varepsilon} = \bigcup_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon D_{x_n^\varepsilon}(\bar{Y}^* + \xi),$$

and

$$\Omega_{\varepsilon, \mathcal{Y}}^* = \text{Int}\left(\bigcup_{n=1}^{N_\varepsilon} \bar{\Omega}_{n, \mathcal{Y}}^{*, \varepsilon}\right), \quad \Lambda_{\varepsilon, \mathcal{Y}}^* = \Omega_\varepsilon^* \setminus \Omega_{\varepsilon, \mathcal{Y}}^*, \quad \text{where } \bar{\Omega}_{n, \mathcal{Y}}^{*, \varepsilon} = \text{Int}\left(\bigcup_{\xi \in \Xi_{n, \mathcal{Y}}^\varepsilon} \varepsilon D_{x_n^\varepsilon}(\bar{Y}^* + \xi)\right),$$

as well as $\tilde{\Omega}_\varepsilon^* = \Omega_\varepsilon^* \cap \tilde{\Omega}_\varepsilon$, where $\tilde{\Omega}_\varepsilon$ is defined as

$$(6.13) \quad \tilde{\Omega}_\varepsilon = \{x \in \Omega : \text{dist}(x, \partial\Omega) > 4\varepsilon \max_{x \in \partial\Omega} \text{diam}(D(x)Y)\}.$$

We also consider $\mathcal{Y}^* = \text{Int}(\bigcup_{k \in \{0, 1\}^d} (\bar{Y}^* + k))$ and $\mathcal{Y}^{*, -} = \text{Int}(\bigcup_{k \in \{0, 1\}^d} (\bar{Y}^* - k))$.

Similar to $\mathcal{Q}_{\mathcal{L}}^\varepsilon$, in the definition of the interpolation operator $\mathcal{Q}_{\mathcal{L}}^{*, \varepsilon}$ we shall distinguish between $\Omega_{\mathcal{Y}}^\varepsilon$ and $\Lambda_{\mathcal{Y}}^\varepsilon \cap \tilde{\Omega}_\varepsilon$. For $x \in \Omega_{\mathcal{Y}}^\varepsilon$ we can consider Q_1 -interpolation between vertices of the corresponding unit cells, whereas for $x \in \Lambda_{\mathcal{Y}}^\varepsilon \cap \tilde{\Omega}_\varepsilon$ we consider triangular Q_1 -interpolation between vertices of unit cells in two neighboring Ω_n^ε and Ω_m^ε . This approach ensures that $\mathcal{Q}_{\mathcal{L}}^{*, \varepsilon}(\phi)$ is continuous in $\tilde{\Omega}_\varepsilon$.

DEFINITION 6.3. *The operator $\mathcal{Q}_{\mathcal{L}}^{*, \varepsilon} : L^p(\Omega_\varepsilon^*) \rightarrow W^{1, \infty}(\tilde{\Omega}_\varepsilon)$ is defined by*

$$(6.14) \quad \mathcal{Q}_{\mathcal{L}}^{*, \varepsilon}(\phi)(\varepsilon \xi) = \int_{Y^*} \phi(D_{x_n^\varepsilon}(\varepsilon \xi + \varepsilon y)) dy \quad \text{for } \xi \in \hat{\Xi}_n^\varepsilon \text{ and } n = 1, \dots, N_\varepsilon,$$

and for $x \in \Omega_{n, \mathcal{Y}}^\varepsilon \cap \tilde{\Omega}_\varepsilon$ we define $\mathcal{Q}_{\mathcal{L}}^{*, \varepsilon}(\phi)(x)$ as the Q_1 -interpolant of the values of $\mathcal{Q}_{\mathcal{L}}^{*, \varepsilon}(\phi)(\varepsilon \xi)$ at vertices of $\varepsilon[D_{x_n^\varepsilon}^{-1}x/\varepsilon]_Y + \varepsilon Y$, where $1 \leq n \leq N_\varepsilon$.

For $x \in \Lambda_{\mathcal{Y}}^\varepsilon \cap \tilde{\Omega}_\varepsilon$ we define $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(x)$ as a triangular Q_1 -interpolant of the values of $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)(\varepsilon\xi)$ at $\xi_n + k_n$ and ξ_m such that $\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon$ for $m \in Z_n$ and $k_n \in \hat{K}_n$, where $1 \leq n \leq N_\varepsilon$, see Figure 4.

In a similar way as for $\mathcal{Q}_{\mathcal{L}}^\varepsilon(\phi)$ and $\mathcal{R}_{\mathcal{L}}^\varepsilon(\phi)$ we obtain estimates for $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)$ and $\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(\phi) = \phi - \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)$.

LEMMA 6.4. For every $\phi \in W^{1,p}(\Omega_\varepsilon^*)$, where $1 \leq p < \infty$, we have

$$\begin{aligned} \|\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)\|_{L^p(\tilde{\Omega}_\varepsilon)} &\leq C\|\phi\|_{L^p(\Omega_\varepsilon^*)}, & \|\nabla\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)\|_{L^p(\tilde{\Omega}_\varepsilon)} &\leq C\|\nabla\phi\|_{L^p(\Omega_\varepsilon^*)}, \\ \|\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(\phi)\|_{L^p(\tilde{\Omega}_\varepsilon^*)} &\leq C\varepsilon\|\nabla\phi\|_{L^p(\Omega_\varepsilon^*)}, & \|\nabla\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(\phi)\|_{L^p(\tilde{\Omega}_\varepsilon^*)} &\leq C\|\nabla\phi\|_{L^p(\Omega_\varepsilon^*)}, \end{aligned}$$

where the constant C is independent of ε .

Proof. The proof for the first estimate follows the same lines as the proof of the corresponding estimate in Lemma 6.2. To show the estimates for $\nabla\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)$ and $\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(\phi)$ we have to estimate the differences $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi) - \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi + k)$ for $\xi \in \tilde{\Xi}_{n,\mathcal{Y}}^\varepsilon$ and $k \in \{0,1\}^d$, and $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi_n + \varepsilon k_n) - \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi_m - \varepsilon k_m)$ for $\xi_n \in \tilde{\Xi}_n^\varepsilon$, $\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon$, with $m \in Z_n$, and $k_n \in \hat{K}_n$, $k_m \in \hat{K}_m^-$, where $1 \leq n \leq N_\varepsilon$. As in the proof of Lemma 6.2, by considering the estimate (6.7), applying the Poincaré inequality and using the estimates similar to (6.10), with Y^* and \mathcal{Y}^* instead of Y and \mathcal{Y} , we obtain

$$\begin{aligned} (6.15) \quad & |\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi) - \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi + k)|^p \leq C\varepsilon^{p-d}\|\nabla\phi\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\mathcal{Y}^* + \xi))}^p, \\ & \|\nabla\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(x)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(Y + \xi))} \leq C\|\nabla\phi\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\mathcal{Y}^* + \xi))}, \\ & \|\phi - \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(Y^* + \xi))} \leq \|\phi - \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(Y^* + \xi))} \\ & \quad + \|\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi) - \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi)\|_{L^p(\varepsilon D_{x_n^\varepsilon}(Y + \xi))} \leq C\varepsilon\|\nabla\phi\|_{L^p(\varepsilon D_{x_n^\varepsilon}(\mathcal{Y}^* + \xi))}, \end{aligned}$$

for $\xi \in \tilde{\Xi}_{n,\mathcal{Y}}^\varepsilon$ and $n = 1, \dots, N_\varepsilon$. For $\xi_n \in \tilde{\Xi}_n^\varepsilon$ and $\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon$, with $m \in Z_n$, we consider sets of $D_{x_j}(Y_0 + \xi)$ for such $D_{x_j}(Y + \xi)$, with $\xi \in \hat{\Xi}_j^\varepsilon$ and $j = n, m$, that have possible nonempty intersections with a triangular element $\omega_{\xi_n, m}^\varepsilon$ between neighboring $\Omega_{n,\mathcal{Y}}^*$ and $\Omega_{m,\mathcal{Y}}^*$

$$\begin{aligned} \hat{\mathcal{Y}}_{\xi_n}^0 &= \bigcup_{k_n^- \in \hat{K}_n^-} \bigcup_{k_n^+ \in \hat{K}_n^+} D_{x_n}(\bar{Y}_0 + \xi_n - k_n^-) \cup D_{x_n}(\bar{Y}_0 + \xi_n + k_n^+), \\ \hat{\mathcal{Y}}_{\xi_n}^{0,-} &= \bigcup_{m \in Z_n, \xi_m \in \tilde{\Xi}_{n,m}^\varepsilon} \bigcup_{\substack{k_m^- \in \hat{K}_m^- \\ l \in Z_m, \xi_l \in \tilde{\Xi}_{l,m}^\varepsilon}} D_{x_m}(\bar{Y}_0 + \xi_m - k_m^-) \cup D_{x_l}(\bar{Y}_0 + \xi_l + k_l^+), \end{aligned}$$

and sets of cells $D_{x_n}(Y + \xi)$ and $D_{x_m}(Y + \xi)$ that have possible nonempty intersections with $\omega_{\xi_n, m}^\varepsilon$

$$\begin{aligned} \hat{\mathcal{Y}}_{\xi_n} &= \text{Int}\left(\bigcup_{k_n^- \in \hat{K}_n^-} D_{x_n}(\bar{Y} + \xi_n - k_n^-) \cup D_{x_n}(\bar{Y} + \xi_n)\right), \\ \hat{\mathcal{Y}}_{\xi_n}^- &= \text{Int}\left(\bigcup_{m \in Z_n, \xi_m \in \tilde{\Xi}_{n,m}^\varepsilon} D_{x_m}(\bar{Y}^- + \xi_m)\right) \end{aligned}$$

and define $\tilde{\mathcal{Y}}_{\xi_n}^* = \text{Int}((\hat{\mathcal{Y}}_{\xi_n} \cup \hat{\mathcal{Y}}_{\xi_n}^-) \setminus (\hat{\mathcal{Y}}_{\xi_n}^0 \cup \hat{\mathcal{Y}}_{\xi_n}^{0,-}))$. We have that $\tilde{\mathcal{Y}}_{\xi_n}^*$ is connected and $\varepsilon\tilde{\mathcal{Y}}_{\xi_n}^* \subset \Omega_\varepsilon^*$ for all $\xi_n \in \tilde{\Xi}_n^\varepsilon$, $n = 1, \dots, N_\varepsilon$. Applying the Poincaré inequality in $\tilde{\mathcal{Y}}_{\xi_n}^*$

and using the regularity of D yields

$$(6.16) \quad \begin{aligned} & \left| \int_{D_{x_n^\varepsilon}(Y^* + \xi_n + k_n)} \phi(y) dy - \int_{\tilde{Y}_{\xi_n}^*} \phi(y) dy \right|^p \leq C \int_{\tilde{Y}_{\xi_n}^*} |\nabla_y \phi(y)|^p dy, \\ & \left| \int_{D_{x_m^\varepsilon}(Y^* + \xi_m - k_m)} \phi(y) dy - \int_{\tilde{Y}_{\xi_n}^*} \phi(y) dy \right|^p \leq C \int_{\tilde{Y}_{\xi_n}^*} |\nabla_y \phi(y)|^p dy, \\ & \left\| \phi - \int_{D_{x_n^\varepsilon}(Y^* + \xi_n)} \phi(y) dy \right\|_{L^p(\tilde{Y}_{\xi_n}^*)} \leq C \|\nabla_y \phi\|_{L^p(\tilde{Y}_{\xi_n}^*)}, \end{aligned}$$

for $\xi_n \in \tilde{\Xi}_n^\varepsilon$, $\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon$, with $m \in Z_n$, and $k_n \in \hat{K}_n$, $k_m \in \hat{K}_m^-$, where the constant C depends on D and is independent of ε , n and m . Then, using a scaling argument in (6.16) implies

$$(6.17) \quad \left| Q_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi_n + \varepsilon k_n) - Q_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi_m - \varepsilon k_m) \right|^p \leq C\varepsilon^{p-d} \|\nabla\phi\|_{L^p(\varepsilon\tilde{Y}_{\xi_n}^*)}^p$$

for $\xi_n \in \tilde{\Xi}_n^\varepsilon$, $\xi_m \in \tilde{\Xi}_{n,m}^\varepsilon$, with $m \in Z_n$, and $k_n \in \hat{K}_n$, $k_m \in \hat{K}_m^-$. Hence, taking into account that $|Z_n| \leq 2^d$ and $|\tilde{\Xi}_{n,m}^\varepsilon| \leq 2^{2(d-1)}$, we obtain

$$(6.18) \quad \|\nabla Q_{\mathcal{L}}^{*,\varepsilon}(\phi)\|_{L^p(\Lambda_{\varepsilon,y}^* \cap \tilde{\Omega}_\varepsilon)}^p \leq C_1 \sum_{n=1}^{N_\varepsilon} \sum_{\xi_n \in \tilde{\Xi}_n^\varepsilon} \|\nabla\phi\|_{L^p(\varepsilon\tilde{Y}_{\xi_n}^*)}^p \leq C_2 \|\nabla\phi\|_{L^p(\Omega_\varepsilon^*)}^p.$$

Applying a scaling argument in (6.16) and using the properties of Q_1 -interpolants and the estimate (6.17) yields

$$(6.19) \quad \begin{aligned} & \|\phi - Q_{\mathcal{L}}^{*,\varepsilon}(\phi)\|_{L^p(\Lambda_{\varepsilon,y}^* \cap \tilde{\Omega}_\varepsilon)} \leq \sum_{n=1}^{N_\varepsilon} \sum_{\xi_n \in \tilde{\Xi}_n^\varepsilon} \left[\|\phi - Q_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi_n)\|_{L^p(\varepsilon\tilde{Y}_{\xi_n}^*)} \right. \\ & \left. + \sum_{m \in Z_n, \xi_m \in \tilde{\Xi}_{n,m}^\varepsilon} \|Q_{\mathcal{L}}^{*,\varepsilon}(\phi)(\varepsilon\xi_n) - Q_{\mathcal{L}}^{*,\varepsilon}(\phi)\|_{L^p(\omega_{\xi_n,m}^\varepsilon)} \right] \leq C\varepsilon \|\nabla\phi\|_{L^p(\Omega_\varepsilon^*)}. \end{aligned}$$

Summing in (6.15) over $\tilde{\Xi}_{n,y}^\varepsilon$ and $1 \leq n \leq N_\varepsilon$, adding (6.18) or (6.19), respectively, and using the definition of $\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(\phi)$ we obtain the estimates stated in the Lemma. \square

7. The l-p unfolding operator in perforated domains: Proofs of convergence results. In this section we prove convergence results for the l-p unfolding operator in domains with locally-periodic perforations. First, we show some properties of the l-p unfolding operator in perforated domains.

LEMMA 7.1.

(i) $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}$ is linear and continuous from $L^p(\Omega_\varepsilon^*)$ to $L^p(\Omega \times Y^*)$, where $p \in [1, \infty)$,

$$\|\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w)\|_{L^p(\Omega \times Y^*)} \leq |Y|^{1/p} \|w\|_{L^p(\Omega_\varepsilon^*)}.$$

(ii) For $w \in L^p(\Omega)$, with $p \in [1, \infty)$, $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w) \rightarrow w$ strongly in $L^p(\Omega \times Y^*)$.

(iii) Let $w^\varepsilon \in L^p(\Omega_\varepsilon^*)$, with $p \in (1, \infty)$, such that $\|w^\varepsilon\|_{L^p(\Omega_\varepsilon^*)} \leq C$. If

$$\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon) \rightharpoonup \hat{w} \quad \text{weakly in } L^p(\Omega \times Y^*),$$

then

$$\tilde{w}^\varepsilon \rightharpoonup \frac{1}{|Y|} \int_{Y^*} \hat{w} dy \quad \text{weakly in } L^p(\Omega).$$

(iv) For $w \in L^p(\Omega; C_{per}(Y_x^*))$ we have $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\mathcal{L}^\varepsilon w) \rightarrow w(\cdot, D_x \cdot)$ in $L^p(\Omega \times Y^*)$, where $p \in [1, \infty)$.

(v) For $w \in C(\bar{\Omega}; L_{per}^p(Y_x^*))$ we have $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\mathcal{L}_0^\varepsilon w) \rightarrow w(\cdot, D_x \cdot)$ in $L^p(\Omega \times Y^*)$, where $p \in [1, \infty)$.

By \tilde{w} we denote the extension of w by zero from Ω_ε^* into Ω .

Proof. [Sketch of the Proof] The proof of (i) follows directly from the definition of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}$ and by using similar calculations as in the proof of Lemma 5.1.

For $w_k \in C_0^\infty(\Omega)$ the convergence in (ii) results from the definition of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}$, the properties of the covering of Ω_ε^* by $\Omega_n^{*,\varepsilon}$ and the following simple calculations

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega \times Y^*} |\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w_k)|^p dy dx = \lim_{\varepsilon \rightarrow 0} \left[\sum_{n=1}^{N_\varepsilon} |\hat{\Omega}_n^\varepsilon| |Y^*| |w_k(x_n^\varepsilon)|^p + \delta_\varepsilon \right] = \int_{\Omega \times Y^*} |w_k(x)|^p dy dx.$$

We used the fact that $|\Lambda_\varepsilon| \rightarrow 0$ as $\varepsilon \rightarrow 0$ and, due to the continuity of w_k , we have

$$\delta^\varepsilon = \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} |Y| \int_{\varepsilon D_{x_n^\varepsilon}(\xi + Y^*)} |w_k(x) - w_k(x_n^\varepsilon)|^p dx \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0.$$

The approximation of $w \in L^p(\Omega)$ by $\{w_k\} \subset C_0^\infty(\Omega)$ and the estimate for the norm of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\varphi)$ in (i), yield the convergence for $w \in L^p(\Omega)$. The proof of the convergence in (iii) is similar to the proof of Lemma 5.2 and the corresponding result for the periodic unfolding operator.

The proof of (iv) follows the same lines as the proof of the corresponding result for $\mathcal{T}_{\mathcal{L}}^\varepsilon$ in Lemma 5.3. In a similar way as in [49, Lemma 3.4] we obtain that

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_{\Omega_\varepsilon^*} |\mathcal{L}_0^\varepsilon(w)(x)|^p dx &= \int_{\Omega} \frac{1}{|Y_x|} \int_{Y_x^*} |w(x, y)|^p dy dx = \int_{\Omega} \frac{1}{|Y|} \int_{Y^*} |w(x, D_x \tilde{y})|^p d\tilde{y} dx, \\ \lim_{\varepsilon \rightarrow 0} \int_{\Lambda_\varepsilon^*} |\mathcal{L}_0^\varepsilon(w)(x)|^p dx &= 0. \end{aligned}$$

Then, the last two convergence results together with the equality

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega \times Y^*} |\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\mathcal{L}_0^\varepsilon w)|^p dy dx = |Y| \lim_{\varepsilon \rightarrow 0} \left[\int_{\Omega_\varepsilon^*} |\mathcal{L}_0^\varepsilon w|^p dx - \int_{\Lambda_\varepsilon^*} |\mathcal{L}_0^\varepsilon w|^p dx \right]$$

imply the convergence result stated in (v). \square

Similar to $\mathcal{T}_{\mathcal{L}}^\varepsilon$ we have $\nabla_y \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w) = \varepsilon \sum_{n=1}^{N_\varepsilon} D_{x_n^\varepsilon}^T \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\nabla w) \chi_{\Omega_n^\varepsilon}$ for $w \in W^{1,p}(\Omega_\varepsilon^*)$. Using the definition and properties of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}$, we prove convergence results for $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$, $\varepsilon \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\nabla w^\varepsilon)$, and $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\nabla w^\varepsilon)$. We start with the proof of Theorem 4.3. Here the proof of the weak convergence follows the same steps as for $\mathcal{T}_{\mathcal{L}}^\varepsilon$ in Theorem 4.1, whereas the periodicity of the limit-function we show in a different way.

Proof. [Proof of Theorem 4.3] The boundedness of $\{\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)\}$ and $\{\nabla_y \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)\}$, ensured by (4.1) and the regularity of D , imply the weak convergences in (4.2). To show the periodicity of w we consider for $\phi \in C_0^\infty(\Omega \times Y^*)$ and $k = 1, \dots, d$

$$\begin{aligned} \int_{\Omega \times Y^*} \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)(x, \tilde{y} + e_k) \phi dx d\tilde{y} &= \int_{\Omega \times Y^*} \sum_{n=1}^{N_\varepsilon} \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon) \phi(x - \varepsilon D_{x_n^\varepsilon} e_k, \tilde{y}) \chi_{\tilde{\Omega}_n^{\varepsilon,k}} dx d\tilde{y} \\ &+ \sum_{n=1}^{N_\varepsilon} \left[\int_{\tilde{\Lambda}_{n,2}^\varepsilon \times Y^*} \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon) \phi dx d\tilde{y} - \int_{\tilde{\Lambda}_{n,1}^\varepsilon \times Y^*} \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)(x, \tilde{y} + e_k) \phi dx d\tilde{y} \right], \end{aligned}$$

where $\tilde{\Omega}_n^{\varepsilon,k}$ and $\tilde{\Lambda}_{n,j}^\varepsilon$, with $j = 1, 2$, are defined in the proof of Theorem 4.1, Section 5, with e_k instead of e_d . Considering the weak convergence of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$ along with the fact that $|\sum_{n=1}^{N_\varepsilon} \tilde{\Lambda}_{n,j}^\varepsilon| \leq C\varepsilon^{1-r}$, for $j = 1, 2$, and taking the limit as $\varepsilon \rightarrow 0$ implies

$$\int_{\Omega \times Y^*} w(x, D_x(\tilde{y} + e_k)) \phi(x, \tilde{y}) d\tilde{y} dx = \int_{\Omega \times Y^*} w(x, D_x \tilde{y}) \phi(x, \tilde{y}) d\tilde{y} dx$$

for all $\phi \in C_0^\infty(\Omega \times Y^*)$ and $k = 1, \dots, d$. Thus, we obtain that $w(x, y)$ is Y_x -periodic. \square

Similar to the periodic case, we use the micro-macro decomposition of a function $\phi \in W^{1,p}(\Omega_\varepsilon^*)$, i.e. $\phi = \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(\phi) + \mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(\phi)$, to show the weak convergence of $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\nabla w^\varepsilon)$. Here we use the fact that for $\{w^\varepsilon\}$ bounded in $W^{1,p}(\Omega_\varepsilon^*)$ the sequence $\{\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)\}$ is bounded in $W^{1,p}(G)$, for every relatively compact open subset $G \subset \Omega$.

Notice that for $w^\varepsilon \in W^{1,p}(\Omega_\varepsilon^*)$ the function $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$ is defined on $\tilde{\Omega}_\varepsilon$. Thus, we can apply $\mathcal{T}_{\mathcal{L}}^\varepsilon$ to $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$ and use the convergence results for the l-p unfolding operator $\mathcal{T}_{\mathcal{L}}^\varepsilon$ (shown in Theorems 4.1 and 4.2) to prove the weak convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)^\sim)$ and $\mathcal{T}_{\mathcal{L}}^\varepsilon([\nabla \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)]^\sim)$, where \sim denotes an extension by zero from $\tilde{\Omega}_\varepsilon$ to Ω .

LEMMA 7.2. *If $\|w^\varepsilon\|_{W^{1,p}(\Omega_\varepsilon^*)} \leq C$, where $p \in (1, \infty)$. Then there exist a subsequence (denoted again by w^ε) and a function $w \in W^{1,p}(\Omega)$ such that*

$$\begin{aligned} \mathcal{T}_{\mathcal{L}}^\varepsilon(\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)^\sim) &\rightarrow w && \text{strongly in } L_{loc}^p(\Omega; W^{1,p}(Y)), \\ \mathcal{T}_{\mathcal{L}}^\varepsilon(\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)^\sim) &\rightharpoonup w && \text{weakly in } L^p(\Omega; W^{1,p}(Y)), \\ \mathcal{T}_{\mathcal{L}}^\varepsilon([\nabla \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)]^\sim) &\rightharpoonup \nabla w && \text{weakly in } L^p(\Omega \times Y). \end{aligned}$$

Proof. Similar to the periodic case [22], the estimates for $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}$ in Lemma 6.4 ensure that there exists a function $w \in W^{1,p}(\Omega)$ such that, up to a subsequence,

$$\begin{aligned} \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)^\sim &\rightarrow w && \text{strongly in } L_{loc}^p(\Omega) \text{ and weakly in } L^p(\Omega), \\ [\nabla \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)]^\sim &\rightharpoonup \nabla w && \text{weakly in } L^p(\Omega). \end{aligned}$$

Then, the first two convergences stated in the Lemma follow directly from the estimates, estimate $\|\nabla_y \mathcal{T}_{\mathcal{L}}^\varepsilon(\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)^\sim)\|_{L^p(\Omega \times Y)} \leq C_1 \varepsilon \|[\nabla \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)]^\sim\|_{L^p(\Omega)} \leq C\varepsilon$, and convergence results for $\mathcal{T}_{\mathcal{L}}^\varepsilon$ in Lemmas 5.1, 5.2 and Theorem 4.1. To prove the final convergence stated in the Lemma we observe that $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)|_G$ is uniformly bounded in $W^{1,p}(G)$, where $G \subset \Omega$ is a relatively compact open set, see Lemma 6.4. Then, by Theorem 4.2 there exists $\hat{w}_{1,G} \in L^p(G; W_{\text{per}}^{1,p}(Y_x))$ such that

$$\mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)|_G) \rightharpoonup \nabla w + D_x^{-T} \nabla_y \hat{w}_{1,G}(\cdot, D_x \cdot) \quad \text{weakly in } L^p(G \times Y).$$

The definition of $\mathcal{Q}_{\mathcal{L}}^\varepsilon$ implies that $\hat{w}_{1,G}$ is a polynomial in y of degree less than or equal to one with respect to each variable y_1, \dots, y_d . Thus, the Y_x -periodicity of $\hat{w}_{1,G}$ yields that it is constant with respect to y and

$$\mathcal{T}_{\mathcal{L}}^\varepsilon([\nabla \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)]^\sim) \rightharpoonup \nabla w \quad \text{weakly in } L_{loc}^p(\Omega; L^p(Y)).$$

The boundedness of $[\nabla \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)]^\sim$ in $L^p(\Omega)$ implies the boundedness of $\mathcal{T}_{\mathcal{L}}^\varepsilon([\nabla \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)]^\sim)$ in $L^p(\Omega \times Y)$ and we obtain the last convergence stated in Lemma. \square

For $\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon) = w^\varepsilon - \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$ we have the following convergence results.

LEMMA 7.3. Consider a sequence $\{w^\varepsilon\} \subset W^{1,p}(\Omega_\varepsilon^*)$, with $p \in (1, \infty)$, satisfying $\|\nabla w^\varepsilon\|_{L^p(\Omega_\varepsilon^*)} \leq C$. Then, there exist a subsequence (denoted again by w^ε) and a function $w_1 \in L^p(\Omega; W_{\text{per}}^{1,p}(Y_x^*))$ such that

$$(7.1) \quad \begin{aligned} \varepsilon^{-1} \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)^\sim) &\rightharpoonup w_1(\cdot, D_x \cdot) && \text{weakly in } L^p(\Omega; W^{1,p}(Y^*)), \\ \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)^\sim) &\rightarrow 0 && \text{strongly in } L^p(\Omega; W^{1,p}(Y^*)), \\ \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}([\nabla \mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)]^\sim) &\rightharpoonup D_x^{-T} \nabla_y w_1(\cdot, D_x \cdot) && \text{weakly in } L^p(\Omega \times Y^*), \end{aligned}$$

where \sim denotes the extension by zero from $\tilde{\Omega}_\varepsilon^*$ to Ω_ε^* .

Proof. The estimates in Lemma 6.4 imply that $\varepsilon^{-1} \mathcal{T}_{\mathcal{L}}^{*,\varepsilon}(\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)^\sim)$ is bounded in $L^p(\Omega; W^{1,p}(Y^*))$ and there exists $\tilde{w}_1 \in L^p(\Omega; W^{1,p}(Y^*))$ and $w_1(x, y) = \tilde{w}_1(x, D_x^{-1}y)$ for $x \in \Omega$, $y \in Y_x^*$, where $Y_x^* = D(x)Y^*$, such that the convergences in (7.1) are satisfied. To show that w_1 is Y_x -periodic we consider the restriction of $\varepsilon^{-1} \mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$ on G_ε^* , which belongs to $W^{1,p}(G_\varepsilon^*)$. Here $G_\varepsilon^* = G \cap \Omega_\varepsilon^*$ and $G \subset \Omega$ is a relatively compact open subset of Ω . Using Lemma 6.4 we obtain

$$\|\varepsilon^{-1} \mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)\|_{L^p(G_\varepsilon^*)} + \varepsilon \|\varepsilon^{-1} \nabla \mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)\|_{L^p(G_\varepsilon^*)} \leq C.$$

Applying Theorem 4.3 to $\varepsilon^{-1} \mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)|_{G_\varepsilon^*}$ yields $w_1|_{G \times Y_x^*} \in L^p(G; W_{\text{per}}^{1,p}(Y_x^*))$. Since G can be chosen arbitrarily we obtain that $w_1 \in L^p(\Omega; W_{\text{per}}^{1,p}(Y_x^*))$. \square

Combining the convergence results from above we obtain directly the main convergence theorem for the l-p unfolding operator in locally-periodic perforated domains.

Proof. [Proof of Theorem 4.4] Similar to the periodic case the convergence results stated in Theorem 4.4 follow directly from the fact that $w^\varepsilon = \mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon) + \mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$ and from the convergence results in Lemmas 7.2 and 7.3. \square

REMARK. In the definition of Ω_ε^* we assumed that there no perforations in layers $(\Omega_n^{*,\varepsilon} \setminus \tilde{\Omega}_n^{*,\varepsilon}) \cap \tilde{\Omega}_{\varepsilon/2}$, with $\tilde{\Omega}_{\varepsilon/2} = \{x \in \Omega : \text{dist}(x, \partial\Omega) > 2\varepsilon \max_{x \in \partial\Omega} \text{diam}(D(x)Y)\}$ and $1 \leq n \leq N_\varepsilon$. In the proofs of convergence results only local estimates for $\mathcal{Q}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$ and $\mathcal{R}_{\mathcal{L}}^{*,\varepsilon}(w^\varepsilon)$ are used, thus no restrictions on the distribution of perforations near $\partial\Omega$ are needed. For the macroscopic description of microscopic processes this assumption is not restrictive since $|\bigcup_{n=1}^{N_\varepsilon} (\Omega_n^{*,\varepsilon} \setminus \tilde{\Omega}_n^{*,\varepsilon}) \cap \Omega| \leq C\varepsilon^{1-r} \rightarrow 0$ as $\varepsilon \rightarrow 0$, $r < 1$. If we allow perforations in layers between two neighboring $\hat{\Omega}_n^{*,\varepsilon}$ and $\hat{\Omega}_m^{*,\varepsilon}$ in $\tilde{\Omega}_{\varepsilon/2}$, then using that $Y^* = Y \setminus \bar{Y}_0$ is connected, the transformation matrix D is Lipschitz continuous and $\text{dist}(\tilde{\Omega}_{\varepsilon/2}, \partial\Omega) > 0$, it is possible to construct an extension of $w^\varepsilon \in W^{1,p}(\Omega_\varepsilon^*)$ from $(\Omega_n^{*,\varepsilon} \setminus \tilde{\Omega}_n^{*,\varepsilon}) \cap \tilde{\Omega}_{\varepsilon/2}$ to $(\Omega_n^\varepsilon \setminus \tilde{\Omega}_n^\varepsilon) \cap \tilde{\Omega}_{\varepsilon/2}$, such that the $W^{1,p}$ -norm of the extension is controlled by the $W^{1,p}$ -norm of the original function, uniform in ε , and apply Lemmas 7.2, 7.3 and Theorem 4.4 to the sequence of extended functions.

8. Two-scale convergence on oscillating surfaces and the l-p boundary unfolding operator. To derive macroscopic equations for the microscopic problems posed on boundaries of locally-periodic microstructures or with non-homogeneous Neumann conditions on boundaries of locally-periodic microstructures we have to show convergence properties for sequences defined on oscillating surfaces. To show the compactness result for l-p two-scale convergence on oscillating surfaces (see Theorem 4.5) we first prove the convergence of the $L^p(\Gamma^\varepsilon)$ -norm of the l-p approximation of $\psi \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$.

LEMMA 8.1. For $\psi \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ and $1 \leq p < \infty$, we have that

$$\lim_{\varepsilon \rightarrow 0} \varepsilon \int_{\Gamma^\varepsilon} |\mathcal{L}^\varepsilon \psi(x)|^p d\sigma_x = \int_{\Omega} \frac{1}{|Y_x|} \int_{\Gamma_x} |\psi(x, y)|^p d\sigma_y dx.$$

Proof. The definition of the l-p approximation \mathcal{L}^ε implies

$$\begin{aligned} \varepsilon \int_{\Gamma^\varepsilon} |\mathcal{L}^\varepsilon \psi|^p d\sigma_x &= \varepsilon \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \int_{\varepsilon \Gamma_{x_n}^\xi} \left| \tilde{\psi}\left(x, \frac{D_{x_n}^{-1} x}{\varepsilon}\right) \right|^p - \left| \tilde{\psi}\left(x_n, \frac{D_{x_n}^{-1} x}{\varepsilon}\right) \right|^p d\sigma_x \\ &+ \varepsilon \sum_{n=1}^{N_\varepsilon} \left[\sum_{\xi \in \hat{\Xi}_n^\varepsilon} \int_{\varepsilon \Gamma_{x_n}^\xi} \left| \tilde{\psi}\left(x_n, \frac{D_{x_n}^{-1} x}{\varepsilon}\right) \right|^p d\sigma_x + \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \int_{\varepsilon \Gamma_{x_n}^\xi} \left| \tilde{\psi}\left(x_n, \frac{D_{x_n}^{-1} x}{\varepsilon}\right) \right|^p \chi_{\Omega_n^\varepsilon} d\sigma_x \right] \\ &= I_1 + I_2 + I_3, \end{aligned}$$

where $\tilde{\Xi}_n^\varepsilon = \Xi_n^\varepsilon \setminus \hat{\Xi}_n^\varepsilon$ and $\Gamma_{x_n}^\xi = D_{x_n}^\varepsilon(\xi + \tilde{\Gamma}_{K_{x_n}^\varepsilon})$. Then, the continuity of ψ , the properties of Ω_n^ε , and the inequality $\left| |a|^p - |b|^p \right| \leq p|a-b|(|a|^{p-1} + |b|^{p-1})$ imply $I_1 \rightarrow 0$ as $\varepsilon \rightarrow 0$. Using the properties of the covering of Ω by $\{\Omega_n^\varepsilon\}_{n=1}^{N_\varepsilon}$ we obtain

$$|I_3| \leq C \sup_{1 \leq n \leq N_\varepsilon} \varepsilon^d |\tilde{\Xi}_n^\varepsilon| |D_{x_n}^\varepsilon \tilde{\Gamma}_{K_{x_n}^\varepsilon}| \leq C \varepsilon^{1-r} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0 \text{ for } 0 \leq r < 1.$$

Considering the properties of the covering of $\hat{\Omega}_n^\varepsilon$ by $D_{x_n}^\varepsilon(Y + \xi)$, where $\xi \in \hat{\Xi}_n^\varepsilon$ and $1 \leq n \leq N_\varepsilon$, and Y -periodicity of $\tilde{\psi}$ the second integral can be rewritten as

$$I_2 = \sum_{n=1}^{N_\varepsilon} \varepsilon^d |\hat{\Xi}_n^\varepsilon| \int_{D_{x_n}^\varepsilon \tilde{\Gamma}_{K_{x_n}^\varepsilon}} |\tilde{\psi}(x_n^\varepsilon, D_{x_n}^{-1} y)|^p d\sigma_y = \sum_{n=1}^{N_\varepsilon} \frac{|\hat{\Omega}_n^\varepsilon|}{|Y_{x_n^\varepsilon}|} \int_{D_{x_n}^\varepsilon \tilde{\Gamma}_{K_{x_n}^\varepsilon}} |\psi(x_n^\varepsilon, y)|^p d\sigma_y.$$

Then, the regularity assumptions on ψ , D and K , the definition of $\hat{\Omega}_n^\varepsilon$ and the properties of the covering of Ω by $\{\Omega_n^\varepsilon\}_{n=1}^{N_\varepsilon}$ imply the convergence result stated in the Lemma. \square

Similar to l-t-s convergence and two-scale convergence for sequences defined on surfaces of periodic microstructures, the convergence of l-p approximations (shown in Lemma 8.1) and the Riesz representation theorem ensure the compactness result for sequences $\{w^\varepsilon\} \subset L^p(\Gamma^\varepsilon)$ with $\varepsilon \|w^\varepsilon\|_{L^p(\Gamma^\varepsilon)} \leq C$.

Proof. [**Proof of Theorem 4.5**] The Banach space $C(\bar{\Omega}; C_{\text{per}}(Y_x))$ is separable and dense in $L^p(\Omega; L^p(\Gamma_x))$. Thus, using the convergence result in Lemma 8.1, the Riesz representation theorem and similar arguments as in [49, Theorem 3.2] we obtain l-t-s convergence of $\{w^\varepsilon\} \subset L^p(\Gamma^\varepsilon)$ to $w \in L^p(Y; L^p_{\text{per}}(\Gamma_x))$, stated in the theorem. \square

Using the structure of $\Omega_{n,K}^{*,\varepsilon}$ and the covering properties of $\Omega_{\varepsilon,K}^*$ by $\{\Omega_{n,K}^{*,\varepsilon}\}_{n=1}^{N_\varepsilon}$ we can derive the trace inequalities for functions defined on Γ^ε . Applying first the trace inequality in $Y_{x_n^\varepsilon, K}^{*,\xi} = D_{x_n}^\varepsilon(\tilde{Y}_{K_{x_n}^\varepsilon}^* + \xi)$, with $\xi \in \hat{\Xi}_n^\varepsilon$, yields

$$\begin{aligned} \|u\|_{L^p(D_{x_n}^\varepsilon(\tilde{\Gamma}_{K_{x_n}^\varepsilon} + \xi))}^p &\leq \mu_\Gamma \left[\|u\|_{L^p(Y_{x_n^\varepsilon, K}^{*,\xi})}^p + \|\nabla u\|_{L^p(Y_{x_n^\varepsilon, K}^{*,\xi})}^p \right], \\ \|u\|_{L^p(D_{x_n}^\varepsilon(\tilde{\Gamma}_{K_{x_n}^\varepsilon} + \xi))}^p &\leq \mu_\Gamma \left[\|u\|_{L^p(Y_{x_n^\varepsilon, K}^{*,\xi})}^p + \int_{Y_{x_n^\varepsilon, K}^{*,\xi}} \int_{Y_{x_n^\varepsilon, K}^{*,\xi}} \frac{|u(y_1) - u(y_2)|^p}{|y_1 - y_2|^{d+\beta p}} dy_1 dy_2 \right], \end{aligned}$$

for $u \in W^{1,p}(Y_{x_n^\varepsilon, K}^{*,\xi})$ or $u \in W^{\beta,p}(Y_{x_n^\varepsilon, K}^{*,\xi})$, for $1/2 < \beta < 1$, respectively, where the constant μ_Γ depends only on D , K , and Y^* , see e.g. [29, 54]. Then, scaling by ε and

summing up over $\xi \in \hat{\Xi}_n^\varepsilon$ and $1 \leq n \leq N_\varepsilon$ implies the estimates

$$(8.1) \quad \varepsilon \|u\|_{L^p(\hat{\Gamma}^\varepsilon)}^p \leq \mu_\Gamma \left[\|u\|_{L^p(\Omega_{\varepsilon,K}^*)}^p + \varepsilon^p \|\nabla u\|_{L^p(\Omega_{\varepsilon,K}^*)}^p \right]$$

for $u \in W^{1,p}(\Omega_{\varepsilon,K}^*)$, $p \in [1, \infty)$,

$$(8.2) \quad \varepsilon \|u\|_{L^p(\hat{\Gamma}^\varepsilon)}^p \leq \mu_\Gamma \left[\|u\|_{L^p(\Omega_{\varepsilon,K}^*)}^p + \varepsilon^{\beta p} \int_{\Omega_{\varepsilon,K}^*} \int_{\Omega_{\varepsilon,K}^*} \frac{|u(x_1) - u(x_2)|^p}{|x_1 - x_2|^{d+\beta p}} dx_1 dx_2 \right]$$

for $u \in W^{\beta,p}(\Omega_{\varepsilon,K}^*)$ with $1/2 < \beta < 1$, $p \in [1, \infty)$,

where the constant μ_Γ depends on D , K , and Y^* and is independent of ε , where

$$\hat{\Gamma}^\varepsilon = \bigcup_{n=1}^{N_\varepsilon} \hat{\Gamma}_n^\varepsilon \quad \text{with} \quad \hat{\Gamma}_n^\varepsilon = \bigcup_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon D_{x_n^\varepsilon}(\tilde{\Gamma}_{K_{x_n^\varepsilon}} + \xi).$$

Since $\Gamma_{x_n^\varepsilon}$ is given by a linear transformation of Γ , for a parametrization $y = y(w)$ of Γ , where $w \in \mathbb{R}^{d-1}$, we obtain by $x(w) = \varepsilon D_{x_n^\varepsilon} K_{x_n^\varepsilon} y(w)$ the parametrization of $\varepsilon \Gamma_{x_n^\varepsilon}$. We consider for Γ that $d\sigma_y = \sqrt{g} dw$ with $w \in \mathbb{R}^{d-1}$ and for $\Gamma_{x_n^\varepsilon}$ we have $d\sigma_x^n = \varepsilon^{d-1} \sqrt{g_{x_n^\varepsilon}} dw$, where $g = \det(g_{ij})$, $g_{x_n^\varepsilon} = \det(g_{x_n^\varepsilon, ij})$ and g_{ij} , $g_{x_n^\varepsilon, ij}$ are the corresponding first fundamental forms (metrics). We have also $\int_{\Gamma^\varepsilon} d\sigma_x^\varepsilon = \sum_{n=1}^{N_\varepsilon} \int_{\Gamma_n^\varepsilon} d\sigma_x^n$ and $\Gamma_x = D(x)K(x)\Gamma$ with $d\sigma_x = \sqrt{g(x)} dw$.

Using the definition of the l-p boundary unfolding operator, the trace inequalities (8.1), and the assumptions on D and K we show the following properties of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}$.

LEMMA 8.2. *For $\psi \in W^{1,p}(\Omega_{\varepsilon,K}^*)$, with $1 \leq p < \infty$, we have*

$$(i) \quad \int_{\Omega \times \Gamma} \sum_{n=1}^{N_\varepsilon} \frac{\sqrt{g_{x_n^\varepsilon}}}{\sqrt{g}|Y_{x_n^\varepsilon}|} |\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\psi)(x, y)|^p \chi_{\Omega_n^\varepsilon} d\sigma_y dx = \varepsilon \int_{\hat{\Gamma}^\varepsilon} |\psi(x)|^p d\sigma_x^\varepsilon,$$

$$(ii) \quad \int_{\Omega \times \Gamma} |\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\psi)(x, y)|^p d\sigma_y dx = \varepsilon \sum_{n=1}^{N_\varepsilon} \int_{\hat{\Gamma}_n^\varepsilon} \frac{\sqrt{g}|Y_{x_n^\varepsilon}|}{\sqrt{g_{x_n^\varepsilon}}} |\psi(x)|^p d\sigma_x^n \leq C \varepsilon \int_{\hat{\Gamma}^\varepsilon} |\psi(x)|^p d\sigma_x^\varepsilon,$$

$$(iii) \quad \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\psi)\|_{L^p(\Omega \times \Gamma)} \leq C (\|\psi\|_{L^p(\Omega_\varepsilon^*)} + \varepsilon \|\nabla \psi\|_{L^p(\Omega_\varepsilon^*)}),$$

where the constant C depends on D and K and is independent of ε .

Proof. Equality (i) follows directly from the definition of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}$, i.e.

$$\begin{aligned} & \int_{\Omega \times \Gamma} \sum_{n=1}^{N_\varepsilon} \frac{\sqrt{g_{x_n^\varepsilon}}}{\sqrt{g}|Y_{x_n^\varepsilon}|} |\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\psi)|^p \chi_{\Omega_n^\varepsilon} d\sigma_y dx \\ &= \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon^d \int_{\Gamma} \frac{\sqrt{g_{x_n^\varepsilon}}}{\sqrt{g}} |\psi(\varepsilon D_{x_n^\varepsilon}(\xi + K_{x_n^\varepsilon} y))|^p d\sigma_y = \varepsilon \int_{\hat{\Gamma}^\varepsilon} |\psi(x)|^p d\sigma_x^\varepsilon. \end{aligned}$$

Similar calculations and the regularity assumptions on D and K imply the equality and the estimate in (ii). The estimate in (iii) is ensured by (ii) and (8.1). \square

Remark. Due to the second estimate in Lemma 8.2 and the assumptions on D and K , the boundedness of $\varepsilon \|w^\varepsilon\|_{L^p(\hat{\Gamma}^\varepsilon)}^p$ implies the boundedness of $\|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(w^\varepsilon)\|_{L^p(\Omega \times \Gamma)}^p$ and, hence, the weak convergence of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(w^\varepsilon)$ in $L^p(\Omega \times \Gamma)$.

Applying the properties of the l-p boundary unfolding operator shown in Lemma 8.2 we prove the relation between the l-t-s convergence on oscillating boundaries and the l-p boundary unfolding operator.

Proof. [**Proof of Theorem 4.6**] Using the definition of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}$ and considering $\psi \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ together with the corresponding $\tilde{\psi} \in C(\bar{\Omega}; C_{\text{per}}(Y))$ yields

$$\begin{aligned} & \int_{\Omega} \int_{\Gamma} \sum_{n=1}^{N_{\varepsilon}} \frac{\sqrt{g_{x_n^{\varepsilon}}}}{\sqrt{g}|Y_{x_n^{\varepsilon}}|} \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(w^{\varepsilon}) \tilde{\psi}(x, K_{x_n^{\varepsilon}} y) \chi_{\Omega_n^{\varepsilon}} d\sigma_y dx \\ &= \sum_{n=1}^{N_{\varepsilon}} \sum_{\xi \in \tilde{\Xi}_{x_n^{\varepsilon}}} \varepsilon \int_{\varepsilon \Gamma_{x_n^{\varepsilon}}^{\xi}} w^{\varepsilon}(z) \tilde{\psi}\left(z, D_{x_n^{\varepsilon}}^{-1} \frac{z}{\varepsilon}\right) d\sigma_z^n \\ &+ \sum_{n=1}^{N_{\varepsilon}} \sum_{\xi \in \tilde{\Xi}_{x_n^{\varepsilon}}} \varepsilon^{1-d} \frac{1}{|Y_{x_n^{\varepsilon}}|} \int_{\varepsilon \Gamma_{x_n^{\varepsilon}}^{\xi}} w^{\varepsilon}(z) \int_{\varepsilon Y_{x_n^{\varepsilon}}^{\xi}} \left[\tilde{\psi}\left(x, D_{x_n^{\varepsilon}}^{-1} \frac{z}{\varepsilon}\right) - \tilde{\psi}\left(z, D_{x_n^{\varepsilon}}^{-1} \frac{z}{\varepsilon}\right) \right] dx d\sigma_z^n, \end{aligned}$$

where $\Gamma_{x_n^{\varepsilon}}^{\xi} = D_{x_n^{\varepsilon}}(\tilde{\Gamma}_{K_{x_n^{\varepsilon}}} + \xi)$ and $Y_{x_n^{\varepsilon}}^{\xi} = D_{x_n^{\varepsilon}}(Y + \xi)$. The continuity of ψ and the boundedness of $\varepsilon \|w^{\varepsilon}\|_{L^p(\Gamma^{\varepsilon})}^p$ ensure the convergence of the last integral to zero as $\varepsilon \rightarrow 0$. Consider first that $w^{\varepsilon} \rightarrow w$ l-t-s. The assumption on w^{ε} , i.e. $\varepsilon \|w^{\varepsilon}\|_{L^p(\Gamma^{\varepsilon})}^p \leq C$, with $p \in (1, \infty)$ ensures that, up to a subsequence, $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(w^{\varepsilon}) \rightharpoonup \hat{w}$ weakly in $L^p(\Omega \times \Gamma)$. Using the continuity of ψ , D , and K , along with $|\Gamma^{\varepsilon} \setminus \hat{\Gamma}^{\varepsilon}| \rightarrow 0$ as $\varepsilon \rightarrow 0$, yields

$$\begin{aligned} & \int_{\Omega} \int_{\Gamma} \frac{\sqrt{g_x}}{|Y_x| \sqrt{g}} \hat{w}(x, D_x K_x y) \tilde{\psi}(x, K_x y) d\sigma_y dx \\ (8.3) \quad &= \lim_{\varepsilon \rightarrow 0} \int_{\Omega} \int_{\Gamma} \sum_{n=1}^{N_{\varepsilon}} \frac{\sqrt{g_{x_n^{\varepsilon}}}}{|Y_{x_n^{\varepsilon}}| \sqrt{g}} \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(w^{\varepsilon}) \tilde{\psi}(x, K_{x_n^{\varepsilon}} y) \chi_{\Omega_n^{\varepsilon}} d\sigma_y dx \\ &= \lim_{\varepsilon \rightarrow 0} \varepsilon \int_{\Gamma^{\varepsilon}} w^{\varepsilon}(x) \mathcal{L}^{\varepsilon}(\psi) d\sigma_x^{\varepsilon} = \int_{\Omega} \frac{1}{|Y_x|} \int_{\Gamma_x} w(x, y) \psi(x, y) d\sigma_y dx \end{aligned}$$

for all $\psi \in C_0^{\infty}(\Omega; C_{\text{per}}^{\infty}(Y_x))$. Applying the coordinate transformation in the integral on the left hand side yields $\hat{w}(x, y) = w(x, y)$ for a.a. $x \in \Omega$, $y \in \Gamma_x$ and, hence the whole sequence $\{\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(w^{\varepsilon})\}$ converges to $w(\cdot, D_x K_x \cdot)$.

Consider $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(w^{\varepsilon}) \rightharpoonup w(\cdot, D_x K_x \cdot)$ in $L^p(\Omega \times \Gamma)$. The boundedness of $\varepsilon \|w^{\varepsilon}\|_{L^p(\Gamma^{\varepsilon})}^p$ implies that, up to a subsequence, $w^{\varepsilon} \rightarrow \hat{w}$ l-t-s and $\hat{w} \in L^p(\Omega; L^p(\Gamma_x))$. Interchanging in (8.3) \hat{w} and w , we obtain that the whole sequence w^{ε} l-t-s converges to w . \square

For functions in $W^{\beta,p}(\Omega)$, with $p \in (1, \infty)$, and $1/2 < \beta$ or for sequences defined on oscillating boundaries and converging in the $L^p(\Gamma^{\varepsilon})$ -norm, scaled by $\varepsilon^{1/p}$, we obtain the strong convergence of the corresponding unfolded sequences.

LEMMA 8.3. *For $u \in W^{\beta,p}(\Omega)$, with $p \in (1, \infty)$, and $1/2 < \beta$, we have*

$$(8.4) \quad \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(u) \rightarrow u \quad \text{strongly in} \quad L^p(\Omega \times \Gamma).$$

If for $\{v^{\varepsilon}\} \subset L^p(\Gamma^{\varepsilon})$ and some $v \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ holds $\varepsilon \|v^{\varepsilon} - \mathcal{L}^{\varepsilon} v\|_{L^p(\Gamma^{\varepsilon})}^p \rightarrow 0$ as $\varepsilon \rightarrow 0$, then

$$(8.5) \quad \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(v^{\varepsilon}) \rightarrow v(\cdot, D_x K_x \cdot) \quad \text{strongly in} \quad L^p(\Omega \times \Gamma).$$

Proof. For an approximation of u by $u_k \in C^1(\bar{\Omega})$ we can write

$$\begin{aligned} \int_{\Omega \times \Gamma} |\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(u_k)|^p d\sigma_y dx &= \sum_{n=1}^{N_\varepsilon} \int_{\Omega \times \Gamma} |u_k(\varepsilon D_{x_n^\varepsilon} [D_{x_n^\varepsilon}^{-1} x / \varepsilon]_Y + \varepsilon D_{x_n^\varepsilon} K_{x_n^\varepsilon} y)|^p \chi_{\hat{\Omega}_n^\varepsilon} d\sigma_y dx \\ &= \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon^d |Y_{x_n^\varepsilon}| \int_{\Gamma} |u_k(\varepsilon D_{x_n^\varepsilon}(\xi + K_{x_n^\varepsilon} y))|^p d\sigma_y = \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} |\varepsilon Y_{x_n^\varepsilon}| |\Gamma| |u_k(\tilde{x}_{n,\xi}^\varepsilon)|^p + \delta_\varepsilon, \end{aligned}$$

for some fixed $\tilde{x}_{n,\xi}^\varepsilon \in \varepsilon D_{x_n^\varepsilon}(K_{x_n^\varepsilon} \Gamma + \xi)$, where, due to the continuity of u_k , we have

$$\delta_\varepsilon = \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon^d |D_{x_n^\varepsilon} Y| \int_{\Gamma} |u_k(\varepsilon D_{x_n^\varepsilon}(\xi + K_{x_n^\varepsilon} y)) - u_k(\tilde{x}_{n,\xi}^\varepsilon)|^p d\sigma_y \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0.$$

The properties of the covering of Ω by $\{\Omega_n^\varepsilon\}_{n=1}^{N_\varepsilon}$ and $|\Omega_n^\varepsilon \setminus \hat{\Omega}_n^\varepsilon| \rightarrow 0$ as $\varepsilon \rightarrow 0$ imply

$$\lim_{\varepsilon \rightarrow 0} \sum_{n=1}^{N_\varepsilon} \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon^d |D_{x_n^\varepsilon} Y| |\Gamma| |u_k(\tilde{x}_{n,\xi}^\varepsilon)|^p = \int_{\Omega} \int_{\Gamma} |u_k(x)|^p d\sigma_y dx.$$

Then, the density of $C^1(\bar{\Omega})$ in $W^{\beta,p}(\Omega)$, the relation (ii) in Lemma 8.2, and the trace estimate (8.2) ensure the convergence result for $u \in W^{\beta,p}(\Omega)$.

To show the convergence in (8.5) we consider

$$\begin{aligned} \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(v^\varepsilon) - v(\cdot, D_x K_x \cdot)\|_{L^p(\Omega \times \Gamma)} &\leq \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(v^\varepsilon) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\mathcal{L}^\varepsilon v)\|_{L^p(\Omega \times \Gamma)} \\ &\quad + \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\mathcal{L}^\varepsilon v) - v(\cdot, D_x K_x \cdot)\|_{L^p(\Omega \times \Gamma)}. \end{aligned}$$

Then, the estimate (ii) in Lemma 8.2, the regularity of v , D , and K , and the convergence

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_{\Omega \times \Gamma} |\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\mathcal{L}^\varepsilon v)|^p d\sigma_y dx &= \lim_{\varepsilon \rightarrow 0} \sum_{n=1}^{N_\varepsilon} |\varepsilon Y_{x_n^\varepsilon}| \sum_{\xi \in \hat{\Xi}_n^\varepsilon} \int_{\Gamma} |\tilde{v}(\varepsilon D_{x_n^\varepsilon}(\xi + K_{x_n^\varepsilon} y), K_{x_n^\varepsilon} y)|^p d\sigma_y \\ &= \int_{\Omega} \int_{\Gamma} |v(x, D_x K_x y)|^p d\sigma_y dx, \end{aligned}$$

where $\tilde{v}(x, y) = v(x, D_x y)$ for $x \in \Omega$ and $y \in Y$, ensure (8.5). \square

The results in Lemma 8.3 ensure the strong convergence of coefficients in equations defined on oscillating boundaries and are used in the derivation of macroscopic problems for microscopic equations defined on surfaces of locally-periodic microstructures.

9. Homogenization of a model for a signaling process in a tissue with locally-periodic distribution of cells. In this section we apply the methods of the l-p unfolding operator and l-t-s convergence on oscillating surfaces to derive macroscopic equations for microscopic models posed in domains with locally-periodic perforations. We consider a generalization of the model for an intercellular signaling process presented in [36] to tissues with locally-periodic microstructures. As examples for tissues with space-dependent changes in the size and shape of cells we consider epithelial and plant cell tissues, see Fig. 3. As an example of a tissue which has a

FIG. 5. *Images of laminar cleavage planes in longitudinal-radial (A) and circumferential-radial (B) tissue sections from basal and apical measurement sites in anterior LV free wall. Reproduced from Costa, Takayana, McCulloch, Covell, 1999[23]. Cardiac muscle fiber orientations vary continuously through the left ventricular wall from a negative angle at the epicardium to positive values toward the endocardium. Reproduced from McCulloch [41].*

plywood-like structure we consider the cardiac muscle tissue of the left ventricular wall, see Fig. 5.

The microstructure of cardiac muscle is described in the same way as a plywood-like structure considered in the introduction, where $D(x) = R^{-1}(\gamma(x_3))$ and the rotation matrix R is as defined in the introduction. Additionally we assume that the radius of fibers may change locally, i.e. $K(x)Y_0 \subset Y$, with $K(x) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \rho(x) & 0 \\ 0 & 0 & \rho(x) \end{pmatrix}$, $Y_0 = \{(y_1, y_2, y_3) \in Y : |y_2 - 1/2|^2 + |y_3 - 1/2|^2 < a^2\}$, $0 < a < 1/2$, and $\rho \in C^1(\bar{\Omega})$ with $0 < \rho_1 \leq \rho(x)a < 1/2$ for all $x \in \bar{\Omega}$. Then, for the plywood-like structure $D_{x_n^\varepsilon} = R^{-1}(\gamma(x_{n,3}^\varepsilon))$, $\tilde{Y}_{K_x}^* = Y \setminus K(x)\bar{Y}_0$, $Y_{x,K}^* = R^{-1}(\gamma(x_3))\tilde{Y}_{K_x}^*$, and the characteristic function of fibers is given by

$$\chi_{\Omega_f^\varepsilon}(x) = \chi_\Omega(x) \sum_{n=1}^{N_\varepsilon} \tilde{\eta}(x_n^\varepsilon, R(\gamma(x_{n,3}^\varepsilon))x/\varepsilon)\chi_{\Omega_n^\varepsilon},$$

where

$$\tilde{\eta}(x, y) = \begin{cases} 1 & \text{for } |\hat{K}(x)^{-1}\hat{y} - (1/2, 1/2)| \leq a, \\ 0 & \text{elsewhere,} \end{cases}$$

and extended \hat{Y} -periodically to the whole of \mathbb{R}^3 . Here $\hat{y} = (y_2, y_3)$, $\hat{Y} = [0, 1]^2$, and $\hat{K}(x) = \rho(x)\mathbf{I}_2$, where \mathbf{I}_2 denotes the identity matrix in $\mathbb{R}^{2 \times 2}$

In the case of an epithelial tissue consider $Y_x = D(x)Y$, with e.g. $D(x) = \begin{pmatrix} \mathbf{I}_2 & \mathbf{0} \\ \mathbf{0}^T & \kappa(x) \end{pmatrix}$, where $\kappa \in C^1(\bar{\Omega})$ and $0 < \kappa_1 \leq \kappa(x) < 1$ for all $x \in \Omega$ defines a compression of cells in x_3 -direction. The changes in the size and shape of cells can be defined by the boundaries of the microstructure $\Gamma_x = S(x)\Gamma \subset Y_x = D_x Y$ for all $x \in \bar{\Omega}$ and $S \in \text{Lip}(\bar{\Omega}; \mathbb{R}^{3 \times 3})$. Then, in the definition of the intercellular space $\Omega_{\varepsilon,K}^*$ we have $Y_{x,K}^* = D(x)\tilde{Y}_{K_x}^* = D(x)(Y \setminus K(x)\bar{Y}_0)$, where $K(x) = D(x)^{-1}S(x)$.

We define the intercellular space in a tissues as

$$\Omega_{\varepsilon,K}^* = \text{Int} \left(\bigcup_{n=1}^{N_\varepsilon} \bar{\Omega}_{n,K}^\varepsilon \right) \cap \Omega, \quad \text{where } \Omega_{n,K}^{*,\varepsilon} = \Omega_n^\varepsilon \setminus \bigcup_{\xi \in \hat{\Xi}_n^\varepsilon} D_{x_n^\varepsilon}(K_{x_n^\varepsilon}\bar{Y}_0 + \xi).$$

We shall use the notation $\hat{\Omega}_{\varepsilon,K}^* = \bigcup_{n=1}^{N_\varepsilon} \bigcup_{\xi \in \hat{\Xi}_n^\varepsilon} \varepsilon D_{x_n^\varepsilon}(\tilde{Y}_{K_{x_n^\varepsilon}}^* + \xi)$ and $\Lambda_{\varepsilon,K}^* = \Omega_{\varepsilon,K}^* \setminus \hat{\Omega}_{\varepsilon,K}^*$.

In the model for a signaling process in a cell tissue we consider diffusion of signaling molecules l^ε in the inter-cellular space and their interactions with free and bound receptors r_f^ε and r_b^ε located on cell surfaces. The microscopic model reads

$$(9.1) \quad \begin{aligned} \partial_t l^\varepsilon - \operatorname{div}(A^\varepsilon(x) \nabla l^\varepsilon) &= F^\varepsilon(x, l^\varepsilon) - d_l^\varepsilon(x) l^\varepsilon && \text{in } (0, T) \times \Omega_{\varepsilon,K}^*, \\ A^\varepsilon(x) \nabla l^\varepsilon \cdot \mathbf{n} &= \varepsilon [\beta^\varepsilon(x) r_b^\varepsilon - \alpha^\varepsilon(x) l^\varepsilon r_f^\varepsilon] && \text{on } (0, T) \times \Gamma^\varepsilon, \\ A^\varepsilon(x) \nabla l^\varepsilon \cdot \mathbf{n} &= 0 && \text{on } (0, T) \times (\partial\Omega \cap \partial\Omega_{\varepsilon,K}^*), \\ l^\varepsilon(0, x) &= l_0(x) && \text{in } \Omega_{\varepsilon,K}^*, \end{aligned}$$

where the dynamics in the concentrations of free and bound receptors on cell surfaces are determined by two ordinary differential equations

$$(9.2) \quad \begin{aligned} \partial_t r_f^\varepsilon &= p^\varepsilon(x, r_f^\varepsilon) - \alpha^\varepsilon(x) l^\varepsilon r_f^\varepsilon + \beta^\varepsilon(x) r_b^\varepsilon - d_f^\varepsilon(x) r_f^\varepsilon && \text{on } (0, T) \times \Gamma^\varepsilon, \\ \partial_t r_b^\varepsilon &= \alpha^\varepsilon(x) l^\varepsilon r_f^\varepsilon - \beta^\varepsilon(x) r_b^\varepsilon - d_b^\varepsilon(x) r_b^\varepsilon && \text{on } (0, T) \times \Gamma^\varepsilon, \\ r_f^\varepsilon(0, x) &= r_{f0}^\varepsilon(x), \quad r_b^\varepsilon(0, x) = r_{b0}^\varepsilon(x) && \text{on } \Gamma^\varepsilon. \end{aligned}$$

The coefficients A^ε , α^ε , β^ε , d_j^ε and the functions $F^\varepsilon(\cdot, \xi)$, $p^\varepsilon(\cdot, \xi)$, r_{i0}^ε are defined as

$$\begin{aligned} A^\varepsilon(x) &= \mathcal{L}_0^\varepsilon(A(x, y)), & F^\varepsilon(x, \xi) &= \mathcal{L}_0^\varepsilon(F(x, y, \xi)), & p^\varepsilon(x, \xi) &= \mathcal{L}_0^\varepsilon(p(x, y, \xi)), \\ \alpha^\varepsilon(x) &= \mathcal{L}_0^\varepsilon(\alpha(x, y)), & \beta^\varepsilon(x) &= \mathcal{L}_0^\varepsilon(\beta(x, y)), & d_j^\varepsilon(x) &= \mathcal{L}_0^\varepsilon(d_j(x, y)), \\ r_{i0}^\varepsilon(x) &= \mathcal{L}^\varepsilon(r_{i0}(x, y)), & & & j &= l, f, b, \quad i = f, b, \end{aligned}$$

for $x \in \Omega$, $y \in Y_x$ and $\xi \in \mathbb{R}$, where $A(x, \cdot)$, $\alpha(x, \cdot)$, $\beta(x, \cdot)$, $d_j(x, \cdot)$, $p(x, \cdot, \xi)$, $F(x, \cdot, \xi)$, and $r_{i0}(x, \cdot)$ are Y_x -periodic functions. We assume also that $\alpha^\varepsilon(x) = 0$ and $\beta^\varepsilon(x) = 0$ for $x \in \Lambda^\varepsilon$. The last assumption is not restrictive, since the domain Λ^ε is very small compared to the size of the whole domain Ω and $|\Lambda^\varepsilon| \leq C\varepsilon^{1-r} \rightarrow 0$ as $\varepsilon \rightarrow 0$ for $0 \leq r < 1$.

Here, $A^\varepsilon : \Omega_T \rightarrow \mathbb{R}$ denotes the diffusion coefficient for signaling molecules (ligands), $F^\varepsilon : \Omega_T \times \mathbb{R} \rightarrow \mathbb{R}$ models the production of new ligands, $p^\varepsilon : \Omega_T \rightarrow \mathbb{R}$ describes the production of new free receptors, $d_j^\varepsilon : \Omega_T \rightarrow \mathbb{R}$, $j = l, f, b$, denote the rates of decay of ligands, free and bound receptors, respectively, $\beta^\varepsilon : \Omega_T \rightarrow \mathbb{R}$ is the rate of dissociation of bound receptors, $\alpha^\varepsilon : \Omega_T \rightarrow \mathbb{R}$ is the rate of binding of ligands to free receptors.

ASSUMPTION 9.1.

- $A \in C(\bar{\Omega}; L_{\text{per}}^\infty(Y_x))$ is symmetric with $(A(x, y)\xi, \xi) \geq d_0|\xi|^2$ for $d_0 > 0$, $\xi \in \mathbb{R}^d$, $x \in \Omega$ and a.a. $y \in Y_x$.
- $F(\cdot, \cdot, \xi) \in C(\bar{\Omega}; L_{\text{per}}^\infty(Y_x))$ is Lipschitz continuous in ξ uniformly in (x, y) and $F(x, y, \xi) \geq 0$ for $\xi \geq 0$, a.a. $x \in \Omega$ and $y \in Y_x$.
- $p(\cdot, \cdot, \xi) \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ is Lipschitz continuous in ξ uniformly in (x, y) and nonnegative for nonnegative ξ .
- Coefficients $\alpha, \beta, d_j \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ are nonnegative, $j = l, f, b$.
- Initial conditions $l_0 \in H^1(\Omega)$, $r_{j0} \in C(\bar{\Omega}; C_{\text{per}}(Y_x))$ are nonnegative, $j = f, b$.

Notice that these assumptions are satisfied by the physical processes described by our model, since for most signaling processes in biological tissues we have that $A = \text{const}$, $F(x, y, \xi) = \mu_1 \xi / (\mu_2 + \mu_3 \xi)$, and $p(x, y, \xi) = \kappa_1 \xi / (\kappa_2 + \kappa_3 \xi)$ with some nonnegative constants μ_i and κ_i , for $i = 1, 2, 3$, and the coefficients α, β , and d_j , with $j = l, f, b$, can be chosen as constant or as some smooth functions.

We shall use the following notations $\hat{\Gamma}_T^\varepsilon = (0, T) \times \hat{\Gamma}^\varepsilon$, $\Gamma_T^\varepsilon = (0, T) \times \Gamma^\varepsilon$, $\Omega_T = (0, T) \times \Omega$, $\Gamma_T = (0, T) \times \Gamma$, and $\Gamma_{x,T} = (0, T) \times \Gamma_x$.

DEFINITION 9.1. *A weak solution of the problem (9.1)–(9.2) are functions $(l^\varepsilon, r_f^\varepsilon, r_b^\varepsilon)$ such that $l^\varepsilon \in L^2(0, T; H^1(\Omega_{\varepsilon, K}^*)) \cap H^1(0, T; L^2(\Omega_{\varepsilon, K}^*))$, $r_j^\varepsilon \in H^1(0, T; L^2(\Gamma^\varepsilon)) \cap L^\infty(\Gamma_T^\varepsilon)$, for $j = f, b$, satisfying the equation (9.1) in the weak form*

$$(9.3) \quad \begin{aligned} \langle \partial_t l^\varepsilon, \phi \rangle_{\Omega_{\varepsilon, K}^*, T} + \langle A^\varepsilon(x) \nabla l^\varepsilon, \nabla \phi \rangle_{\Omega_{\varepsilon, K}^*, T} &= \langle F^\varepsilon(x, l^\varepsilon) - d_l^\varepsilon(x) l^\varepsilon, \phi \rangle_{\Omega_{\varepsilon, K}^*, T} \\ &+ \varepsilon \langle \beta^\varepsilon(x) r_b^\varepsilon - \alpha^\varepsilon(x) l^\varepsilon r_f^\varepsilon, \phi \rangle_{\hat{\Gamma}^\varepsilon, T}, \end{aligned}$$

for all $\phi \in L^2(0, T; H^1(\Omega_{\varepsilon, K}^*))$, the equations (9.2) are satisfied a.e. on Γ_T^ε , and $l^\varepsilon(t, \cdot) \rightarrow l_0(\cdot)$ in $L^2(\Omega_{\varepsilon, K}^*)$, $r_j^\varepsilon(t, \cdot) \rightarrow r_{j0}^\varepsilon(\cdot)$ in $L^2(\Gamma^\varepsilon)$ as $t \rightarrow 0$.

Here for $v, w \in L^2((0, \sigma) \times A)$ we use the notation $\langle v, w \rangle_{A, \sigma} = \int_0^\sigma \int_A v w \, dx dt$.

In a similar way as in [16, 36] we obtain the existence and uniqueness results and *a priori* estimates for weak solutions of the microscopic model (9.1)–(9.2).

LEMMA 9.2. *Under Assumption 9.1 there exists a unique non-negative weak solution of the microscopic model (9.1)–(9.2) satisfying a priori estimates*

$$(9.4) \quad \begin{aligned} \|l^\varepsilon\|_{L^\infty(0, T; L^2(\Omega_{\varepsilon, K}^*))} + \|\nabla l^\varepsilon\|_{L^\infty(0, T; L^2(\Omega_{\varepsilon, K}^*))} + \|\partial_t l^\varepsilon\|_{L^2((0, T) \times \Omega_{\varepsilon, K}^*)} &\leq C, \\ \varepsilon^{1/2} \|l^\varepsilon\|_{L^2(\hat{\Gamma}^\varepsilon)} + \|r_j^\varepsilon\|_{L^\infty(0, T; L^\infty(\Gamma^\varepsilon))} + \varepsilon^{1/2} \|\partial_t r_j^\varepsilon\|_{L^2(\Gamma_T^\varepsilon)} &\leq C, \end{aligned}$$

with $j = f, b$, where the constant C is independent of ε . Additionally, we have that

$$(9.5) \quad \|(l^\varepsilon - Me^{Bt})^+\|_{L^\infty(0, T; L^2(\Omega_{\varepsilon, K}^*))} + \|\nabla(l^\varepsilon - Me^{Bt})^+\|_{L^2((0, T) \times \Omega_{\varepsilon, K}^*)} \leq C\varepsilon^{1/2},$$

where $v^+ = \max\{0, v\}$, $M \geq \sup_\Omega l_0(x)$, $B = B(F, \beta, p)$, and C is independent of ε .

Proof. [Proof Sketch] To prove the existence of a solution of the microscopic model we show the existence of a fix point of an operator \mathcal{B} defined on $L^2(0, T; H^\varsigma(\Omega_{\varepsilon, K}^*))$, with $1/2 < \varsigma < 1$, by $l_n^\varepsilon = \mathcal{B}(l_{n-1}^\varepsilon)$ given as a solution of (9.1)–(9.2) with l_{n-1}^ε in the equations (9.2) and in the nonlinear function $F^\varepsilon(x, l^\varepsilon)$ instead of l_n^ε . For a given non-negative $l_{n-1}^\varepsilon \in L^2(0, T; H^\varsigma(\Omega_{\varepsilon, K}^*))$ there exists a non-negative solution $(r_{f,n}^\varepsilon, r_{b,n}^\varepsilon)$ of (9.2). Then, the non-negativity of solutions, the equality

$$\partial_t(r_{f,n}^\varepsilon + r_{b,n}^\varepsilon) = p^\varepsilon(x, r_{f,n}^\varepsilon) - d_b^\varepsilon(x) r_{b,n}^\varepsilon - d_f^\varepsilon(x) r_{f,n}^\varepsilon,$$

and the Lipschitz continuity of p ensure the boundedness of $r_{f,n}^\varepsilon$ and $r_{b,n}^\varepsilon$. Considering $l_n^{\varepsilon, -} = \min\{0, l_n^\varepsilon\}$ as a test function in (9.3) and using the non-negativity of $r_{f,n}^\varepsilon$, $r_{b,n}^\varepsilon$ and the initial data we obtain the non-negativity of l_n^ε . Applying Galerkin's method and using a priori estimates similar to these in (9.4) we obtain the existence of a weak non-negative solution $l_n^\varepsilon \in H^1(0, T; L^2(\Omega_{\varepsilon, K}^*)) \cap L^2(0, T; H^1(\Omega_{\varepsilon, K}^*))$. The compactness of the embedding $H^1(0, T; L^2(\Omega_{\varepsilon, K}^*)) \cap L^2(0, T; H^1(\Omega_{\varepsilon, K}^*)) \subset L^2(0, T; H^\varsigma(\Omega_{\varepsilon, K}^*))$ and Schauder's theorem imply the existence of a fixed point l^ε of \mathcal{B} . Notice that the strong convergence of l_n^ε in $L^2(\Gamma_\varepsilon)$, as $n \rightarrow \infty$, implies the strong convergence of $r_{j,n}^\varepsilon$, $j = f, b$. Taking l_n^ε and $\partial_t l_n^\varepsilon$ as test functions in (9.3) and using the trace estimate (8.1) we obtain a priori estimates for l_n^ε . Testing (9.2) by $\partial_t r_{f,n}^\varepsilon$ and $\partial_t r_{b,n}^\varepsilon$, respectively, yields the estimates for the time derivatives. Then, using the lower semicontinuity of the norm we obtain the *a priori* estimates (9.4) for l^ε , r_f^ε and r_b^ε .

Especially for the derivation of *a priori* estimates for $\partial_t l^\varepsilon$ we consider

$$\begin{aligned} \varepsilon \int_{\Gamma^\varepsilon} (\beta^\varepsilon r_b^\varepsilon - \alpha^\varepsilon r_f^\varepsilon l^\varepsilon) \partial_t l^\varepsilon \, d\sigma_x &= \varepsilon \frac{d}{dt} \int_{\Gamma^\varepsilon} \beta^\varepsilon r_b^\varepsilon l^\varepsilon \, d\sigma_x - \varepsilon \int_{\Gamma^\varepsilon} \beta^\varepsilon \partial_t r_b^\varepsilon l^\varepsilon \, d\sigma_x \\ &- \frac{\varepsilon}{2} \frac{d}{dt} \int_{\Gamma^\varepsilon} \alpha^\varepsilon r_f^\varepsilon |l^\varepsilon|^2 \, d\sigma_x + \frac{\varepsilon}{2} \int_{\Gamma^\varepsilon} \alpha^\varepsilon \partial_t r_f^\varepsilon |l^\varepsilon|^2 \, d\sigma_x. \end{aligned}$$

Using the equation for $\partial_t r_f^\varepsilon$, the last integral can be rewritten as

$$\frac{\varepsilon}{2} \int_{\Gamma^\varepsilon} \alpha^\varepsilon (p^\varepsilon(x, r_f^\varepsilon) - \alpha^\varepsilon l^\varepsilon r_f^\varepsilon + \beta^\varepsilon r_b^\varepsilon - d_f^\varepsilon r_f^\varepsilon) |l^\varepsilon|^2 d\sigma_x.$$

Applying the trace estimate (8.1) and using the assumptions on α^ε and β^ε , along with the non-negativity of l^ε and r_j^ε , the boundedness of r_j^ε , uniform in ε , and the estimate $\varepsilon \|\partial_t r_b^\varepsilon\|_{L^2(\Gamma_T^\varepsilon)}^2 \leq C$, we obtain

$$\begin{aligned} \varepsilon \int_0^\tau \int_{\Gamma^\varepsilon} (\beta^\varepsilon r_b^\varepsilon - \alpha^\varepsilon r_f^\varepsilon l^\varepsilon) \partial_t l^\varepsilon d\sigma_x dt &\leq C_1 [\|l^\varepsilon(\tau)\|_{L^2(\Omega_{\varepsilon,K}^*)}^2 + \varepsilon^2 \|\nabla l^\varepsilon(\tau)\|_{L^2(\Omega_{\varepsilon,K}^*)}^2] \\ &\quad + C_2 [\|l^\varepsilon\|_{L^2((0,\tau) \times \Omega_{\varepsilon,K}^*)}^2 + \varepsilon^2 \|\nabla l^\varepsilon\|_{L^2((0,\tau) \times \Omega_{\varepsilon,K}^*)}^2] + C_3 \end{aligned}$$

for $\tau \in (0, T]$. Standard arguments pertaining to the difference of two solutions $l_1^\varepsilon - l_2^\varepsilon$, $r_{j,1}^\varepsilon - r_{j,2}^\varepsilon$, with $j = f, b$, imply the uniqueness of a weak solution of the microscopic model (9.1)–(9.2). In particular, the non-negativity of α^ε , r_f^ε and l^ε along with the boundedness of r_f^ε ensures

$$(9.6) \quad \partial_t \|r_{f,1}^\varepsilon - r_{f,2}^\varepsilon\|_{L^2(\Gamma^\varepsilon)}^2 \leq C \left(\sum_{j=f,b} \|r_{j,1}^\varepsilon - r_{j,2}^\varepsilon\|_{L^2(\Gamma^\varepsilon)}^2 + \|l_1^\varepsilon - l_2^\varepsilon\|_{L^2(\hat{\Gamma}^\varepsilon)}^2 \right).$$

Testing the difference of the equations for $r_{f,1}^\varepsilon + r_{b,1}^\varepsilon$ and $r_{f,2}^\varepsilon + r_{b,2}^\varepsilon$ by $r_{f,1}^\varepsilon + r_{b,1}^\varepsilon - r_{f,2}^\varepsilon - r_{b,2}^\varepsilon$ yields

$$(9.7) \quad \|r_{b,1}^\varepsilon(\tau) - r_{b,2}^\varepsilon(\tau)\|_{L^2(\Gamma^\varepsilon)}^2 \leq C \int_0^\tau \sum_{j=f,b} \|r_{j,1}^\varepsilon - r_{j,2}^\varepsilon\|_{L^2(\Gamma^\varepsilon)}^2 + \|l_1^\varepsilon - l_2^\varepsilon\|_{L^2(\hat{\Gamma}^\varepsilon)}^2 dt.$$

Applying the Gronwall Lemma yields the estimate for $\|r_{j,1}^\varepsilon(\tau) - r_{j,2}^\varepsilon(\tau)\|_{L^2(\Gamma^\varepsilon)}^2$, with $\tau \in (0, T]$ and $j = f, b$, in terms of $\|l_1^\varepsilon - l_2^\varepsilon\|_{L^2(\hat{\Gamma}_T^\varepsilon)}^2$. Taking $(l^\varepsilon - S)^+$ as a test function in (9.3) and using the boundedness of r_j^ε we obtain

$$\|(l^\varepsilon - S)^+\|_{L^\infty(0,T;L^2(\Omega_{\varepsilon,K}^*))} + \|\nabla(l^\varepsilon - S)^+\|_{L^2((0,T) \times \Omega_{\varepsilon,K}^*)} \leq 2S \left(\int_0^T |\Omega_{\varepsilon,K}^{*,S}(t)| dt \right)^{\frac{1}{2}},$$

where $S \geq \max\{\sup_\Omega l_0(x), \sup_{\Omega \times Y_x} |\beta(x, y)|, \sup_{\Omega \times Y_x} |\alpha(x, y)|, \|r_j^\varepsilon\|_{L^\infty(\Gamma_T^\varepsilon)}\}$ and $\Omega_{\varepsilon,K}^{*,S}(t) = \{x \in \Omega_{\varepsilon,K}^* : l^\varepsilon(t, x) > S\}$ for a.a. $t \in (0, T)$. Then, applying Theorem II.6.1 in [35] yields the boundedness of l^ε for every fixed $\varepsilon > 0$. Considering equation (9.3) for l_1^ε and l_2^ε we obtain the estimate for $\|l_1^\varepsilon - l_2^\varepsilon\|_{L^2(0,T;H^1(\Omega_{\varepsilon,K}^*))}$, with $\tau \in (0, T]$, in terms of $\varepsilon^{1/2} \|r_{j,1}^\varepsilon - r_{j,2}^\varepsilon\|_{L^2(\Gamma_T^\varepsilon)}$, with $j = f, b$. Then, using the estimates for $\|r_{j,1}^\varepsilon(\tau) - r_{j,2}^\varepsilon(\tau)\|_{L^2(\Gamma^\varepsilon)}$, with $\tau \in (0, T]$, in (9.6) and (9.7) yields that $l_1^\varepsilon = l_2^\varepsilon$ a.e. in $(0, T) \times \Omega_{\varepsilon,K}^*$ and hence $r_{j,1}^\varepsilon = r_{j,2}^\varepsilon$ a.e. in Γ_T^ε , where $j = f, b$.

To show (9.5), we consider $(l^\varepsilon - Me^{Bt})^+$ as a test function in (9.3). Using the boundedness of r_j^ε , uniform in ε , and the trace estimate (8.1) we obtain for $\tau \in (0, T)$

$$\begin{aligned} \|(l^\varepsilon(\tau) - Me^{B\tau})^+\|_{L^2(\Omega_{\varepsilon,K}^*)}^2 + \|\nabla(l^\varepsilon - Me^{Bt})^+\|_{L^2((0,\tau) \times \Omega_{\varepsilon,K}^*)}^2 \\ \leq C_1 \|(l^\varepsilon - Me^{Bt})^+\|_{L^2((0,\tau) \times \Omega_{\varepsilon,K}^*)}^2 + C_2 \varepsilon, \end{aligned}$$

where $M \geq \sup_\Omega l_0(x)$, $MB \geq \left(\sup_{\Omega \times Y_x} |F(x, y, 0)| + \mu_\Gamma \sup_{\Omega \times Y_x} \beta(x, y) \|r_b^\varepsilon\|_{L^\infty(\hat{\Gamma}_T^\varepsilon)} \right)$, with μ_Γ as in (8.1). Applying Gronwall's Lemma in the last inequality yields (9.5). \square

Notice, that in the case of a perforated domain where the periodicity and the shape of perforations vary in space, i.e. $K \neq \mathbf{I}$, we can not apply the l-p unfolding operator to functions defined on $\Omega_{\varepsilon,K}^*$ directly. To overcome this problem we consider a local extension of a function from $\hat{\Omega}_{n,K}^\varepsilon$ to $\hat{\Omega}_n^\varepsilon$ and then apply the l-p unfolding operator $\mathcal{T}_\mathcal{L}^\varepsilon$, determined for functions defined on $\hat{\Omega}^\varepsilon$. Applying the assumptions on the microstructure of $\Omega_{\varepsilon,K}^*$ considered here, i.e. $K_x \bar{Y}_0 \subset Y$ or fibrous microstructure, we obtain

LEMMA 9.3. For $x_n^\varepsilon \in \hat{\Omega}_n^\varepsilon$, where $1 \leq n \leq N_\varepsilon$, and $u \in W^{1,p}(Y_{x_n^\varepsilon}^*)$, with $p \in (1, \infty)$, there exists an extension $\hat{u} \in W^{1,p}(Y_{x_n^\varepsilon})$ such that

$$(9.8) \quad \|\hat{u}\|_{L^p(Y_{x_n^\varepsilon})} \leq \mu \|u\|_{L^p(Y_{x_n^\varepsilon}^*)}, \quad \|\nabla \hat{u}\|_{L^p(Y_{x_n^\varepsilon})} \leq \mu \|\nabla u\|_{L^p(Y_{x_n^\varepsilon}^*)},$$

where μ depends on Y, Y_0, D and K and is independent of ε and n .

For $u \in W^{1,p}(\Omega_{\varepsilon,K}^*)$ we have an extension $\hat{u} \in W^{1,p}(\hat{\Omega}_\varepsilon)$ from $\hat{\Omega}_{\varepsilon,K}^*$ to $\hat{\Omega}^\varepsilon$ such that

$$(9.9) \quad \|\hat{u}\|_{L^p(\hat{\Omega}^\varepsilon)} \leq \mu \|u\|_{L^p(\hat{\Omega}_{\varepsilon,K}^*)}, \quad \|\nabla \hat{u}\|_{L^p(\hat{\Omega}^\varepsilon)} \leq \mu \|\nabla u\|_{L^p(\hat{\Omega}_{\varepsilon,K}^*)},$$

where μ depends on Y, Y_0, D and K and is independent of ε .

Proof. [Sketch of the Proof] The proof follows the same lines as in the periodic case, see e.g. [15, 19]. The only difference is that the extension depends on the Lipschitz continuity of K and D and the uniform boundedness from above and below of $|\det K(x)|$ and $|\det D(x)|$. To show (9.9), we first consider an extension from $D_{x_n^\varepsilon}(\tilde{Y}_{K_{x_n^\varepsilon}^*} + \xi)$ to $D_{x_n^\varepsilon}(Y + \xi)$ satisfying estimates (9.8), where $\xi \in \hat{\Xi}_n^\varepsilon$. Then, scaling by ε and summing up over $\xi \in \hat{\Xi}_n^\varepsilon$ and $n = 1, \dots, N_\varepsilon$ imply the estimates (9.9). \square

REMARK. Notice that the definition of $\Omega_{\varepsilon,K}^*$ implies that there no perforations in $(\Omega_{n,K}^{*,\varepsilon} \setminus \hat{\Omega}_{n,K}^{*,\varepsilon}) \cap \tilde{\Omega}_{\varepsilon/2}$, with $\tilde{\Omega}_{\varepsilon/2} = \{x \in \Omega : \text{dist}(x, \partial\Omega) > 2\varepsilon \max_{x \in \partial\Omega} \text{diam}(D(x)Y)\}$. Also in the case of a plywood-like structure the fibres are orthogonal to the boundaries of Ω_n^ε and near $\partial\Omega_n^\varepsilon$ we need to extend l^ε only in the directions parallel to $\partial\Omega_n^\varepsilon$. Thus, applying Lemma 9.3 we can extend l^ε from $\Omega_{n,K}^{*,\varepsilon}$ into $\hat{\Omega}_n^\varepsilon \cup (\Omega_n^\varepsilon \cap \tilde{\Omega}_{\varepsilon/2})$, for $n = 1, \dots, N_\varepsilon$.

THEOREM 9.4. A sequence of solutions of the microscopic model (9.1)–(9.2) converges to a solution (l, r_f, r_b) with $l \in H^1(\Omega_T)$ and $r_j \in H^1(0, T; L^2(\Omega; L^2(\Gamma_x)))$ of the macroscopic equations

$$(9.10) \quad \begin{aligned} \frac{|Y_{x,K}^*|}{|Y_x|} \partial_t l - \text{div}(\mathcal{A}(x)\nabla l) &= \frac{1}{|Y_x|} \int_{Y_{x,K}^*} F(x, y, l) dy \\ &+ \frac{1}{|Y_x|} \int_{\Gamma_x} (\beta(x, y) r_b - \alpha(x, y) r_f l) d\sigma_y \quad \text{in } \Omega_T, \\ \mathcal{A}(x)\nabla l \cdot \mathbf{n} &= 0 \quad \text{on } \partial\Omega, \\ \partial_t r_f &= p(x, y, r_f) - \alpha(x, y) l r_f + \beta(x, y) r_b - d_f(x, y) r_f \quad \text{for } y \in \Gamma_x, \\ \partial_t r_b &= \alpha(x, y) l r_f - \beta(x, y) r_b - d_b(x, y) r_b \quad \text{for } y \in \Gamma_x, \end{aligned}$$

and for $(t, x) \in \Omega_T$, where $Y_{x,K}^* = D_x(Y \setminus K_x Y_0)$ and the macroscopic diffusion matrix is defined as

$$\mathcal{A}_{ij}(x) = \frac{1}{|Y_x|} \int_{Y_{x,K}^*} [A_{ij}(x, y) + (A(x, y)\nabla_y \omega^j(x, y))_i] dy \quad \text{for } x \in \Omega,$$

for $i, j = 1, \dots, d$, with

$$(9.11) \quad \begin{aligned} \operatorname{div}_y(A(x, y)(\nabla_y \omega^j + e_j)) &= 0 && \text{in } Y_{x, K}^*, \\ A(x, y)(\nabla_y \omega^j + e_j) \cdot \mathbf{n} &= 0 && \text{on } \Gamma_x, \quad \omega^j \text{ } Y_x \text{-periodic.} \end{aligned}$$

We have that $\hat{l}^\varepsilon \rightarrow l$ in $L^2(\Omega_T)$, $\partial_t l^\varepsilon \rightharpoonup \partial_t l$ and $\partial_t r_j^\varepsilon \rightharpoonup \partial_t r_j$ locally-periodic two-scale, $r_j^\varepsilon \rightarrow r_j$ strongly locally-periodic two-scale, $j = f, b$, and

$$\begin{aligned} \nabla l^\varepsilon &\rightharpoonup \nabla l + \nabla_y l_1 && \text{l-t-s,} && l_1 \in L^2(\Omega_T; H_{\text{per}}^1(Y_{x, K}^*)), \\ \lim_{\varepsilon \rightarrow 0} \langle A^\varepsilon \nabla l^\varepsilon, \nabla l^\varepsilon \rangle_{\Omega_{\varepsilon, K}^*, T} &= \langle |Y_x|^{-1} A(x, y) (\nabla l + \nabla_y l_1), \nabla l + \nabla_y l_1 \rangle_{\Omega_T, Y_{x, K}^*}, \end{aligned}$$

where $l_1(t, x, y) = \sum_{j=1}^d \frac{\partial l}{\partial x_j}(t, x) \omega^j(x, y)$. Here $\hat{\phi}$ denotes the extension as in Lemma 9.3

from $(0, T) \times \Omega_{\varepsilon, K}^*$ to $(0, T) \times (\tilde{\Omega}_{\varepsilon/2} \cup \Omega_{\varepsilon, K}^*)$ and then by zero to Ω_T .

Proof. Applying Lemma 9.3 we can extend l^ε from $\Omega_{\varepsilon, K}^*$ into $\hat{\Omega}^\varepsilon \cup \Lambda_{\varepsilon, K}^*$. We shall use the same notations for original functions and their extensions. The *a priori* estimates in Lemma 9.2 imply

$$(9.12) \quad \|l^\varepsilon\|_{L^2(0, T; H^1(\hat{\Omega}^\varepsilon \cup \Lambda_{\varepsilon, K}^*))} + \|\partial_t l^\varepsilon\|_{L^2((0, T) \times (\hat{\Omega}^\varepsilon \cup \Lambda_{\varepsilon, K}^*))} \leq C,$$

where the constant C depends on D and K and is independent of ε . Then the sequences $\{l^\varepsilon\}$, $\{\nabla l^\varepsilon\}$, and $\{\partial_t l^\varepsilon\}$ are defined on $\hat{\Omega}_\varepsilon$ and we can determine $\mathcal{T}_\mathcal{L}^\varepsilon(l^\varepsilon)$, $\mathcal{T}_\mathcal{L}^\varepsilon(\nabla l^\varepsilon)$ and $\partial_t \mathcal{T}_\mathcal{L}^\varepsilon(l^\varepsilon)$. The properties of $\mathcal{T}_\mathcal{L}^\varepsilon$ together with (9.12) ensure

$$\|\mathcal{T}_\mathcal{L}^\varepsilon(l^\varepsilon)\|_{L^2(\Omega_T \times Y)} + \|\mathcal{T}_\mathcal{L}^\varepsilon(\nabla l^\varepsilon)\|_{L^2(\Omega_T \times Y)} + \|\partial_t \mathcal{T}_\mathcal{L}^\varepsilon(l^\varepsilon)\|_{L^2(\Omega_T \times Y)} \leq C.$$

The *a priori* estimates in Lemma 9.2 yield the estimates for the l-p boundary unfolding operator

$$\|\mathcal{T}_\mathcal{L}^{\varepsilon, b}(l^\varepsilon)\|_{L^2(\Omega_T \times \Gamma)} + \|\mathcal{T}_\mathcal{L}^{\varepsilon, b}(r_f^\varepsilon)\|_{H^1(0, T; L^2(\Omega \times \Gamma))} + \|\mathcal{T}_\mathcal{L}^{\varepsilon, b}(r_b^\varepsilon)\|_{H^1(0, T; L^2(\Omega \times \Gamma))} \leq C.$$

Notice that due to the assumptions on $\Omega_{\varepsilon, K}^*$ we have that $\tilde{\Omega}_{\varepsilon/2} \subset \hat{\Omega}^\varepsilon \cup \Lambda_{\varepsilon, K}^*$.

Then, the convergence results in Theorems 4.2, 4.4, 4.5, and 4.6 imply that there exist subsequences (denoted again by l^ε , r_f^ε , r_b^ε) and the functions $l \in L^2(0, T; H^1(\Omega)) \cap H^1(0, T; L^2(\Omega))$, $l_1 \in L^2(\Omega_T; H_{\text{per}}^1(Y_x))$, and $r_j \in H^1(0, T; L^2(\Omega; L^2(\Gamma_x)))$ such that

$$(9.13) \quad \begin{aligned} \mathcal{T}_\mathcal{L}^\varepsilon(l^\varepsilon) &\rightharpoonup l && \text{weakly in } L^2(\Omega_T; H^1(Y)), \\ \mathcal{T}_\mathcal{L}^\varepsilon(l^\varepsilon) &\rightarrow l && \text{strongly in } L^2(0, T; L_{\text{loc}}^2(\Omega; H^1(Y))), \\ \partial_t \mathcal{T}_\mathcal{L}^\varepsilon(l^\varepsilon) &\rightharpoonup \partial_t l && \text{weakly in } L^2(\Omega_T \times Y), \\ \mathcal{T}_\mathcal{L}^\varepsilon(\nabla l^\varepsilon) &\rightharpoonup \nabla l + D_x^{-T} \nabla_{\tilde{y}} l_1(\cdot, D_x \cdot) && \text{weakly in } L^2(\Omega_T \times Y), \\ \mathcal{T}_\mathcal{L}^{b, \varepsilon}(l^\varepsilon) &\rightharpoonup l && \text{weakly in } L^2(\Omega_T \times \Gamma), \\ \mathcal{T}_\mathcal{L}^{b, \varepsilon}(l^\varepsilon) &\rightarrow l && \text{strongly in } L^2(0, T; L_{\text{loc}}^2(\Omega; L^2(\Gamma))), \\ r_j^\varepsilon &\rightharpoonup r_j, \quad \partial_t r_j^\varepsilon \rightharpoonup \partial_t r_j && \text{l-t-s, } r_j, \partial_t r_j \in L^2(\Omega_T; L^2(\Gamma_x)), \\ \mathcal{T}_\mathcal{L}^{b, \varepsilon}(r_j^\varepsilon) &\rightharpoonup r_j(\cdot, D_x K_x \cdot) && \text{weakly in } L^2(\Omega_T \times \Gamma), \\ \partial_t \mathcal{T}_\mathcal{L}^{b, \varepsilon}(r_j^\varepsilon) &\rightharpoonup \partial_t r_j(\cdot, D_x K_x \cdot) && \text{weakly in } L^2(\Omega_T \times \Gamma), \quad j = f, b. \end{aligned}$$

Notice that for l^ε we have *a priori* estimates only in $L^2(0, T; H^1(\hat{\Omega}^\varepsilon \cup \Lambda_{\varepsilon, K}^*))$ and not in $L^2(0, T; H^1(\Omega))$ and can not apply the convergence results in Theorem 4.2 directly. However using $\|l^\varepsilon\|_{L^2(0, T; H^1(\tilde{\Omega}_{\varepsilon/2}))} + \|\partial_t l^\varepsilon\|_{L^2((0, T) \times \tilde{\Omega}_{\varepsilon/2})} \leq C$, ensured by (9.12), applying Lemmas 7.2 and 7.3 to $\mathcal{Q}_{\mathcal{L}}^\varepsilon(l^\varepsilon)$ and $\mathcal{R}_{\mathcal{L}}^\varepsilon(l^\varepsilon)$, respectively, and considering the proof of Theorem 4.4 we obtain the convergences for $\mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)$, $\partial_t \mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)$, and $\mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon)$ in (9.13). Lemma 5.4 implies that $\nabla l^\varepsilon \rightharpoonup \nabla l + \nabla_y l_1$ l-t-s and $\partial_t l^\varepsilon \rightharpoonup \partial_t l$ l-t-s. The local strong convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)$ together with the estimate $\|(l^\varepsilon - Me^{Bt})^+\|_{L^2((0, T) \times \Omega_{\varepsilon, K}^*)} \leq C\varepsilon^{1/2}$, shown in Lemma 9.4, yields the strong convergence of \hat{l}^ε in $L^2(\Omega_T)$.

To derive macroscopic equations for l^ε we consider $\psi^\varepsilon(x) = \psi_1(x) + \varepsilon \mathcal{L}_\rho^\varepsilon(\psi_2)(x)$ with $\psi_1 \in C^1(\bar{\Omega})$ and $\psi_2 \in C_0^1(\Omega; C_{\text{per}}^1(Y_x))$ as a test function in (9.3). Applying the l-p unfolding operator and the l-p boundary unfolding operator implies

$$\begin{aligned} & \frac{1}{|Y|} \left[\langle \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon, K}^*}^\varepsilon) \partial_t \mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon), \mathcal{T}_{\mathcal{L}}^\varepsilon(\psi^\varepsilon) \rangle_{\Omega_T \times Y} + \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon, K}^*}^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon), \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla \psi^\varepsilon) \rangle_{\Omega_T \times Y} \right] \\ &= |Y|^{-1} \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon, K}^*}^\varepsilon) \hat{F}^\varepsilon(x, \tilde{y}, \mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)), \mathcal{T}_{\mathcal{L}}^\varepsilon(\psi^\varepsilon) \rangle_{\Omega_T \times Y} \\ &+ \left\langle \sum_{n=1}^{N_\varepsilon} \frac{\sqrt{g_{x_n^\varepsilon}}}{\sqrt{g}|Y_{x_n^\varepsilon}|} [\mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\beta^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(r_b^\varepsilon) - \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\alpha^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(l^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(r_f^\varepsilon)] \chi_{\Omega_n^\varepsilon}, \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\psi^\varepsilon) \right\rangle_{\Omega_T \times \Gamma} \\ &\quad - \langle \partial_t l^\varepsilon, \psi^\varepsilon \rangle_{\Lambda_{\varepsilon, T}^*} - \langle A^\varepsilon(x) \nabla l^\varepsilon, \nabla \psi^\varepsilon \rangle_{\Lambda_{\varepsilon, T}^*} + \langle F^\varepsilon(x, l^\varepsilon), \psi^\varepsilon \rangle_{\Lambda_{\varepsilon, T}^*}, \end{aligned}$$

where $\hat{F}^\varepsilon(x, \tilde{y}, l^\varepsilon) = \sum_{n=1}^{N_\varepsilon} F(x_n^\varepsilon, D_{x_n^\varepsilon} \tilde{y}, \mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)) \chi_{\Omega_n^\varepsilon}(x)$ for $\tilde{y} \in Y$, $x \in \Omega$ and $\chi_{\Omega_{\varepsilon, K}^*}^\varepsilon = \mathcal{L}_0^\varepsilon(\chi_{Y_{x, K}^*})$. Here $\chi_{Y_{x, K}^*}$ is the characteristic function of $Y_{x, K}^* = D_x(Y \setminus K_x Y_0)$, extended Y_x -periodically to \mathbb{R}^d . We notice that $\hat{F}^\varepsilon(x, \tilde{y}, \xi) = \mathcal{T}_{\mathcal{L}}^\varepsilon(\mathcal{L}_0^\varepsilon(F(x, y, \xi)))$.

Applying Lemma 5.3 yields $\mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon, K}^*}^\varepsilon)(x, \tilde{y}) \rightarrow \chi_{Y_{x, K}^*}(x, D_x \tilde{y})$, $\mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon)(x, \tilde{y}) \rightarrow A(x, D_x \tilde{y})$, and $\hat{F}^\varepsilon(x, \tilde{y}, l) \rightarrow F(x, D_x \tilde{y}, l)$ in $L^p(\Omega_T \times Y)$, for $1 < p < \infty$, as $\varepsilon \rightarrow 0$. Lemma 8.3 ensures $\mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\phi^\varepsilon)(x, \hat{y}) \rightarrow \phi(x, D_x K_x \hat{y})$ in $L^p(\Omega_T \times \Gamma)$ as $\varepsilon \rightarrow 0$, where $\phi^\varepsilon(x) = \beta^\varepsilon(x)$, $\alpha^\varepsilon(x)$, or $d_j^\varepsilon(x)$ and $\phi(x, y) = \alpha(x, y)$, $\beta(x, y)$, or $d_j(x, y)$, with $j = f, b$, respectively.

For an arbitrary $\delta > 0$ we consider $\Omega^\delta = \{x \in \Omega : \text{dist}(x, \partial\Omega) > \delta\}$ and rewrite the boundary integral in the form

$$\begin{aligned} & \left\langle \sum_{n=1}^{N_\varepsilon} \frac{\sqrt{g_{x_n^\varepsilon}}}{\sqrt{g}|Y_{x_n^\varepsilon}|} \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\alpha^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(l^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(r_f^\varepsilon) \chi_{\Omega_n^\varepsilon}, \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\psi^\varepsilon) \right\rangle_{\Omega^\delta \times \Gamma_T} \\ &+ \left\langle \sum_{n=1}^{N_\varepsilon} \frac{\sqrt{g_{x_n^\varepsilon}}}{\sqrt{g}|Y_{x_n^\varepsilon}|} \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\alpha^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(l^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(r_f^\varepsilon) \chi_{\Omega_n^\varepsilon}, \mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\psi^\varepsilon) \right\rangle_{(\Omega \setminus \Omega^\delta) \times \Gamma_T} = I_1 + I_2. \end{aligned}$$

Using the *a priori* estimates for l^ε and r_j^ε , the weak convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)$ in $L^2(\Omega_T; H^1(Y))$ and the strong convergence in $L^2(0, T; L_{\text{loc}}^2(\Omega; H^1(Y)))$ we obtain

$$(9.14) \quad \begin{aligned} \lim_{\delta \rightarrow 0} \lim_{\varepsilon \rightarrow 0} I_1 &= \left\langle \frac{\sqrt{g_x}}{\sqrt{g}|Y_x|} \alpha(x, D_x K_x \hat{y}) r_f(x, D_x K_x \hat{y}) l(x), \psi_1(x) \right\rangle_{\Omega_T \times \Gamma}, \\ \lim_{\delta \rightarrow 0} \lim_{\varepsilon \rightarrow 0} I_2 &= 0. \end{aligned}$$

To obtain (9.14) we also used the strong convergence and boundedness of $\mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(\alpha^\varepsilon)$, the weak convergence and boundedness of $\mathcal{T}_{\mathcal{L}}^{b, \varepsilon}(r_f^\varepsilon)$, the regularity of D and K , and the

strong convergence of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\psi^\varepsilon)$. Similar arguments along with the Lipschitz continuity of F and the strong convergence of $\hat{F}^\varepsilon(x, \tilde{y}, l)$ and $\mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\hat{\Omega}_{\varepsilon,K}^*}) = \mathcal{T}_{\mathcal{L}}^\varepsilon(\mathcal{L}_0^\varepsilon(\chi_{Y_{x,K}^*}))$ ensure

$$\langle \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\hat{\Omega}_{\varepsilon,K}^*}) \hat{F}^\varepsilon(x, \tilde{y}, \mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)), \mathcal{T}_{\mathcal{L}}^\varepsilon(\psi^\varepsilon) \rangle_{\Omega_T \times Y} \rightarrow \langle \chi_{Y_{x,K}^*}(x, D_x \tilde{y}) F(x, D_x \tilde{y}, l), \psi_1 \rangle_{\Omega_T \times Y}$$

as $\varepsilon \rightarrow 0$ and $\delta \rightarrow 0$. Using the convergence results (9.13), the strong convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(\psi^\varepsilon)$ and $\mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla \psi^\varepsilon)$ and the fact that $|\Lambda_{\varepsilon,K}^*| \rightarrow 0$ as $\varepsilon \rightarrow 0$, taking the limit as $\varepsilon \rightarrow 0$, and considering the transformation of variables $y = D_x \tilde{y}$ for $\tilde{y} \in Y$ and $y = D_x K_x \hat{y}$ for $\hat{y} \in \Gamma$ we obtain

$$\begin{aligned} & \langle |Y_x|^{-1} l, \psi_1 \rangle_{Y_{x,K}^* \times \Omega_T} + \langle |Y_x|^{-1} A(x, y) (\nabla l + \nabla_y l_1), \nabla \psi_1 + \nabla_y \psi_2 \rangle_{Y_{x,K}^* \times \Omega_T} \\ & + \langle |Y_x|^{-1} [\alpha(x, y) r_f l - \beta(x, y) r_b], \psi_1 \rangle_{\Gamma_x \times \Omega_T} = \langle |Y_x|^{-1} F(x, y, l), \psi_1 \rangle_{Y_{x,K}^* \times \Omega_T}. \end{aligned}$$

Considering $\psi_1(t, x) = 0$ for $(t, x) \in \Omega_T$ we obtain $l_1(t, x, y) = \sum_{j=1}^d \partial_{x_j} l(t, x) \omega^j(x, y)$, where ω^j are solutions of (9.11). Choosing $\psi_2(t, x, y) = 0$ for $x \in \Omega_T$ and $y \in Y_x$ yields the macroscopic equation for l . Applying the l-p boundary unfolding operator to the equations on Γ^ε we obtain

$$\begin{aligned} (9.15) \quad & \partial_t \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon) = \hat{p}^\varepsilon(x, \hat{y}, \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon)) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\alpha^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(l^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon) \\ & + \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\beta^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_b^\varepsilon) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(d_f^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon), \\ & \partial_t \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_b^\varepsilon) = \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\alpha^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(l^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\beta^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_b^\varepsilon) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(d_b^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_b^\varepsilon), \end{aligned}$$

in $\Omega_T \times \Gamma$, where $\hat{p}^\varepsilon(x, \hat{y}, \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon)) = \sum_{n=1}^{N_\varepsilon} p(x_n^\varepsilon, D_{x_n^\varepsilon} K_{x_n^\varepsilon} \hat{y}, \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon)) \chi_{\hat{\Omega}_n^\varepsilon}(x)$ for $\hat{y} \in \Gamma$ and $x \in \Omega$. In order to pass to the limit in the nonlinear function $\hat{p}^\varepsilon(x, \hat{y}, \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon))$ we have to show the strong convergence of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon)$. We consider the difference of the equations for $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_f^{\varepsilon_k})$ and $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_f^{\varepsilon_m})$ and use $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_f^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_f^{\varepsilon_m})$ as a test function. Applying the Lipschitz continuity of p along with the strong convergence of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\alpha^\varepsilon)$, $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\beta^\varepsilon)$, and $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(d_j^\varepsilon)$, and the non-negativity of l^ε and α^ε yields

$$\begin{aligned} & \frac{d}{dt} \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_f^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_f^{\varepsilon_m})\|_{L^2(\Omega \times \Gamma)}^2 \leq C \left[\sum_{j=f,b} \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_j^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_j^{\varepsilon_m})\|_{L^2(\Omega \times \Gamma)}^2 \right. \\ & + \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(l^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(l^{\varepsilon_m})\|_{L^2(\Omega^\delta \times \Gamma)}^2 + \delta^{\frac{1}{2}} \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(l^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(l^{\varepsilon_m})\|_{L^2((\Omega \setminus \Omega_\delta) \times \Gamma)} \\ & \left. + \sigma(\varepsilon_k, \varepsilon_m) \right], \end{aligned}$$

where $\sigma(\varepsilon_k, \varepsilon_m) \rightarrow 0$ as $\varepsilon_k, \varepsilon_m \rightarrow 0$. Considering the sum of the equations for $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_j^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_j^{\varepsilon_m})$, with $j = f, b$, using $\sum_{j=f,b} \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_j^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_j^{\varepsilon_m})$ as a test function, and applying the Lipschitz continuity of p imply

$$\begin{aligned} & \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_b^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_b^{\varepsilon_m})\|_{L^2(\Omega \times \Gamma)}^2 \leq C_1 \int_0^t \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(l^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(l^{\varepsilon_m})\|_{L^2(\Omega_\delta \times \Gamma)}^2 d\tau \\ & + C_2 \int_0^t \sum_{j=f,b} \|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_j^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_j^{\varepsilon_m})\|_{L^2(\Omega \times \Gamma)}^2 d\tau + \sigma(\varepsilon_k, \varepsilon_m) + C_3 \delta^{\frac{1}{2}}. \end{aligned}$$

Using the a priori estimates for l^ε and the local strong convergence of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(l^\varepsilon)$, collecting the estimates from above, and applying the Gronwall inequality we obtain

$$\|\mathcal{T}_{\mathcal{L}}^{b,\varepsilon_k}(r_j^{\varepsilon_k}) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon_m}(r_j^{\varepsilon_m})\|_{L^2(\Omega \times \Gamma)} \leq C(\sigma(\varepsilon_k, \varepsilon_m) + \delta^{\frac{1}{4}}) \quad \text{for } j = f, b,$$

where $\sigma(\varepsilon_k, \varepsilon_m) \rightarrow 0$ as $\varepsilon_k, \varepsilon_m \rightarrow 0$ and $\delta > 0$ is arbitrary. Thus, we conclude that $\{\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_j^\varepsilon)\}$, for $j = f, b$, are Cauchy sequences in $L^2(\Omega_T \times \Gamma)$. Using the strong convergence of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon)$ and the Lipschitz continuity of p we obtain $\hat{p}^\varepsilon(x, \hat{y}, \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon)) \rightharpoonup p(x, D_x K_x \hat{y}, r_f)$ in $L^2(\Omega_T \times \Gamma)$. Then, passing in (9.15) to the limit as $\varepsilon \rightarrow 0$ implies the macroscopic equations (9.10) for r_f and r_b . This concludes the proof of the convergence up to sub-sequences. The strong convergence of $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_j^\varepsilon)$ together with the estimates in Lemma 8.2, the boundedness of r_j^ε , with $j = f, b$, and the regularity of D and K ensure the strong l-t-s convergence of r_j^ε , i.e.

$$\lim_{\varepsilon \rightarrow 0} \varepsilon \|r_j^\varepsilon\|_{L^2(\Gamma_T^\varepsilon)}^2 = \int_{\Omega_T} \frac{1}{|Y_x|} \int_{\Gamma_x} |r_j(t, x, y)|^2 d\sigma_x dx dt, \quad \text{for } j = f, b.$$

The non-negativity of l^ε and r_j^ε and the uniform boundedness of r_j^ε , with $j = f, b$ (see Lemma 9.2) along with the weak convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(r_j^\varepsilon)$ and l^ε ensure the non-negativity of r_j and l and the boundedness of $r_j(t, x, y)$ for a.a. $(t, x) \in \Omega_T$ and $y \in \Gamma_x$. Considering $(l - M_1 e^{M_2 t})^+$ as a test function in the weak formulation of the macroscopic model (9.10) and using the boundedness of r_f and r_b we obtain

$$\|(l - M_1 e^{M_2 t})^+\|_{L^\infty(0, T; L^2(\Omega))} + \|\nabla(l - M_1 e^{M_2 t})^+\|_{L^2(\Omega_T)} \leq 0.$$

Hence, $0 \leq l(t, x) \leq M_1 e^{M_2 T}$ for a.a. $(t, x) \in \Omega_T$, where $M_1 \geq \sup_\Omega l_0(x)$ and $M_1 M_2 \geq (\|F(x, y, 0)\|_{L^\infty(\Omega; L^\infty(Y_x))} + |Y_{x,K}^*|^{-1} \|\beta(x, y)\|_{L^\infty(\Omega; L^\infty(Y_x))} \|r_b\|_{L^\infty(\Omega; L^1(\Gamma_x))})$.

Considering equations for the difference of two solutions of (9.10), taking $l_1 - l_2$, $r_{f,1} - r_{f,2}$, and $r_{b,1} - r_{b,2}$ as test functions in the weak formulation of the macroscopic model, and using the Lipschitz continuity of F and p along with boundedness of r_j and l , we obtain uniqueness of a weak solution of the model (9.10). Thus, we have that the entire sequence of weak solutions $(l^\varepsilon, r_f^\varepsilon, r_b^\varepsilon)$ of the microscopic model (9.1)–(9.2) convergences to the weak solution of the macroscopic model (9.10).

Applying the lower-semicontinuity of a norm, the ellipticity of A , and the strong convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon)$ and $\mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*})$ in $L^p(\Omega_T \times Y)$ for any $p \in (1, +\infty)$, yields

$$\begin{aligned} & \langle |Y_x|^{-1} A(x, y) (\nabla l + \nabla_y l_1), \nabla l + \nabla_y l_1 \rangle_{\Omega_T, Y_{x,K}^*} \\ & \leq \liminf_{\varepsilon \rightarrow 0} |Y|^{-1} \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon), \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon) \rangle_{\Omega_T, Y} \\ & \leq \limsup_{\varepsilon \rightarrow 0} |Y|^{-1} \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon), \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon) \rangle_{\Omega_T, Y} \\ & \leq \limsup_{\varepsilon \rightarrow 0} \langle A^\varepsilon \nabla l^\varepsilon, \nabla l^\varepsilon \rangle_{\Omega_{\varepsilon,K}^*, T} = \limsup_{\varepsilon \rightarrow 0} [I_1 + I_2 + I_3], \end{aligned}$$

where

$$\begin{aligned} I_1 &= |Y|^{-1} \langle \hat{F}^\varepsilon(x, \tilde{y}, \mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)) - \partial_t \mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon), \mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon) \rangle_{\Omega_T, Y}, \\ I_2 &= \int_{\Omega_T \times \Gamma} \sum_{n=1}^{N_\varepsilon} \frac{\sqrt{g_{x_n^\varepsilon}}}{\sqrt{g}|Y_{x_n^\varepsilon}|} \left[\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\beta^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_b^\varepsilon) - \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(\alpha^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(l^\varepsilon) \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_f^\varepsilon) \right] \mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(l^\varepsilon) \chi_{\Omega_n^\varepsilon} d\sigma_y, \\ I_3 &= \langle F^\varepsilon(x, l^\varepsilon) - \partial_t l^\varepsilon, l^\varepsilon \rangle_{\Lambda_{\varepsilon,K}^*, T}. \end{aligned}$$

Using the estimates in Lemma 9.2, together with $0 \leq l^\varepsilon \leq M + (l^\varepsilon - M)^+$ and the definition of $\Lambda_{\varepsilon,K}^*$, we obtain $\lim_{\varepsilon \rightarrow 0} I_3 = 0$.

Considering the strong convergence $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(r_j^\varepsilon)$, with $j = f, b$, and the local strong convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(l^\varepsilon)$ and $\mathcal{T}_{\mathcal{L}}^{b,\varepsilon}(l^\varepsilon)$, together with (9.5), taking l as a test function in (9.3) and using the fact that l_1 is a solution of the unit cell problem yields

$$\lim_{\varepsilon \rightarrow 0} [I_1 + I_2] = \langle |Y_x|^{-1} A(x, y) (\nabla l + \nabla_y l_1), \nabla l + \nabla_y l_1 \rangle_{\Omega_T, Y_{x,K}^*}.$$

Hence, we conclude the convergence of the energy

$$(9.16) \quad \lim_{\varepsilon \rightarrow 0} \langle A^\varepsilon \nabla l^\varepsilon, \nabla l^\varepsilon \rangle_{\Omega_{\varepsilon,K}^*, T} = \langle |Y_x|^{-1} A(x, y) (\nabla l + \nabla_y l_1), \nabla l + \nabla_y l_1 \rangle_{\Omega_T, Y_{x,K}^*},$$

as well as

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} |Y|^{-1} \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon), \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon) \rangle_{\Omega_T, Y} \\ &= \langle |Y_x|^{-1} A(x, y) (\nabla l + \nabla_y l_1), \nabla l + \nabla_y l_1 \rangle_{\Omega_T, Y_{x,K}^*}. \end{aligned}$$

This implies also the strong convergence of the unfolded gradient

$$(9.17) \quad \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon) \rightarrow \chi_{Y_{x,K}^*}(D_x \cdot) (\nabla l + D_x^{-T} \nabla_{\tilde{y}} l_1(\cdot, D_x \cdot)) \quad \text{in } L^2(\Omega_T \times Y).$$

To show the strong convergence in (9.17) we consider

$$\begin{aligned} & \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}^\varepsilon) (\mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon) - \nabla l - D_x^{-T} \nabla_{\tilde{y}} l_1), \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon) - \nabla l - D_x^{-T} \nabla_{\tilde{y}} l_1 \rangle_{\Omega_T \times Y} \\ &= \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon), \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon) \rangle_{\Omega_T \times Y} \\ &\quad - \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon), \nabla l + D_x^{-T} \nabla_{\tilde{y}} l_1 \rangle_{\Omega_T \times Y} \\ &\quad - \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}^\varepsilon) (\nabla l + D_x^{-T} \nabla_{\tilde{y}} l_1), \mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon) \rangle_{\Omega_T \times Y} \\ &\quad + \langle \mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon) \mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}^\varepsilon) (\nabla l + D_x^{-T} \nabla_{\tilde{y}} l_1), \nabla l + D_x^{-T} \nabla_{\tilde{y}} l_1 \rangle_{\Omega_T \times Y}. \end{aligned}$$

Applying the strong convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(A^\varepsilon)$ and $\mathcal{T}_{\mathcal{L}}^\varepsilon(\chi_{\Omega_{\varepsilon,K}^*}^\varepsilon)$ along with the weak convergence of $\mathcal{T}_{\mathcal{L}}^\varepsilon(\nabla l^\varepsilon)$, the convergence of the energy (9.16), and the uniform ellipticity of $A(x, y)$, implies the convergence (9.17). \square

REMARK. Since in $\Omega_{\varepsilon,K}^*$ we have both spatial changes in the periodicity of the microstructure and in the shape of perforations, the l-p unfolding operator $\mathcal{T}_{\mathcal{L}}^{*,\varepsilon}$ is not defined on $\Omega_{\varepsilon,K}^*$ directly and in the derivation of the macroscopic equations we used a local extension of l^ε from $\hat{\Omega}_K^{*,\varepsilon}$ to $\hat{\Omega}^\varepsilon$. The local extension allows us to apply the l-p unfolding operator $\mathcal{T}_{\mathcal{L}}^\varepsilon$ to l^ε . If we have changes only in the periodicity and no additional changes in the shape of perforations, then we can apply the l-p unfolding operator defined in a perforated domain Ω_ε^* directly, without considering an extension from $\hat{\Omega}_\varepsilon^*$ to $\hat{\Omega}_\varepsilon$, and derive macroscopic equations in the same way as in the proof of Theorem 9.4.

10. Discussions. The macroscopic model (9.10) derived from the microscopic description of a signaling process in a domain with locally-periodic perforations reflects spatial changes in the microscopic structure of a cell tissue. The effective coefficients of the macroscopic model describe the impact of changes in the microstructure on the movement (diffusion) of signaling molecules (ligands) and on interactions between ligands and receptors in a biological tissue. The multiscale analysis also allows us to consider the influence of non-homogeneous distribution of receptors in a cell membrane as well as non-homogeneous membrane properties (e.g. cells with top-bottom

and front-back polarities) on the signaling process. The dependence of the coefficients on the macroscopic variables represents the difference in the signaling properties of cells depending on the size and/or position. For example, the changes in the size and shape of cells in epithelium tissues are caused by the maturation process and, hence cells of different age may show different activity in a signaling process. Expanding the microscopic model by including equations for cell biomechanics and using the proposed multiscale analysis techniques we can also consider the impact of mechanical properties of a biological tissue with a non-periodic microstructure on signaling processes.

Techniques of locally-periodic homogenization allow us to consider a wider range of composite and perforated materials than the methods of periodic homogenization. The structures of macroscopic equations obtained for microscopic problems posed in domains with periodic and locally-periodic microstructures are similar. If we consider the microscopic model (9.1)–(9.2) in a domain with periodic microstructure, i.e. $D(x) = \mathbf{I}$ and $K(x) = \mathbf{I}$, where \mathbf{I} denotes the identity matrix, then the macroscopic equations (9.10) with $D(x) = \mathbf{I}$ and $K(x) = \mathbf{I}$ correspond to the macroscopic equations obtained in [36] by considering the periodic distribution of cells and applying methods of periodic homogenization. For some locally-periodic microstructures, e.g. domains consisting of periodic cells with smoothly changing perforations, it is possible to derive the same macroscopic equations by applying periodic and locally-periodic homogenization techniques, see e.g. [37, 38, 49]. However, as mentioned in the introduction, for the microscopic description and homogenization of processes defined in domains with e.g. plywood-like microstructures or on oscillating surfaces of locally-periodic microstructures the techniques of locally-periodic homogenization are essential. Methods of locally-periodic homogenization are applied to analyse microscopic problems posed in domains with non-periodic but deterministic microstructures, in contrast to stochastic homogenization techniques used to derive macroscopic equations for problems posed in domains with random microstructures.

The corrector function l_1 and the macroscopic diffusion coefficient in the macroscopic problem (9.10) are determined by solutions of the unit cell problems (9.11), which depend on the macroscopic variables x . This dependence corresponds to spatial changes in the structure of the microscopic domains. To compute solutions of the unit cell problems (9.11) (and hence the effective macroscopic coefficients and the corrector l_1) numerically approaches from the two-scale finite element method [40] or the heterogeneous multiscale method [1, 2, 26] can be applied. Using heterogeneous multiscale methods one would have to compute the solutions of (9.11) only at the grid points of a discretisation of the macroscopic domain, which requires much lower spatial resolution than computing the microscopic model on the scale of a single cell. Similar approach can be applied for numerical simulations of the ordinary differential equations determining the dynamics of receptor densities, which depend on the macroscopic x and the microscopic y variables as parameters.

REFERENCES

- [1] ABDULLE, A. *The Finite Element Heterogeneous Multiscale Method: a computational strategy for multiscale PDEs*. GAKUTO International Series Math. Sci. Appl., 31 (2009), pp. 133–181.
- [2] ABDULLE, A., WEINAN, E., ENGQUIST, B., VANDEN-ELJNDEN, E. *The heterogeneous multiscale method* Acta Numerica, 21 (2012), pp. 1–87.
- [3] ALEXANDRE, R. *Homogenization and $\theta-2$ convergence*. Proceeding of Roy. Soc. of Edinburgh, 127A (1997), pp. 441–455.

- [4] ALLAIRE, G. *Homogenization and two-scale convergence*. SIAM J Math. Anal., 23 (1992), pp. 1482–1518.
- [5] ALLAIRE, G., DAMLAMIAN, A., HORNUNG, U. *Two-scale convergence on periodic surfaces and applications*, in Proceedings of the International Conference on Mathematical Modelling of Flow through Porous Media, A. Bourgeat et al., eds., World Scientific, Singapore, 1996, pp. 15–25.
- [6] BABAROSIE, C., TOADER, A.-M. *Optimization of bodies with locally periodic microstructures*. Mechanics of Advanced Materials and Structures, 19 (2012), pp. 290–301.
- [7] BABAROSIE, C., TOADER, A.-M. *Optimization of bodies with locally periodic microstructures by varying the periodicity pattern*. Network and Heterogeneous Media, 9 (2014), pp. 433–451.
- [8] BELHADJ M., CANCÈS E., GERBEAU, J.-F., MIKELIĆ, A. *Homogenization approach to filtration through a fibrous medium*. INRIA, 5277, 2004.
- [9] BELYAEV, A.G., PYATNITSKII, A.L., CHECHKIN, G.A. *Asymptotic behaviour of a solution to a boundary value problem in a perforated domain with oscillating boundary*. Siberian Math. J., 39 (1998), pp 621–644.
- [10] BENSOUSSAN, A., LIONS, J.-L., PAPANICOLAOU, G. *Asymptotic analysis for periodic structures*. AMS, Providence, 2011.
- [11] BOURGEAT, A., MIKELIĆ, A., WRIGHT, S. *Stochastic two-scale convergence in the mean and applications*. J Reine Angew. Math., 456 (1994) pp 19–51.
- [12] BRIANE, M. *Homogénéisation de matériaux fibrés et multi-couches*. PhD thesis, Université Paris VI, 1990.
- [13] BRIANE, M. *Three models of non periodic fibrous materials obtained by homogenization*. RAIRO Modél. Math.Anal.Numer., 27 (1993), pp. 759–775.
- [14] BRIANE, M. *Homogenization of a non-periodic material*. J Math. Pures Appl., 73 (1994), pp. 47–66.
- [15] CAPDEBOSQ, Y., PTASHNYK, M. *Root growth: Homogenization in domains with time dependent partial perforations*. ESAIM: COCV, 18 (2012), pp. 856876
- [16] CHAVARRÍA-KRAUSER, A., PTASHNYK, M. *Homogenization approach to water transport in plant tissues with periodic microstructures*. Math. Model. Nat. Phenom. 8 (2013), pp. 80–111.
- [17] CHECHKIN, G.A., PIATNITSKI, A.L. *Homogenization of boundary-value problem in a locally periodic perforated domain*. Applicable Analysis, 71 (1999), pp. 215–235.
- [18] CHENAIS, D., MASCARENHAS, M. L., TRABUCHO, L. *On the optimization of nonperiodic homogenized microstructures*. RAIRO Modél. Math. Anal. Numér., 31 (1997), pp. 559–597.
- [19] CIORANESCU, D., J. SAINT JEAN PAULIN, J. *Homogenization of reticulated structures*, Springer, 1999.
- [20] CIORANESCU, D., DAMLAMIAN, A., GRISO, G. *The periodic unfolding method in homogenization*. SIAM J Math. Anal., 40 (2008), pp. 1585–1620.
- [21] CIORANESCU, D., DONATO, P., ZAKI, R. *The periodic unfolding method in perforated domains*. Portugaliae Mathematica, 63 (2006), pp. 467–496.
- [22] CIORANESCU, D., DAMLAMIAN, A., DONATO, P., GRISO, G., ZAKI, R., *The periodic unfolding method in domains with holes*. SIAM J. Math. Anal. 44 (2012), pp. 718–760.
- [23] COSTA, K.D., TAKAYAMA, Y., MCCULLOCH, A.D., COVELL, J.W., *Laminar fiber architecture and three-dimensional systolic mechanics in canine ventricular myocardium*. American Physiological Society, 1999, pp. H595–H607.
- [24] CIORANESCU D., DAMLAMIAN A., GRISO G., ONOFREI D. *The Periodic Unfolding Method for perforated domains and Neumann sieve models*, Jour. Math. Pures et Appl. 89 (2008), pp. 248–277.
- [25] E., W., ENGQUIST B. *The Heterogeneous Multi-Scale Methods*. Commun. Math. Sci. 1 (2003), pp. 87–132.
- [26] EFENDIEV Y., HOU, T.Y. *Multiscale Finite Element Methods: Theory and Applications*, Springer, New York, 2009.
- [27] FABRITIUS, H.-O., SACHS, CH., TRIGUERO, P.R., RAABE, D. *Influence of structural principles on the mechanics of a biological fiber-based composite material with hierarchical organization: the exoskeleton of the lobster Homorus americanus*. Adv. Materials, 21 (2009), pp. 391–400.
- [28] FATIMA, T., MUNTEAN, A., PTASHNYK, M. *Unfolding-based corrector estimates for a reaction-diffusion system predicting concrete corrosion*. Applicable Analysis, 91 (2012) pp. 1129–1154.
- [29] EVANS L.C. *Partial Differential Equations* AMS, Providence, 2002.
- [30] HORNUNG, J., JÄGER, W. *Diffusion, convection, adsorption and reaction of chemicals in porous media*. J Differ. Equations, 92 (1991), pp. 199–225.
- [31] GRISO, G. *Error estimate and unfolding for periodic homogenization* Asymptotic Analysis 40

- (2004), pp. 269–286.
- [32] GRISO G., *Interior error estimate for periodic homogenization* Anal. Appl. 4 (2006), pp. 61–80.
 - [33] GRISO G. *Error estimates in periodic homogenization with a non-homogeneous Dirichlet condition*. ArXiv, 2014.
 - [34] LUKKASSEN, D., NGUETSENG, G., WALL, P. *Two-scale convergence*. Int. J Pure Appl. Math., 2 (2002), pp. 35-86.
 - [35] LADYZHENSKAJA O.A., SOLONNIKOV V.A., URALCEVA N. N. *Linear and quasi-linear equations of parabolic type*. American Mathematical Society, 1968.
 - [36] MARCINIAK-CZOCZRA, A., PTASHNYK, M. *Derivation of a macroscopic receptor-based model using homogenization techniques*. SIAM J. Math Anal, 40 (2008), 215-237.
 - [37] MASCARENHAS, M. L., POLIŠEVSKI, D. *The warping, the torsion and the Neumann problems in a quasi-periodically perforated domain*. RAIRO Modél. Math. Anal. Numér., 28 (1994), pp. 37–57.
 - [38] MASCARENHAS, M.L., TOADER, A.-M. *Scale convergence in homogenization*. Numer. Funct. Anal. Optimiz., 22 (2001), pp. 127–158.
 - [39] MASCARENHAS, M.L. *Homogenization problems in locally periodic perforated domains*. Asymptotic methods for elastic structures (Proc. of the Internat. Conference, Lisbon, Portugal, 1993), 141–149, de Gruyter, Berlin, 1995.
 - [40] MATACHE, A.-M., SCWAB C. *Two-scale FEM for homogenisation problems*. ESAIM: M2AN, 36 (2002), pp. 537–572.
 - [41] MCCULLOCH A.D. *Cardiac Biomechanics*. The Biomedical Engineering Handbook: Second Edition. Ed. Joseph D. Bronzino, Boca Raton: CRC Press LLC, 2000
 - [42] MURAT, F, TARTAR, L. *H-convergence*. in Topics in the mathematical modelling of composite materials, 21–43, Progr. Nonlinear Differential Equations Appl., 31, Birkhäuser Boston, Boston, MA, 1997.
 - [43] NEUSS-RADU M. *Some extensions of two-scale convergence*. C.R. Acad. Sci. Paris, 322, Serie I (1996), pp. 899-904.
 - [44] NGUETSENG G. *A general convergence result for a functional related to the theory of homogenization*. SIAM J Math. Anal., 20 (1989), pp. 608–623.
 - [45] VAN NOORDEN, T.L., MUNTEAN, A. *Homogenization of a locally-periodic medium with areas of low and high diffusivity*. European J Appl. Math., 22 (2011), pp. 493 –516.
 - [46] ONOFREI, D. *The Unfolding operator near a hyperplane and its application to the Neumann sieve model*. Adv. in Math. Sci. and Appl., vol 16 (2006), pp. 239-258
 - [47] ONOFREI D., VERNESCU B. *Error estimates for periodic homogenization with non-smooth coefficients*. Asymptotic Anal., 54 (2007), pp. 103–123
 - [48] PESKIN, C.S. *Fiber architecture of the left ventricular wall: an asymptotic analysis*. Comm. Pure and Appl. Math., 42 (1989), pp. 79–113.
 - [49] PTASHNYK, M. *Two-scale convergence for locally-periodic microstructures and homogenization of plywood structures*. SIAM Multiscale Modelling and Simulations, 11 (2013), pp 92-117.
 - [50] ROY, D.M., IDORN, G.M. *Concrete Microstructure*. SHRP, Nat.Resear.Coun.Washington, 1993.
 - [51] SCHOCH,W.,HELLER,I.,SCHWEINGRUBER,F.H.,KIENAST,F. *Wood anatomy of central European Species*. 2004, Online version: www.woodanatomy.ch
 - [52] SCHWEERS, E., LOFFLER,F. *Realistic modelling of the behaviour of fibrous filters through consideration of filter structure*. Powder Technol., 80 (1994) pp. 191–206.
 - [53] SHKOLLER, S. *An approximate homogenization scheme for nonperiodic materials*. Comp. Math. Applic., 33 (1997), pp 15–34.
 - [54] WLOKA, J. *Partielle Differentialgleichungen*, Teubner Verlag, Stuttgart, 1982.
 - [55] ZHIKOV, V.V. *On an extension of the method of two-scale convergence and its applications*. Sbornik: Mathematics, 191 (2000) pp. 973-1014.